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1. INTRODUCTION

Engineering compositional disorder into materials has attracted immense interest since the introduction of the high-entropy (HE) concept and the first intentional synthesis of HE alloys in 2004. High-entropy materials (HEMs) represent a new class of solid solutions containing at least five different elements. Their compositional diversity makes them promising as platforms for the development of functional materials. We synthesized new HEMs in a mullite-type structure and present five compounds, i.e., Bi₄(Al₀.₂₅Ga₀.₇₅Fe₀.₂₅Mn₀.₂₅)O₉ and AₓMₓMₙO₈ with variations of A = Nd, Sm, Y, Er, Eu, Ce, and Bi, demonstrating the vast accessible composition space. By combining scattering, microscopy, and spectroscopy techniques, we show that our materials are mixed solid solutions. Remarkably, when following their crystallization in situ using X-ray diffraction and X-ray absorption spectroscopy, we find that the HEMs form through a metastable amorphous phase without the formation of any crystalline intermediates. We expect that our synthesis is excellently suited to synthesizing diverse HEMs and therefore will have a significant impact on their future exploration.

ABSTRACT: High-entropy materials (HEMs) represent a new class of solid solutions containing at least five different elements. Their compositional diversity makes them promising as platforms for the development of functional materials. We synthesized new HEMs in a mullite-type structure and present five compounds, i.e., Bi₄(Al₀.₂₅Ga₀.₇₅Fe₀.₂₅Mn₀.₂₅)O₉ and AₓMₓMₙO₈ with variations of A = Nd, Sm, Y, Er, Eu, Ce, and Bi, demonstrating the vast accessible composition space. By combining scattering, microscopy, and spectroscopy techniques, we show that our materials are mixed solid solutions. Remarkably, when following their crystallization in situ using X-ray diffraction and X-ray absorption spectroscopy, we find that the HEMs form through a metastable amorphous phase without the formation of any crystalline intermediates. We expect that our synthesis is excellently suited to synthesizing diverse HEMs and therefore will have a significant impact on their future exploration.

1. INTRODUCTION

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solution rather than a simple physical mixture of the parent compounds Bi$_2$Al$_2$O$_{10}$, Bi$_2$Ga$_2$O$_{10}$, Bi$_2$Fe$_2$O$_{10}$, and Bi$_2$Mn$_2$O$_{10}$. Comparing the synchrotron XRD data of Bi-HEO and its parent compounds clearly shows significantly shifted reflections, e.g., (001), (121), and (211) as highlighted by gray lines in Figure 2a. The separated (001) reflection shows a symmetrical profile without any observable shoulders (Figure 2b), and the pattern can be excellently fitted with space group Pbam using the Pawley method ($R_{wp}$ = 1.9%, Figure S1a). The refined cell dimensions $a = 7.782(2)$ Å, $b = 8.368(3)$ Å, and $c = 5.848(2)$ Å (Table S1) are close to a simple calculated average of the parent compounds in the respective dimension $a = 7.7967$ Å, $b = 8.3631$ Å, and $c = 5.8472$ Å (Table S2). Rietveld refinement using a structural model with equally distributed Mn, Al, Fe, and Ga over the octahedral and tetrahedral sites results in $R_{wp}$ = 4.1% (Figure 2e) and validates the unit cell dimensions as $a = 7.7853(7)$ Å, $b = 8.3696(7)$ Å, and $c = 5.8423(6)$ Å. The refined average crystallite sizes of 57(3) nm (Rietveld) and 70(9) nm (Pawley) suggest rather small coherently scattering domains, despite a long calcination time of 20 h at 700 °C. This implies that the growth of this material is inherently inhibited. A reason could be the significant microstrain that was refined to $\varepsilon_0 = 0.0030(2)$. This value is an order of magnitude higher compared to those of the parent compounds calcined for the same amount of time (Table S2) and can be expected for such a complex network of differently connected polyhedral environments with differing ionic radii (Al$^{3+}$ 0.53 Å, Fe$^{3+}$ 0.79 Å$^{28}$).

Temperature-dependent XRD studies (Figure 2f) did not reveal any obvious peaks or shoulders on the HEO reflections. Therefore, we can exclude any hidden impurities with differing thermal expansion behavior.

The parent compounds Bi$_2$Fe$_2$O$_{10}$ and Bi$_2$Mn$_2$O$_{10}$ are well-known for their pentagonal frustrated magnetic structures with ordering temperatures of $T_N = 244$ K$^{17}$ and $T_N = 39$ K$^{29}$. However, the NPD patterns of Bi-HEO do not show the appearance of any long-range magnetic order even when the sample is cooled to 1.7 K. This is seen in Figure 2f,g, where the scattering patterns of Bi-HEO and Bi$_2$Fe$_2$O$_{10}$ are compared in the relevant Q-range. Here, Bi$_2$Fe$_2$O$_{10}$ shows strong magnetic reflections at $Q = 1.3$ Å$^{-1}$ already at 100 K. In the same region, Bi-HEO shows a magnetic diffuse signal (arrow, Figure 2f) at 1.7 K, indicating a short-range magnetic order. Such a diffuse signal was also reported for Bi$_2$Fe$_2$O$_{10}$ at 280 K, which sluggishly grew into well-defined magnetic Bragg peaks at lower temperatures.$^{30}$ The absence of any long-range magnetic order in Bi-HEO provides strong evidence that our material represents a solid solution rather than a mixture of separated phases. Otherwise, the strong magnetic signals of Bi$_2$Fe$_2$O$_{10}$ or Bi$_2$Mn$_2$O$_{10}$ would appear below their respective ordering temperatures. A Pawley fit of the 300 K neutron data (Figure S3a) yields an $R_{wp}$ = 4.4% and unit cell dimensions $a = 7.77(1)$ Å, $b = 8.37(1)$ Å, and $c = 5.85(1)$ Å in accordance with the XRD refinements (Table S1). However, magnification of the background shows at least three peaks (Figure S3b), which clearly cannot be explained with the Pbam symmetry. Since we did not observe these in the XRD data, we assume that they arise from one or more minor impurity phases in low concentrations containing weak scatterers, e.g., some Al/Ga oxides or hydroxides. Nevertheless, Rietveld refinement of the neutron data (Figure 2h) yields a reasonable $R_{wp}$ = 6.2% and matches the refined $a$, $b$, $c$, and microstrain $\varepsilon_0$ of the XRD data (Table S1). For both data sets, some atomic displacement
Parameters (ADPs) refine to quite high values ($B_{\text{iso}} = 0.4(2) - 2.5(2)$ Å$^2$), which we relate to the substantial chemical disorder around the sites rather than to abnormally large thermal vibrations. As high-entropy materials contain many elements with different sizes, a slight displacement from their average atomic position in the crystal is expected. In a Rietveld refinement, such a static displacive disorder would add to the displacement caused by thermal vibrations and therefore lead to higher $B$-values. To explore the local structure of Bi-HEO, we collected X-ray total scattering data and performed PDF analysis. Bi-HEO clearly shows shorter correlation lengths compared to its parent compounds (Figure S4). We attribute this to microstrain rather than to size effects since the Rietveld analysis suggests that the average crystallite size is too big to cause any significant damping of the PDFs in the $r$-range available from the X-ray rapid-acquisition (RA)-PDF experiments. In Figure S5, we compare a calculated PDF with $1/4$ Bi$_2$Fe$_4$O$_9$, $1/4$ Bi$_2$Al$_4$O$_9$, $1/4$ Bi$_2$Ga$_4$O$_9$, and $1/4$ Bi$_2$Mn$_4$O$_{10}$ simulating a simple physical mix of the parent compounds (details in the Methods Section) with the experimental PDF of Bi-HEO. The PDFs are almost identical until ca. 8.4 Å, which is on the order of the dimension of the unit cell. This means that the local structure of the material is well described by a simple mixture of the parents. However, beyond that distance, the PDFs clearly start to differ from each other as the longer range order is more affected by the differing lattice parameters for the parent phases in the simulated physical mixture. The whole range of the Bi-HEO PDF ($r = 1 - 60$ Å) is well described ($R_{wp}$ = 18.9%, Figure S6) by the mullite structure. However, individual fits of the $r$-ranges from 1 to 25 Å and from 25 to 60 Å lead to significantly improved fits with $R_{wp}$ values of 15.7 and 11.6% (Figure S7, Table S1). This improvement relates to the increased PDF peak broadening in the high $r$-region and considerably higher refined values for the ADP (e.g., $B_{\text{iso}}$ of Bi doubles from 1.1 to 2 Å$^2$, Table S1). Consistent with the results from the Rietveld analysis, these high ADP values likely represent static disorder around the sites and could be caused by a mismatch of the different ion sizes in the lattice leading to significant microstrain in the material.26

To further evaluate the local structure, we collected the infrared (IR) (Figure 2c) and Raman (Figure 2d) spectra of Bi-HEO and its parent phases. The IR spectrum of mullite-type materials can be divided into three main regions: 400–450 cm$^{-1}$, 650–750 cm$^{-1}$, and 1000–1600 cm$^{-1}$.
600 cm\(^{-1}\), M-O stretching modes of MO\(_6\) octahedra; 600–800 cm\(^{-1}\), low-energy stretching vibrations of MO\(_4\) tetrahedra and M–O–M bending vibrations of the double-tetrahedral units; and 800–900 cm\(^{-1}\), high-energy stretching vibrations of MO\(_4\) tetrahedra.

The latter are characteristic and distinguishable for the parent compounds, e.g., Bi\(_2\)Fe\(_4\)O\(_9\); 812 cm\(^{-1}\), Bi\(_2\)Ga\(_4\)O\(_9\); 848 cm\(^{-1}\), and Bi\(_2\)Al\(_4\)O\(_9\); 922 cm\(^{-1}\) (Figure 2c). In contrast, the IR spectrum of Bi-HEO shows very broad features in all three regions. Two seemingly broad bands with maxima at 850 and 888 cm\(^{-1}\) show up in the characteristic MO\(_4\) stretching region, but a closer look reveals the overlap of several distinct peaks (Figure S8). Such a spectrum would be expected assuming a broad distribution of bond lengths and different combinations of vibrating oscillators in the double tetrahedral and/or pyramidal units. Considering only the M-O-M bending mode, we could expect 10 different combinations of oscillators in Bi-HEO, e.g., Mn–O–Mn, Fe–O–Fe, Al–O–Al, Ga–O–Ga, Ga–O–Al, Ga–O–Fe, Ga–O–Mn, Fe–O–Al, Fe–O–Mn, and Al–O–Mn. A clear band splitting in the IR spectra for Bi\(_2\)Al\(_{1-x}\)Fe\(_x\)O\(_9\) solid solutions was also observed by Voll et al.\(^{31}\)

While the IR spectrum of Bi-HEO mainly resembles the ones of the Bi\(_2\)M\(_{4}\)O\(_9\) (M = Al, Ga, and Fe) parents, the Raman spectrum of Bi-HEO (Figure 2d) is quite similar to the one of Bi\(_2\)Mn\(_4\)O\(_{10}\). However, compared to Bi\(_2\)Mn\(_4\)O\(_{10}\) our HEO shows a significantly broader band distribution especially around the maxima at 608 and 644 cm\(^{-1}\). In Bi\(_2\)Mn\(_4\)O\(_{10}\) these bands are assigned to stretching modes of MnO\(_6\) in the ab-plane and reported to be 605 and 647 cm\(^{-1}\), respectively.\(^{33}\) We assume that the broad distribution near these two mode frequencies in Bi-HEO arises due to local changes around the MnO\(_6\) octahedra, significantly affecting its vibrational properties.

The elemental distribution of the various elements was mapped by STEM-EDX. The representative micrographs shown in Figure 2i and Figure S9 demonstrate that the particles have a broad size distribution ranging between ca. 100 and 2000 nm and contain all of the expected elements: Bi, Al, Fe, Ga, Mn, and O. Apart from well-mixed regions (Figure 2i and Figure S9a), some particles clearly show segregation (Figure S9b,c). Especially Al and Mn show some degree of anticorrelation, i.e., they prefer not to mix as can be seen in Figure S9c. Some particles show Ga-rich areas, which might be...
related to the fact that Ga$_2$O$_3$ was used as the Ga source in the synthesis. This may lead to regions of higher heterogeneity due to slower reaction kinetics. Bi, Fe, and O seem to be well mixed in all particles. The observed heterogeneities could be small nanoparticles in the form of, e.g., binary oxides or hydroxides. Due to their low scattering power compared to Bi-containing compounds, such contributions would vanish in XRD, especially when the particles are small and in low concentration, e.g., in the range of 1–3%. This would be following the unexplained reflections observed in NPD, where Al ($b = 3.449(5)$ barn)\textsuperscript{34} and Ga ($b = 7.288(2)$ barn)\textsuperscript{34} impurities would be observable next to Bi ($b = 8.532(2)$ barn),\textsuperscript{34} Fe ($b = 9.45(2)$ barn),\textsuperscript{34} and Mn ($b = -3.73(2)$ barn)\textsuperscript{34} compounds due to their distinguishable coherent scattering lengths $b$. However, we expect these observed impurities to be minor phases in the material (<3%).

\subsection*{2.2. A-Site-Substituted A$_3$Mn$_4$O$_{10}$ High-Entropy Oxides.} To further map out the synthetic accessibility of HEOs in the mullite-type system, we synthesized another four compounds, substituting the A-site of the Bi$_3$Mn$_4$O$_{10}$ parent (Figure 1b, overview of compounds in Figure 1f). Here, we chose lanthanides of various sizes, as well as Bi$^{3+}$ and Y$^{3+}$. Three of the compounds contain Bi, and only one element (A) is exchanged A-HEO (A = Er, Eu, and Ce and HEO = Y, Sm, Nd, and Bi). The fourth compound RE-HEO, however, does not contain Bi, thus eliminating a possible influence of its stereoeactive lone electron pair. Like Bi-HEO, the materials were synthesized by the sol–gel method. The phase purity and structure of the compounds were confirmed by synchrotron XRD (Figure 3a) and Rietveld refinements (Figures S10–S13, Table S3), shown for RE-HEO in Figure 3d. The data can be fitted with a structural model assuming a statistical distribution of the A-site cations and $Pbam$ symmetry yielding $R_w$ values between 2 and 2.9%. With substitution, the reflections shift (Figure 3b). The refined unit cell parameters of the new HEOs are compared to the respective parent compounds A$_3$Mn$_4$O$_{10}$ (A = Er,\textsuperscript{35} Y,\textsuperscript{35} Eu,\textsuperscript{36} Sm,\textsuperscript{36} Nd,\textsuperscript{37} and Bi) in Figure 3c (volume) and Figure S14 ($a$, $b$, and $c$). The values for Ce$_3$Mn$_4$O$_{10}$ are not given since this compound has not been reported so far, and our synthesis attempts were not successful. With the introduction of bigger ions into the crystal lattice, $a$, $b$, and $c$ increase, as expected. The introduction of Bi$^{3+}$ leads to a clear drop in $b$ following the expected trends of the parent compounds (Figure S14b). The average crystallite sizes were refined to be 21(1)–37(2) nm for Er-HEO, Eu-HEO, and Ce-HEO and significantly bigger for RE-HEO with 221(17) nm, which was heated to a higher temperature (900 °C vs 700 °C). At the same time, the microstrain ($\varepsilon_M$) is an order of magnitude smaller (Table S3) than those in Er-HEO, Eu-HEO, and Ce-HEO, which might be related to the elimination of the stereochemically active lone electron pair of Bi in RE-HEO. The microstrain of all A-site-substituted HEOs, in general, is significantly smaller than the one determined for Bi-HEO (see Tables S1 and S3), and the reduction in correlation lengths in the PDFs is not observed (Figure S15). This is understandable considering the complex connectivity of substituted polyhedral sites in Bi-HEO compared to the substitution of lanthanides with close ionic radii in the A-site-substituted HEOs.

When comparing RE-HEO to Bi$_3$Mn$_4$O$_{10}$, the substitution of Bi$^{3+}$ by RE$^{3+}$ leads to strong shifts of vibrational frequencies in the IR (∼25 cm$^{-1}$) and Raman spectra (∼41 cm$^{-1}$), as can be seen in Figure 3e.f. This can be explained by structural changes involving the coordination polyhedra, e.g., bond lengths and polyhedral distortions, which these spectroscopy methods are very sensitive to.\textsuperscript{38} We can follow the effect of the A-site substitution in our HEOs by the frequency shifts and induced peak broadening although the observable modes mainly involve the Mn coordinations.\textsuperscript{33} As an example, the Raman spectrum of Bi$_3$Mn$_4$O$_{10}$ shows two signals at 187 and 199 cm$^{-1}$, whereas we observe only a single sharp feature at
217 cm$^{-1}$ for RE-HEO, which does not contain Bi. However, Er-HEO, Eu-HEO, and Ce-HEO, which contain Bi, also show broad asymmetric peak shoulders in that region. The same asymmetry can be observed in the higher energy region ($\sim$600 cm$^{-1}$), which we attribute to the influence of the Bi lone electron pair. In general, the spectral features observed for the A-site-substituted HEOs are significantly sharper than those for Bi-HEO, which highlights the large influence of the B-site substitution on the peak broadening (Figure 2d).

STEM-EDX analysis reveals a homogeneous distribution of the RE elements in the particles, as can be seen in Figure 3g (RE-HEO) and Figure S16 (other compositions). The materials possess an interesting open network structure of sintered nanoparticles with individual sizes of ca. 20 nm (Er-HEO, Eu-HEO, and Ce-HEO) and 150 nm (RE-HEO) matching the results of the Rietveld refinements (Table S3). The particles are highly crystalline (Figure S17e−h) as atomic column contrast can be observed in the high-resolution transmission micrographs (HRTEM) and Fast Fourier transformations (FFT) of selected areas (Figure S17, yellow squares) that clearly show sharp spots indicative of a highly ordered structure.

2.3. Formation of High-Entropy Mullite-Type Oxides.
In contrast to most reports on the synthesis of HEOs, we use the sol–gel method instead of the classical solid-state approach. As described in detail in the Methods Section, metal nitrates are dissolved in stoichiometric amounts in diluted HNO$_3$ together with D-sorbitol, which is used as the complexing agent. From the solutions, a well-mixed precursor is thus produced, which is afterward calcined in air at 700 °C for 20 h. This temperature is sufficient for the HEOs to form, except for RE-HEO, which is heated to 900 °C for 20 h to yield the HEO product. This indicates that the use of Bi greatly promotes HEO crystallization, possibly due to its volatile nature at higher temperatures.39

To gain further insights into the crystallization mechanism, we followed the formation of the HEOs by in situ XRD and found that all compounds form through an amorphous intermediate without the formation of any crystalline phases, as seen in Figure 4a−c (Bi-HEO composition) and Figure S18 (other compositions). In all cases, the data show that the precursor is amorphous, except for Ga$_2$O$_3$, which is seen as a crystalline phase. The formation of the HEOs passes through two stages. First, the amorphous precursor transforms into a second amorphous phase of a different structure at ca. 300 °C, which then crystallizes directly into the respective HEO. The low temperature of the first transition indicates that the precursor powder consists of components containing nitrates/organics as these have typical decomposition temperatures in this temperature range. The transformation of the amorphous precursor into the amorphous intermediate is accompanied by a drastic change in position of the main scattering from $\sim$1.3 to $\sim$2.2 Å$^{-1}$ (see Figure 4a). The increase in the scattering angle implies that a larger phase transforms into a smaller one, which fits into the picture of a decomposing phase through the loss of its contained organic components. In the case of Bi-HEO, the Bragg peaks from Ga$_2$O$_3$ disappear when the mullite-type structure forms. While Bi-HEO needs a reaction time of about 14 min at 720 °C (52 min in total) to form, Er-HEO, Eu-HEO, and Ce-HEO start to crystallize already at 670 °C. The first reflections of RE-HEO can be observed after 22 min at 720 °C (60 min in total); however, even after 100 min at 720 °C (140 min in total), the intermediate phase did not completely convert into the mullite-type structure (Figure S18d). In addition, we followed the changes in the oxidation states of Mn and Fe during the formation of Bi-HEO via X-ray absorption spectroscopy (XAS) (Figure 4d−f). Both Mn and Fe oxidize between ca. 250 and 420 °C as determined by linear component analysis and demonstrated in Figure 4f. Remarkably, they oxidize exactly at the same time, as the curves of both species almost overlay (Figure 4f). Thereafter, no significant change can be observed, suggesting that the transition of the precursor into the amorphous intermediate lays the foundation for the later crystallization of the final products. Of note, the crystallization of Bi-HEO in the coupled XAS/XRD experiments is observed at ca. 640 °C (Figure S19) compared to 720 °C in the XRD experiments (Figure 4a). This can be explained by the significant difference in heating rates (XAS/XRD = 1.7 K min$^{-1}$; XRD = 20 K min$^{-1}$) likely changing the kinetics of the crystallization process. To follow the structural changes during the phase formation of the HEOs in more detail, we performed sequential Rietveld refinements on the diffraction patterns collected at 720 °C. For all materials, we observe a reduction of unit cell volume during growth, as presented in Figure S20c. This surprising result is in line with earlier studies on Bi$_2$Fe$_3$O$_9$40 and suggests that this volume contraction is a general phenomenon during the growth of mullite-type compounds as it seems to be independent of chemical composition. The crystallites of the A-site-substituted HEOs grow very sluggishly, as can be seen in Figure S20b. We have fitted the crystallite growth curves using a saturation growth rate equation of the form

$$y = \frac{a - x - c}{b + (x - c)}$$

with fitting parameters $a$, $b$, and $c$. Here, $a$ represents the saturation crystallite size, whereas $c$ can be seen as the starting point of the crystal growth. The predicted saturation sizes by the fitted model and the refined crystallite sizes of the heated HEOs (700 °C for 20 h) show excellent agreement, as can be seen in Table S4. As an example, for Ce-HEO, we observe a crystallite size of 35(1) nm after heating at 700 °C for 20 h compared to a predicted size by the model of 30.3(1) nm. We take this result as a confirmation that the crystal growth of our calcined HEOs have reached saturation in the nanoregime although calcination at high temperatures. This observation is surprising and offers the possibility to control the sizes of the crystallites solely through the adjustment of calcination temperature. In a broader context, synthesizing a material in a high-entropy version might be a route to produce nanomaterials of structures usually only accessible at high temperatures. The combination of the sluggish crystal growth and the saturation growth behavior in the nanoregime suggests that the chemical gradients in the material are very small, i.e., the precursor is highly homogeneous. The high homogeneity combined with the high number of different chemical constituents may lead to a slow atomic diffusion rate due to low chemical potentials. From the local perspective of an atom, the surrounding potential might not differ much, and as such, material diffusion might be inhibited. As soon as the temperature is high enough, the material crystallizes very suddenly, and many grains are formed at the same time. This in turn could explain the lack of Ostwald ripening as all crystallites would be of ca. the same size. However, this hypothesis is of speculative nature. The exact mechanism of the sluggish growth behavior is out of the scope of this paper.
and will need further evaluation. To investigate if any impurities would form upon cooling, we continued to collect data in the XRD experiments after the HEOs had formed at 720 °C and the heating process was stopped. None of the $\mathrm{A}_x\mathrm{Mn}_y\mathrm{O}_{z+2}$ HEOs show any phase separation as can be seen in Figure S21. For Bi-HEO, however, we observe the appearance of new reflections (Figure S22) during cooling. While the pattern at 720 °C can be described excellently with the Pbam lattice symmetry (Figure 4c), the patterns below 690 °C clearly cannot. This led us to repeat the synthesis several times, and although every synthesis attempt showed the formation of Bi-HEO as the main phase, the formation of this impurity phase was regularly observed in varying amounts (see Figure S23). We identified this phase as sillenite-type $\mathrm{Bi}_2\mathrm{Ga}_2\mathrm{O}_7$ which can be dissolved by diluted $\mathrm{HNO}_3$ (Figure S23, inset). The formation of $\mathrm{Bi}_2\mathrm{Ga}_2\mathrm{O}_7$ may be related to the insolubility of $\mathrm{Ga}_2\mathrm{O}_3$ in the precursor mixture leading to Ga-rich regions in the material (see the section on B-Site-Substituted $\mathrm{Bi}_5(\mathrm{Al}_{0.25}\mathrm{Fe}_{0.25}\mathrm{Ga}_{0.25}\mathrm{Mn}_{0.25})\mathrm{O}_3$).

3. CONCLUSIONS

With this work, we report the first successful synthesis of HEOs in a mullite-type structure and thereby introduce a completely new family for the design of HEOs with potentially exciting properties. To demonstrate the wide accessible compound space, we present five samples with varying compositions spanning different ionic sizes. Applying the HE concept to mullite-type materials offers a tremendous platform to study and tune a diversity of desired functional properties by exchange of either the A- or B-sites of these compounds. As an example, we demonstrate a change of magnetic behavior for B-site-substituted Bi-HEO using NPD. Further investigations of the magnetic properties and structures of $\mathrm{A}_x\mathrm{Mn}_y\mathrm{O}_{z+2}$ HEOs hold great potential for new discoveries.

We revealed the wide synthetic accessibility of only two compounds in the mullite-type family. However, the compound space is immense, and we expect that the synthesis of other HEOs of this family is possible including, e.g., $\mathrm{RECrGeO}_x$, $\mathrm{REAlGeO}_x$, $\mathrm{REFeMnO}_x$, and $\mathrm{RECrMnO}_x$. As such, we created a potentially large space of interesting new HEOs, which are worth exploring toward a variety of applications.

All compounds studied here crystallize via a two-step process. First, the amorphous precursor transforms into an amorphous phase with a different structure, from which the HEOs crystallize directly without the formation of any crystalline oxides as intermediates, proving that the homogeneity is maintained throughout the synthesis. This observation is surprising as thermodynamically stable compounds are known to exist in this temperature regime, e.g., several polymorphs of manganese oxide. This let us conclude that the pathway through the intermediate amorphous phase allows bypassing possible thermodynamic wells in the energy landscape of formation enthalpy. This might explain why there are so few HEO structure-types known so far as the traditional solid-state route has mainly been used for their synthesis. We expect that the presented synthesis is very effective also in the fabrication of HEOs of other structure types, and many reports of new materials with interesting properties will follow.

4. METHODS

4.1. Syntheses. Syntheses were performed by using the sol–gel method. First, an acidic solution with 26 mL of deionized $\mathrm{H}_2\mathrm{O}$, 3 mL of conc. $\mathrm{HNO}_3$ (TraceSELECT for trace analysis, ≥99.0%, Honeywell), and 1.29 g of $\mathrm{D-sorbitol}$ (≥98%, Sigma-Aldrich) is produced. After that, a total of 7.5 mmol of the respective metal nitrates in the given stoichiometry is dissolved. Heating the solution at 90 °C in a glycerin bath for 2 h leads to a gel, which is subsequently heated in air at 250 °C for 18 h in an oven to obtain a solid “foam” precursor. This precursor is afterward ground and calcined in a porcelain crucible in air at 700 °C for 20 h (900 °C for 20 h; RE-HEO) to form the nanocrystalline products. We did not observe any reaction of the materials with the porcelain crucible under the given conditions. All reagents were of analytical grade and used as received: bismuth(III) nitrate pentahydrate ($\mathrm{AC}$ reagent, ≥98.0%, Sigma-Aldrich), aluminum(III) nitrate nonahydrate ($\mathrm{AC}$ reagent, ≥98.0%, Sigma-Aldrich), gallium(III) oxide (99.99+% metals basis, Sigma-Aldrich), manganese(II) nitrate tetrahydrate ($\mathrm{AC}$ reagent, ≥98.0%, Sigma-Aldrich), petalite (99.9% Y, abcr), samarium(III) nitrate hexahydrate (99.99% trace metals basis, Sigma-Aldrich), neodymium(III) nitrate hexahydrate (99.9% trace metals basis, Sigma-Aldrich), erbium(III) nitrate pentahydrate (99.9% trace metals basis, Sigma-Aldrich), and ammonium cerium(IV) nitrate (for analysis ESMURE ACS, Reag. Ph Eur, Sigma-Aldrich).

4.2. Structural Characterization. Room temperature X-ray diffraction and total scattering data were collected at P02.1 at the PETRA III/DESY in Hamburg, Germany. Measurements have been carried out on samples in 1 mm Kapton capillaries using a wavelength of 0.2074 Å and a PerkinElmer XRD1621 CN3 area detector with a sample detector distance of about 600 mm (XRD, center geometry) and 270 mm (PDF, corner geometry). The exposure times were 60 and 300 s for XRD and PDF measurements. Data were integrated using pyFAI and the PDFs obtained using xpDFS Suite with a Qmax and rPoly of 21.5 Å and 0.9, respectively. The scattering pattern of an empty Kapton capillary was used for background subtraction.

In situ X-ray diffraction data were collected at P21.1 at PETRA III/DESY in Hamburg, Germany using a wavelength of 0.1220319 Å and a PerkinElmer XRD1621 area detector with a sample detector distance of about 600 mm (XRD, center geometry) and about 270 mm (PDF, corner geometry). The exposure times were 60 and 300 s for XRD and PDF measurements. Data were integrated using pyFAI and the PDFs obtained using xpDFS Suite with a Qmax and rPoly of 21.5 Å and 0.9, respectively. The scattering pattern of an empty Kapton capillary was used for background subtraction.

For the crystallization experiments, the samples were heated from room temperature to 720 °C and kept there for ca. 2 h. After that, the samples were cooled naturally while the XRD patterns were still collected. In the heating/cooling experiments on crystalline $\mathrm{Bi}_5(\mathrm{Al}_{0.25}\mathrm{Fe}_{0.25}\mathrm{Ga}_{0.25}\mathrm{Mn}_{0.25})\mathrm{O}_3$, the sample was heated from room temperature to 720 °C and cooled back to room temperature with a ramp of 20 K min$^{-1}$. The procedure was performed twice in total.

Neutron powder diffraction data on a $\mathrm{Bi}_5(\mathrm{Al}_{0.25}\mathrm{Fe}_{0.25}\mathrm{Ga}_{0.25}\mathrm{Mn}_{0.25})\mathrm{O}_3$ compound were collected with a high-resolution powder neutron diffractometer HRPT at the SINQ Spallation Neutron Source at PSI, Switzerland. The instrument was used in high intensity mode and with the neutron wavelength 1.8857 Å. A sample was closed in a vanadium container, and the data were collected at 1.7 and 300 K. The NPD data of Bi$_2$FeO$_3$ were collected on samples at 100 and 300 K on the NOMAD instrument at the Spallation Neutron Source (SNS) at Oak Ridge National Laboratory, Tennessee, USA.

Refinements of XRD, PDF, and NPD data were performed using Topas Academic. The instrumental resolution functions (IRF) were determined by fitting against standards of LaB$_6$ (room temperature XRD and PDF, P02.1), Ni (in situ XRD, P21.1), and Na$_2$Ca$_3$AlF$_{14}$ (NPD, HRPT). XRD and NPD data were fitted by the Pawley and Rietveld method. Rietveld and PDF refinements of the X-ray data.
were performed by using a model with a statistical distribution of the nominal compositions over the respective crystallographic sites. Refined parameters (Table S1) are the scale factor, lattice parameters, atomic displacement parameters (ADP), atomic coordinates, one Lorentzian profile parameter to calculate \( LVol-IB \) for the apparent average crystallite size, and one Gaussian profile parameter for calculating \( e_0 \) as the microstrain. In the refinements of the NPD data, the occupancy of the individual elements was refined on the 4e and 4h sites, however constrained to the weighted nominal composition and alternatingly with the respective ADPs. The temperature-dependent XRD data of the crystallized Bi-HEO was refined sequentially using the Pawley method. The time-dependent XRD data collected at 720 °C during the formation of Bi-HEO, Eu-HEO, Ce-HEO, and Er-HEO was refined using the Rietveld method. We performed sequential refinements starting from the last pattern of the most crystalline sample and going backward toward the least crystalline sample. The resulting crystal growth curves for Eu-HEO, Er-HEO, and CE-HEO were fitted using a saturation growth rate equation of the form \( \text{growth rate} = a(x - c)/b + (x - c) \) (eq 1) with fitting parameters \( a, b, \) and \( c \). The PPDFs were refined in the ranges 1-60, 1-25, and 25-60 Å. Refined parameters were the scale factor, lattice parameters, atomic displacement parameters (ADP), atomic coordinates, and sp-diameter. The PPDFs of the “simulated physical mixtures” were calculated by PDFgui using the experimental instrumental parameters and assuming each 1/4 for Bi\(_2\)Al\(_2\)O\(_4\), Bi\(_2\)Ga\(_2\)O\(_4\), Bi\(_2\)Fe\(_2\)O\(_4\), and Bi\(_2\)Mn\(_2\)O\(_4\). The total scaling factor for the respective parent compounds.

4.3. Spectroscopy. Raman spectra were measured on an inVia Qontor confocal Raman microscope using the Wire v. 5.4 software (Renishaw, Gloucestershire, UK) with a CCD detector. A laser excitation of 785 nm (yellow laser), laser power of 10 mW, dispersed by a 1200 lines mm\(^{-1}\) grating, and exposure time for 10 s and 10 accumulations were used for Bi\(_2\)(Al\(_{2}\)Fe\(_{0.5}\)Ga\(_{0.5}\)Mn\(_{25}\))O\(_{4}\), Bi\(_2\)Al\(_{2}\)O\(_{4}\), Bi\(_2\)Ga\(_{2}\)O\(_{4}\), Bi\(_2\)Fe\(_{2}\)O\(_{4}\), and Bi\(_2\)Mn\(_{2}\)O\(_{4}\). Data for the A-site-substituted A\(_{1/2}\)Mn\(_{1/2}\)O\(_{10}\) HEOs were collected using a laser excitation of 532 nm (green laser), laser power of 5 mW, dispersed by a 1800 lines mm\(^{-1}\) grating, and exposure for 10 s and 5 accumulations. Data were collected between 100 and 3200 cm\(^{-1}\).

Infrared spectra were measured with 64 scans from 374 to 4000 cm\(^{-1}\) using a Bruker Alpha FT IR spectrometer in attenuated total reflectance (ATR) detection. Coupled in situ XRD and XAS data were collected at Swiss Norwegian Beamline (SNBL) BM31, ESRF, Grenoble. The solid precursors were placed with quartz wool plugs in both side opened quartz capillaries (0.8 mm diameter). The capillary was glued into an in situ flow-through reactor available at the beamline that was connected to gas lines by Swagelok connections. \(^{54,55}\) Synthetic air was flowed through the capillary with a flow rate of 5 mL min\(^{-1}\). The capillary was heated with a gas blower at a constant heating ramp of 1.7 K min\(^{-1}\), while XRD and XAS data were collected consecutively.

XAS-XRD data were collected in a quasi-simultaneous manner using a Si(111) double crystal monochromator (DCM). XANES spectra were acquired at the Fe K-edge (7112 eV) and Mn K-edge (6539 eV) in fluorescence mode with a Vortex one-element silicon drift detector (SDD) and FalconX electronics for extremely high count rate performance. In situ X-ray diffraction data were collected with a wavelength of 0.338591 Å, using a 2D DEXELA CMOS (complementary metal-oxide semiconductor) pixel detector with a sample detector distance of ca. 420 mm.

XANES spectra were treated with the ATHENA\(^{56}\) and LARCH program package.\(^{57}\) The linear combination analysis (LCA) was performed with LARCH using the initial precursor spectra and the final oxide spectra as the respective components for the LCA fits.

4.4. Electron Microscopy and STEM-EDX Analysis. The nanoparticle (NP) samples were suspended in ethanol and subsequently drop-casted onto Cu 300 mesh grids with 2 nm carbon (Quantifoil). High-resolution transmission electron microscopy (HR-TEM) images were collected on an FEI Talos FX-200 equipped with a Ceta 16 M camera with samples mounted in a double tilt holder. The FIJI software was used to crop out quadratic regions of interest around selected NPs oriented with their major zone axes aligned with the optical axis.\(^{58}\) These smaller cropped regions were apodized using a Hann window and, subsequently, fast Fourier transformed (FFT) using HyperSpy.\(^{59}\) The resulting power spectra of the FFT images are displayed on a log scale with the minimum value set to display the 30th percentile. Scanning transmission electron microscopy (STEM) images were collected using the same electron microscope in STEM mode and by using a 10.5 mrad converged electron probe. The collection angle of the STEM high-angle annular dark-field (HAADF) images was 60–200 mrad. STEM energy-dispersive X-ray spectroscopy (EDX) data sets were collected using the ChemiSTEM Gen1 system installed on the Talos Microscope. Data analysis and plotting were subsequently performed using HyperSpy and FIJI. Specifically, the elemental distribution maps presented in the manuscript are produced by using HyperSpy to rebin the collected data spatially to increase signal-to-noise ratio and subsequently subtract background counts by performing linear interpolation around the relevant elemental lines. Finally, the color maps representing the spatial distribution of different elements shown in the manuscript are produced by using FIJI.

**ASSOCIATED CONTENT**

Data Availability Statement
All relevant data within the article are available from the corresponding authors on reasonable request.

Supporting Information
The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acs.chemmater.3c01830.

Rietveld and Pawley refinements of synchrotron X-ray diffraction data from all compounds including the results in a tabulated form; temperature-dependent lattice parameters of Bi-HEO; Pawley refinement of neutron diffraction data for Bi-HEO; pair distribution functions of all synthesized HEOs and parent compounds Bi\(_2\)M\(_2\)O\(_9\); Gaussian fits of FTIR data; comparison of lattice parameters from synthesized HEOs and parent compounds; STEM-EDX and HAADF images of the synthesized HEOs; in situ diffraction data of the formation of the HEOs; crystallite size growth curves along with volume evolution during the formation of the HEOs; comparison between XRD patterns of products formed during the in situ heating experiments and after cooling; reproducibility test of Bi-HEO (PDF).

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Author Contributions
A.K. and K.M.O.J. planned the project. A.K. designed, synthesized, characterized (Raman and IR) the compounds, and analyzed the Raman, IR, XRD, PDF, and NPD data. A.K., N.L., J.K.M., and S.L.S. acquired the in situ XRD data at beamline P21.1. N.L. helped with data visualization. E.D.B. acquired and analyzed the TEM data with the help of R.L. D.S. acquired the NPD data at PSI, HRPT, Switzerland. R.K.P. acquired and analyzed the XANES data. The manuscript was written by A.K. and revised by K.M.O.J. with the help of all other authors.

Notes
The authors declare no competing financial interest.

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REFERENCES