Rigidity and non-existence results for collapsed translators

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RIGIDITY AND NON-EXISTENCE RESULTS FOR COLLAPSED TRANSLATORS

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Abstract. We prove a rigidity result for mean curvature self-translating solitons, characterizing the grim reaper cylinder as the only finite entropy self-translating 2-surface in $\mathbb{R}^3$ of width $\pi$ and bounded from below. The proof makes use of parabolicity in a weighted setting applied to a suitable universally $L$-superharmonic function defined on translators in such slabs.

1. Introduction

In this paper, a translator $\Sigma^2 \subseteq \mathbb{R}^3$ is a smooth complete connected surface in $\mathbb{R}^3$ which evolves according to the mean curvature flow as unit speed translation in the direction of the positive $x_3$-axis.

Therefore $\Sigma^2$ satisfies the following equation

$$H = \langle e_3, \nu \rangle \nu,$$

where $H = H\nu$ is the mean curvature vector and $\nu$ is the unit normal vector field on $\Sigma^2$. See Section 2 for some examples, with references.

J. Spruck and L. Xiao proved in [SX20] that every complete mean convex translator in $\mathbb{R}^3$ is actually convex. X. J. Wang proved in [Wan11] that any convex translator in $\mathbb{R}^{n+1}$ which is not an entire graph lies in a vertical slab of finite width (i.e. is collapsed). Hence, using the grim reaper cylinder as a barrier, the following result can be shown (see e.g [HIMW19], [BLT20]).

Proposition 1.1. Let $\Sigma^2 \subseteq \mathbb{R}^3$ be a translator, which is a graph $\{(x, y, u(x, y))\}$ and contained in a vertical slab of width $\pi$. Then, up to rigid motions, $\Sigma$ is a standard grim reaper cylinder $\Gamma \times \mathbb{R}$.

Recall that D. Hoffman, T. Ilmanen, F. Martín, and B. White have actually obtained in [HIMW19] the complete classification of complete graphical translators in $\mathbb{R}^3$. These are either tilted grim reaper cylinders, $\Delta$-wings, or the bowl soliton (the unique entire graph over $\mathbb{R}^2$). Both the (non-trivially) tilted grim reaper cylinders and the $\Delta$-wings lie in vertical slab regions of widths strictly greater than $\pi$.

In this paper we do not assume graphicality of the translator. This more general setting opens the possibility for a much wider variety of examples. For an overview of known examples of translators which are confined to some vertical slab, we refer to Section 2 below.

Our main result is the following rigidity theorem for the grim reaper cylinder among low-width translators confined in possibly slanted upper-halfspaces.

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Theorem 1.2. Let $\Sigma^2 \subseteq \mathbb{R}^3$ be a translator with finite entropy contained in a vertical slab of width $\lambda \pi$ for some $\lambda \in [1, \sqrt{2}]$. If $\Sigma \subseteq \{ \left(1 - \frac{\lambda^2}{2}\right)x_3 + \alpha x_2 \geq c \}$ for some $\alpha^2 \geq \lambda^2 - 1$ and $c \in \mathbb{R}$, then necessarily $\lambda = 1$, $\alpha = 0$, and $\Sigma$ is a grim reaper cylinder $\Gamma \times \mathbb{R}$.

We emphasize in particular the following rigidity result which is a direct consequence of Theorem 1.2.

Corollary 1.3. Let $\Sigma^2 \subseteq \mathbb{R}^3$ be a translator with finite entropy and contained in a vertical slab of width $\omega = \pi$. Suppose that $\Sigma$ is contained in an upper halfspace. Then, up to rigid motions, $\Sigma$ is the standard grim reaper cylinder $\Gamma \times \mathbb{R}$.

Remark 1.4. Concerning the width bound in Corollary 1.3, note that the so-called $\Delta$-wings (see [HIMW19]) give examples of translators of finite entropy which are bounded from below and contained in vertical slabs of width $w > \pi$. This shows the necessity/sharpness of the condition $w = \pi$ for the uniqueness in Corollary 1.3.

To the authors' knowledge, it is still an open problem whether there could exist examples of translators in a slab of width $< \pi$, which are not bounded from below. By Proposition 1.1, such examples could certainly not be graphs. However one might naively imagine connecting two planes by a catenoidal neck and adding a (perhaps large) perturbation; see also [HMW23, Remark 1.7].

As another consequence of Theorem 1.2 we obtain in particular the following non-existence result.

Corollary 1.5. There do not exist any 2-dimensional translators $\Sigma^2 \subseteq \mathbb{R}^3$ with finite entropy, contained in a slab of width $\omega$ and satisfying one of the following assumptions:

(a) $\omega < \pi$, and $\Sigma$ is contained in an upper halfspace w.r.t. the translating direction;
(b) $\omega = \pi$, and $\Sigma$ is contained in a slanted halfspace of the form $\{x_3 + \beta x_2 \geq c\}$, with $\beta \neq 0$.

Note that a similar result as the first part of the previous Corollary 1.5, was obtained in [GMM22, Proposition 9.1]. In this latter paper the same conclusion was reached imposing, additionally to the finite entropy requirement, topological assumptions (finite genus and only one end) instead of the confinement to an upper halfspace, and with different techniques. These topological properties are features of stable translators. See [KS19, IR19] and Section 3.3

We finally note that the assumption of collapsedness (the translators being contained in vertical slabs) is a very natural one, for example in the light of the classification of the projected convex hulls of translators in [CM19], namely as: slabs (incl. their degenerate cases of flat planes, halfspaces or the entire space).

Appendix A, of independent interest, contains the characterization of equidistant surfaces to a translator with respect to the Ilmanen conformal metric. This yields an interesting foliation of the vertical $\pi$-slab.

Added note

During the writing of the manuscript, the authors were made aware that some results in the same direction as Theorem 1.2 were also obtained in the very recent preprint by D. Hoffman, F. Martín and B. White, [HMW23]. The two proofs are however clearly different in spirit and the results as well have different features. The most evident advantage of the approach in [HMW23], using suitable nonlinear barriers, is that they only assume that the slab-confined translator is lower bounded on one $x_2$-slice, while our upper halfspace assumption is asking that there are no wings going down.

On the other hand, our direct approach permits to easily get a sharp rigidity statement also in the boundary-less setting, which seems to be more difficult to be handled with the techniques in [HMW23], where the translators with boundary in their Theorem 1.3 is assumed
to obey the Dirichlet-type condition that the boundary consists of straight lines, while in their Theorem 1.4 there is no rigidity or uniqueness statement, which we do obtain in our Theorem 1.2. Furthermore we explicitly note that, when considering slanted halfspaces, our Theorem 1.2 even gives non-existence also in slightly wider slabs, up to the width $\sqrt{2\pi}$.

2. Examples

We list here all the examples of complete translators, i.e. solutions to (1), in slabs of $\mathbb{R}^3$ which we are aware of at the moment.

1. Vertical planes.
2. Grim reaper cylinders $\Gamma \times \mathbb{R}$, where $\Gamma$ is the grim reaper curve, as well as the so-called tilted grim reaper cylinders (see e.g. [SX20]).
3. The so-called $\Delta$-wings [HIMW19].
4. Nguyen’s periodic examples aka. “tridents” [Ngu09], [HMW22b].
5. Pitchforks (by Hoffman, Martín and White [HMW22a])
6. Prongs (by Hoffman, Martín and White, [HMW]).
7. Capped and uncapped translating annuloids (by Hoffman, Martín and White [HMW], [HMW23]).

Among these examples: note that the capped translating annuloids are collapsed and bounded from below. In particular, our theorem shows that the width of any example with those two properties has at least the width of the grim reaper cylinder.

Apart from these, there are of course numerous examples of translators not contained in slabs, e.g. the bowl soliton, the Clutterbuck-Schnürer-Schulze examples, [CSS07], and gluings of these, [DdPN17].

3. Some preliminary material

3.1. Background on weighted manifolds and parabolicity. Let $(M^m, g_M)$ be a complete $m$-dimensional Riemannian manifold with (possibly empty) smooth boundary $\partial M$. Given $h \in C^\infty(M)$ we can consider the weighted manifold $M^m_h = (M^m, g_M, e^{-h}d\text{vol}_M)$, where $d\text{vol}_M$ denotes the canonical Riemannian volume form on $M$. Associated to the weighted manifold $M^m_h$ there is a natural divergence form second order diffusion operator, the $h$-Laplacian, defined on $u$ by

$$\Delta_h u = e^h \text{div} \left( e^{-h} \nabla u \right) = \Delta u - g_M(\nabla u, \nabla h).$$

This is clearly symmetric on the space $L^2(M^m_h)$, endowed with the inner product

$$(u, v)_{L^2(M^m_h)} = \int_M u v e^{-h}d\text{vol}_M.$$

Let $\eta$ be the outward pointing unit normal to $\partial M$. The weighted manifold $M^m_h$ is said to be (Neumann) $h$-parabolic if for every $u \in C^0(M) \cap W^{1,2}_{\text{loc}}(\text{int}(M), e^{-h}d\text{vol}_M)$ such that

$$\begin{cases} 
\Delta_h u \geq 0 & \text{on int}(M) \\
\frac{\partial u}{\partial \eta} \leq 0 & \text{on } \partial M \\
u^* = \sup_{M} u < +\infty 
\end{cases}$$

$u$ is necessarily constant. In the boundary-less setting we just leave out the boundary condition.

As a matter of fact, $h$-parabolicity is related to a wide class of equivalent properties involving the recurrence of the Brownian motion, $h$-capacities of condensers, the heat kernel associated to the drifted Laplacian, weighted volume growth, function theoretic tests, global divergence theorems and many other geometric and potential-analytic properties. Referring to [IdLPS18] and [IPS17] for a general presentation of the theory and for detailed proofs, here we limit ourselves to pointing out a few characterizations.
Proposition 3.1 ([IPS17], [IdLPS18]). Let $M_h$ be a weighted manifold with boundary $\partial M \neq \emptyset$ and outward pointing unit normal $\eta$. Then $M_h$ is $h$-parabolic if and only if one of the following equivalent conditions holds:

(a) (Ahlfors maximum principle) For every domain $D \subseteq M$ with $\partial_0 D = \partial D \cap \text{int}(M) \neq \emptyset$ and for every $u \in C^0(D) \cap W^{1,2}_{\text{loc}}(D, e^{-h} \text{dvol}_M)$ satisfying

\[ \begin{cases} \Delta_h u \geq 0 & \text{on int}(D) \\ \frac{\partial u}{\partial \eta} \leq 0 & \text{on } \partial M \cap D \\ \sup_D u < +\infty \end{cases} \]

in the weak sense, it holds that

\[ \sup_D u = \sup_{\partial_0 D} u. \]

(b) (Global $L^2$-Stokes theorem) Let $X$ be a vector field on $M$ satisfying

\[ |X| \in L^2(M_h), \]

\[ g_M(X, \eta) \in L^1((\partial M)_h), \]

\[ \text{div} X \in L^1_{\text{loc}}(M), \quad (\text{div} X)_- \in L^1(M_h). \]

Then

\[ \int_M e^h \text{div} \left( e^{-h} X \right) e^{-h} \text{dvol}_M = \int_{\partial M} g_M(X, \eta) e^{-h} \text{dvol}_{m-1}, \]

where $\text{dvol}_{m-1}$ is the $(m - 1)$-dimensional Hausdorff measure.

Moreover, when $D = M$, if $M_h$ is $h$-parabolic and $u \in C^0(M) \cap W^{1,2}_{\text{loc}}(D, e^{-h} \text{dvol}_M)$ satisfies

\[ \begin{cases} \Delta_h u \geq 0 & \text{on int}(M) \\ \sup_M u < +\infty \end{cases} \]

then

\[ \sup_M u = \sup_{\partial M} u. \]

From a geometric point of view, $h$-parabolicity is related to the growth-rate of the weighted volumes of intrinsic metric objects. Indeed, exploiting a standard result due to Grigor’yan, [Gri99, Theorem 7.3], in [IdLPS18] is given a direct proof of the following result, which will be sufficient for our scope here.

Proposition 3.2 (see e.g. [IdLPS18]). Let $M_h$ be a complete weighted manifold with (possibly empty) boundary $\partial M$. If, for some reference point $o \in M$,

\[ \text{vol}_h(B_r^M(o)) = \mathcal{O}(r^3), \quad r \to +\infty \]

then $M_h$ is $h$-parabolic.

Here we are denoting by $\text{vol}_h(B_r^M(o))$ the weighted volume of the geodesic ball of radius $r$ centered at a reference point $o \in M$.

3.2. Useful facts about translators. In this subsection we will fix some notation and collect some general results about translators that will be used throughout the paper.

First of all, let $x : \Sigma^2 \to \mathbb{R}^3$ be a translator of the mean curvature flow i.e. satisfying (1). Hereafter, $(x_1, x_2, x_3)$ are the standard coordinates of $\mathbb{R}^3$ and $(e_1, e_2, e_3)$ is the standard orthonormal basis of $\mathbb{R}^3$. We will write the standard Euclidean metric on $\mathbb{R}^3$ as $g_{\mathbb{R}^3} = \delta_{ij}$, but we will also often use the more standard notation $\langle , \rangle$. Moreover, with a slight abuse of notation, given a translator $x : \Sigma^2 \to \mathbb{R}^3$, we will use $x_i$ also to denote the restriction on $\Sigma$ of the $i$-th coordinate function $x_i$, that is we will identify $x_i(p)$ with $\langle x(p), e_i \rangle$, for any $p \in \Sigma$. 
The translator equation (1) turns out to be the Euler-Lagrange equation for the weighted area functional
\[ \text{vol}_f(\Sigma) = \int_\Sigma e^{-f} \text{dvol}_\Sigma, \]
when choosing \( f = -x_3 \). Clearly, the translator \( \Sigma \) inherits the weighted structure of the ambient space \( \mathbb{R}^3_f \). Indeed, we can consider intrinsically the weighted manifold
\[ \Sigma^2_\tilde{f} = \left( \Sigma, x^*((\cdot), \cdot), e^{-\tilde{f}} \text{dvol}_\Sigma \right), \]
where \( \tilde{f} = f \circ x \). It is customary to drop the “tilde” in the weight function and to write \( \Sigma^2_f \).

In the following lemma we recall a basic identity holding on translators that will be useful in what follows. For a proof we refer to [IR19, Lemma 2.1].

**Lemma 3.3.** Set \( f = -x_3 \) and denote by \( \nu \) the unit normal vector to \( \Sigma \), by \( A \) the second fundamental form of the immersion and by \( e_i^T = e_i - (e_i, \nu) \nu \) the projection on \( \Sigma \) of the unit vector \( e_i \in \mathbb{R}^3 \), \( i = 1, 2, 3 \). Then
\[ \Delta_{\tilde{f}} x_i = (e_i, e_3), \quad i = 1, 2, 3. \]

Since throughout the paper we will focus our attention on translators with finite entropy, it is useful to recall the definition of entropy and some related properties.

**Definition 3.4 ([MM09], [CM12]).** Let \( \Sigma^n \subset \mathbb{R}^{n+k} \) be a submanifold. Given \( x_0 \in \mathbb{R}^{n+k} \) and \( t > 0 \), set
\[ F_{x_0, t}(\Sigma) = \frac{1}{(4\pi t)^{\frac{n+k}{2}}} \int_\Sigma e^{-\frac{d_0(x_0)^2}{4t}} \text{dvol}_\Sigma. \]
The entropy \( \lambda(\Sigma) \) of \( \Sigma \) is then defined as
\[ \lambda(\Sigma) = \sup_{x_0 \in \mathbb{R}^{n+k}, t > 0} F_{x_0, t}(\Sigma). \]

Note that not all translators have finite entropy. It is well known (see [Sun18]) that having finite entropy is equivalent to have some control on the extrinsic volume growth. More specifically, a submanifold \( \Sigma^k \subset \mathbb{R}^{n+k} \) has finite entropy if and only if
\[ \sup_{B_{R}^{n+k}(x)} \frac{\text{vol} \left( \Sigma \cap B_{R}^{n+k}(x) \right)}{R^n} < +\infty. \]

As a consequence of the latter equivalence, we are able to relate the finite entropy assumption to potential-analytic properties of 2-dimensional translators. More specifically, we have the validity of the following proposition.

**Proposition 3.5.** Let \( \Sigma^2 \subset \mathbb{R}^3 \) be a translator. If \( \Sigma \) has finite entropy then it is parabolic. Moreover, if \( g \in C^\infty(\Sigma) \) is any function satisfying \( \inf_\Sigma g > -\infty \), then \( \Sigma \) is also \( g \)-parabolic.

**Proof.** Fix a reference point \( o \in \Sigma \). Since \( \Sigma \) has finite entropy and extrinsic volumes are bigger than intrinsic volumes, it follows that
\[ \text{vol} \left( B^\Sigma_R(o) \right) \leq \text{vol} \left( \Sigma \cap B^{\mathbb{R}^3}_R(o) \right) \leq C \lambda(\Sigma) R^2. \]
The conclusion is now a direct consequence of Proposition 3.2. Finally, if \( g \in C^\infty(\Sigma) \) is any function bounded from below, then
\[ \text{vol}_g(B^\Sigma_R(o)) = \int_{B^\Sigma_R(o)} e^{-g} \text{dvol}_\Sigma \leq e^{-\inf_\Sigma g} \int_{B^\Sigma_R(o)} \text{dvol}_\Sigma = C \text{vol}(B^\Sigma_R(o)) \leq \hat{C} R^2. \]
with \( \hat{C} \) a constant depending on the entropy of \( \Sigma \). \( \square \)
3.3. Stability for translators. Translators being critical points for the weighted area functional it makes sense to investigate stability properties. Given a translator $x : \Sigma^m \to \mathbb{R}^{m+1}$ and a compactly supported normal variation $\nu$, for some $u \in C^\infty_c(\Sigma)$, the second variation formula for the weighted area functional yields the quadratic form

$$Q_f(u, u) = \int_\Sigma \left( |\nabla u|^2 - |A|^2 u^2 \right) e^{-f} \mathrm{dvol}_\Sigma.$$ 

In particular we say that the translator is $f$-stable if $Q_f(u, u)$ is non-negative for any $u \in C^\infty_c(\Sigma)$.

An interesting fact, first observed by T. Ilmanen, is that $x : \Sigma^m \to \mathbb{R}^{m+1}$ is a translator if and only if it is minimal in $(\mathbb{R}^{m+1}, \langle \cdot, \cdot \rangle)$. Therefore, letting $\bar{g} = x^* \langle \cdot, \cdot \rangle$, one could also consider stability of $x : (\Sigma^m, \bar{g}) \to (\mathbb{R}^{m+1}, \langle \cdot, \cdot \rangle)$ as a critical point of the area functional associated to this conformal metric. It turns out that these two notions of stability are equivalent; see [CMZ15, Proposition 14].

4. Rigidity in slabs

By Proposition 3.5, Theorem 1.2 follows from the following more general result

**Theorem 4.1.** Let $\Sigma^2 \subset \mathbb{R}^3$ be a translator which is parabolic and contained in a vertical slab of width $\pi \lambda$ for some $\lambda \in [1, \sqrt{2}]$. If $\Sigma \subseteq \left\{ \left(1 - \frac{x^2}{2}\right)x_3 + \alpha x_2 \geq c \right\}$ for some $\alpha^2 \geq \lambda^2 - 1$ and $c \in \mathbb{R}$, then necessarily $\lambda = 1$, $\alpha = 0$, and $\Sigma$ is a grim reaper cylinder $\Gamma \times \mathbb{R}$.

**Proof.** Let $x : \Sigma^2 \to \mathbb{R}^3$ be a complete translator with finite entropy contained in the slab $(0, \lambda \pi) \times \mathbb{R}^2$ and assume that $\inf_{\Sigma} \left( \left(1 - \frac{x^2}{2}\right)x_3 + \alpha x_2 \right) > -\infty$.

Consider the function $u : \Sigma \to \mathbb{R}$ given by

$$(2) \quad u = \frac{x_3}{\lambda^2} + \frac{\alpha}{\lambda^2} x_2 + \ln \sin \left( \frac{x_1}{\lambda} \right).$$

Since $x_1 \in (0, \lambda \pi)$ we have that $\sin(\frac{x_1}{\lambda}) > 0$. Moreover

$$\nabla u = \frac{\nabla x_3}{\lambda^2} + \frac{\alpha}{\lambda^2} \nabla x_2 + \frac{1}{\lambda \sin(x_1/\lambda)} \nabla x_1$$

and, as a consequence of Lemma 3.3,

$$\Delta f u = \frac{1}{\lambda^2} \Delta f + \Delta f \ln \sin \left( \frac{x_1}{\lambda} \right) = \frac{1}{\lambda^2} - \frac{1}{\lambda^2} \frac{1}{\sin^2(x_1/\lambda)} |\nabla x_1|^2.$$

Then

$$\Delta f e^u = e^u (\Delta f u + |\nabla u|^2)$$

$$= e^u \left( 1 - \frac{|\nabla x_1|^2}{\sin^2(x_1/\lambda)} + \frac{|\nabla x_2|^2}{\lambda^2} + \frac{\cos^2(x_1/\lambda)}{\sin^2(x_1/\lambda)} |\nabla x_1|^2 \right)$$

$$+ \frac{\alpha^2}{\lambda^2} |\nabla x_2|^2 + \frac{2 \cos(x_1/\lambda)}{\lambda \sin(x_1/\lambda)} \langle \nabla x_1, \nabla (x_3 + \alpha x_2) \rangle + \frac{2 \alpha}{\lambda^2} \langle \nabla x_3, \nabla x_2 \rangle.$$

On the other hand

$$\frac{2 \cos(x_1/\lambda)}{\lambda \sin(x_1/\lambda)} \langle \nabla x_1, \nabla (x_3 + \alpha x_2) \rangle = 2 \langle \nabla u, \nabla (x_3 + \alpha x_2) \rangle - \frac{2}{\lambda^2} |\nabla (x_3 + \alpha x_2)|^2.$$
Hence we deduce that
\[
\Delta f e^{\frac{u}{\lambda^2}} = \frac{e^u}{\lambda^2} \left( 1 - |\nabla x_1|^2 + \frac{|\nabla x_3|^2}{\lambda^2} + \frac{\alpha^2}{\lambda^2} |\nabla x_2|^2 + \frac{2\alpha^2}{\lambda^2} \langle \nabla x_3, \nabla x_2 \rangle 
\right) 
\]
\[
- \frac{2}{\lambda^2} |\nabla x_3|^2 \left( - \frac{4\alpha}{\lambda^2} \langle \nabla x_3, \nabla x_2 \rangle - \frac{2\alpha^2}{\lambda^2} |\nabla x_2|^2 \right) 
\]
\[
+ 2\langle \nabla e^u, \nabla x_3 + \alpha x_2 \rangle 
\]
\[
= \frac{e^u}{\lambda^2} \left( 1 - |\nabla x_1|^2 - \frac{|\nabla x_3|^2}{\lambda^2} - \frac{\alpha^2}{\lambda^2} |\nabla x_2|^2 - \frac{2\alpha^2}{\lambda^2} \langle \nabla x_3, \nabla x_2 \rangle 
\right) 
\]
\[
+ 2\langle \nabla e^u, \nabla x_3 + \alpha x_2 \rangle. 
\]

Set \( g = \frac{2\alpha}{\lambda^2} x_2 + \frac{2-\lambda^2}{\lambda^2} x_3 \). Then, using the fact that \( |\nabla x_i|^2 = 1 - \langle e_i, \nu \rangle^2 \), it is immediate to see that
\[
(3) \quad \Delta g e^{\frac{u}{\lambda^2}} = e^g \text{div}(e^{-g} \nabla e^{\frac{u}{\lambda^2}}) 
\]
\[
= \frac{e^u}{\lambda^4} \left( (\lambda^2 - 1 - \alpha^2) \langle e_1, \nu \rangle^2 - (\alpha \langle e_3, \nu \rangle - \langle e_2, \nu \rangle)^2 \right). 
\]

Note that by assumption \( \inf \Sigma g > -\infty \). Hence, Proposition 3.5 implies that \( \Sigma \) is \( g \)-parabolic and:

- **Case 1.** If \( \alpha^2 > \lambda^2 - 1 \) then \( \Delta g e^{\frac{u}{\lambda^2}} \leq 0 \) and \( u \) has to be constant. The constancy of \( u \) forces \( \langle e_1, \nu \rangle = 0 \) and \( \alpha \langle e_3, \nu \rangle = \langle e_2, \nu \rangle \). Since \( H = \langle e_3, \nu \rangle \), using the identity
  \[
  1 = \langle e_1, \nu \rangle^2 + \langle e_2, \nu \rangle^2 + \langle e_3, \nu \rangle^2, 
  \]
this would imply that \( H \) should be a positive constant, which is absurd by the translator equation \((1)\).

- **Case 2.** Let us assume that \( \alpha^2 = \lambda^2 - 1 \). Then, again \( \Delta g e^{\frac{u}{\lambda^2}} \leq 0 \) and \( u \) has to be constant. In this case either \( \lambda = 1 \) and we conclude that \( \Sigma \) must be a grim reaper cylinder, or \( \lambda \in (1, \sqrt{2}) \) (see the next Remark 4.3). However, in this case, the constancy of \( u \) implies that \( \Sigma \) is a tilted grim. However any tilted grim in our slab intersects the planes of the form \( \frac{2\alpha}{\lambda^2} x_2 + \frac{2-\lambda^2}{\lambda^2} x_3 = c \).

\[ \square \]

**Remark 4.2.** We remind the reader that the crucial property that \( \Delta g e^{\frac{u}{\lambda^2}} \leq 0 \), in the above proof, means that \( (2) \) is a so-called universally \( L \)-superharmonic function, i.e. superharmonic for this operator on any self-translating soliton.

**Remark 4.3.** Concerning the assumptions in Theorem 1.2, note that in case \( \lambda > \sqrt{2} \) the assumption that \( \Sigma \subseteq \left\{ \left(1 - \frac{\lambda^2}{2}\right)x_3 + \alpha x_2 \geq c \right\} \) implies that the function \( \left(\frac{\lambda^2}{2} - 1\right)x_3 - \alpha x_2 \) is bounded from above. However, since \( \Sigma \) is properly immersed and
\[
\Delta f \left( \left(\frac{\lambda^2}{2} - 1\right)x_3 - \alpha x_2 \right) = \frac{\lambda^2}{2} - 1 > 0, 
\]
one immediately reaches a contradiction as an application of the Omori-Yau maximum principle for the \( f \)-Laplacian, holding in this setting since \( |\nabla f| \) is bounded and \( \Sigma \) is properly immersed by our assumption on the entropy; see e.g. [PRS05, Example 1.14]

**Appendix A. Equidistant surfaces to translators with respect to the conformal metric**

**A.1. Geodesics of \((\mathbb{R}^3, \langle , \rangle)\).** Consider on \( \mathbb{R}^3 \) the conformal metric \( \langle , \rangle = e^{x_3} \langle , \rangle \). We recall that \( \Sigma^2 \) is a translator in \((\mathbb{R}^3, \langle , \rangle)\) if and only if it is minimal in \((\mathbb{R}^3, \langle , \rangle)\).
The Christoffel symbols of \((\mathbb{R}^3, \langle , \rangle)\) are given by
\[\Gamma^k_{ij} = \partial_i \frac{x^k}{2} \delta^i_j + \partial_j \frac{x^k}{2} \delta^i_i - \partial_k \frac{x^i}{2} \delta^i_j.\]
In particular, the non-vanishing ones are
\[
\begin{cases}
\Gamma^1_{13} = \Gamma^2_{23} = \Gamma^3_{31} = \Gamma^3_{32} = \frac{1}{2}, \\
\Gamma^3_{11} = \Gamma^3_{22} = -\frac{1}{2}.
\end{cases}
\]
A curve \(\gamma = (\gamma_1, \gamma_2, \gamma_3)\) in \(\mathbb{R}^3\) is a geodesic in \((\mathbb{R}^3, \langle , \rangle)\) if and only if its component functions satisfy the geodesic equations
\[
\begin{cases}
\dot{\gamma}_1 + \dot{\gamma}_1 \dot{\gamma}_3 = 0 \\
\dot{\gamma}_2 + \dot{\gamma}_2 \dot{\gamma}_3 = 0 \\
\dot{\gamma}_3 + \frac{1}{2} \{(\ddot{\gamma}_3)^2 - (\dot{\gamma}_1)^2 - (\dot{\gamma}_2)^2\} = 0.
\end{cases}
\]
Let \(p \in \mathbb{R}^3\) and \(v \in \mathbb{R}^3\) such that \(\langle v, v \rangle = 1\) and let us denote by \(\gamma\) the unique geodesic of \((\mathbb{R}^3, \langle , \rangle)\) satisfying \(\gamma(0) = p, \gamma(0) = v\). Then
\[
\dot{\gamma}_1^2 + \dot{\gamma}_2^2 = e^{-\gamma_3} - \dot{\gamma}_3^2.
\]
Hence the third equation in (4) reads as
\[
\dot{\gamma}_3 + \frac{1}{2} \{-e^{-\gamma_3} + 2 \dot{\gamma}_3^2\} = 0.
\]
Multiplying by \(e^{\gamma_3}\) on both sides and integrating between 0 and \(t\), gives
\[
\gamma_3(t) = \ln \left( \frac{t^2 + 4 \dot{\gamma}_3(0)e^{\gamma_3(0)}t + 4e^{\gamma_3(0)}}{4} \right).
\]
Note that this expression is defined for all \(t > 0\).

We now integrate the first and the second equation in (4). For \(i = 1, 2\) we have that
\[
\dot{\gamma}_i(t) = \dot{\gamma}_i(0)e^{-\gamma_3(t)+\gamma_3(0)} = \frac{4\dot{\gamma}_i(0)e^{\gamma_3(0)}}{t^2 + 4\dot{\gamma}_3(0)e^{\gamma_3(0)}t + 4e^{\gamma_3(0)}}
\]
\[
\gamma_i(t) = \gamma_i(0) + 4\dot{\gamma}_i(0)e^{\gamma_3(0)} \int_0^t \frac{du}{u^2 + 4\dot{\gamma}_3(0)e^{\gamma_3(0)}u + 4e^{\gamma_3(0)}}
\]
Note that
\[
\int_0^t \frac{du}{u^2 + 4\dot{\gamma}_3(0)e^{\gamma_3(0)}u + 4e^{\gamma_3(0)}} = \int_0^t \frac{du}{(u + 2\dot{\gamma}_3(0)e^{\gamma_3(0)})^2 + 4e^{\gamma_3(0)}(1 - \dot{\gamma}_3(0)^2)e^{\gamma_3(0)}} = \frac{2\sqrt{1 - \dot{\gamma}_3^2} \gamma_3 e^{\gamma_3}}{2e^{\gamma_3(0)} \sqrt{\gamma_3^2 + \gamma_2^2}} \left[ \arctan \left( \frac{t + 2\dot{\gamma}_3(0)e^{\gamma_3(0)}}{2e^{\gamma_3(0)} \sqrt{\gamma_3^2 + \gamma_2^2}} \right) - \arctan \left( \frac{\dot{\gamma}_3(0)}{\sqrt{\gamma_1^2 + \gamma_2^2}} \right) \right],
\]
where in the third identity we have used (5). Hence, recalling that
\[\arctan t_1 - \arctan t_2 = \arctan \left( \frac{t_1 - t_2}{1 + t_1t_2} \right),\]
a simple computation yields
\[
\gamma_i(t) = \gamma_i(0) + 2\frac{\dot{\gamma}_i(0)}{\sqrt{\gamma_1(0) + \gamma_2(0)}} \arctan \left( \frac{t \sqrt{\gamma_1(0) + \gamma_2(0)}}{t \gamma_3(0) + 2} \right).
\]
Summarising, the geodesic $\gamma$ can be parametrized by
\[
\begin{align*}
\gamma_1(t) &= \gamma_1(0) + 2\frac{\dot{\gamma}_1(0)}{\sqrt{\dot{\gamma}_1(0) + \dot{\gamma}_2(0)}} \arctan\left(\frac{t\sqrt{\dot{\gamma}_1(0) + \dot{\gamma}_2(0)}}{t\dot{\gamma}_3(0) + 2}\right) \\
\gamma_2(t) &= \gamma_2(0) + 2\frac{\dot{\gamma}_2(0)}{\sqrt{\dot{\gamma}_1(0) + \dot{\gamma}_2(0)}} \arctan\left(\frac{t\sqrt{\dot{\gamma}_1(0) + \dot{\gamma}_2(0)}}{t\dot{\gamma}_3(0) + 2}\right) \\
\gamma_3(t) &= \ln\left(\frac{t^2 + 4\dot{\gamma}_3(0)e^{\gamma_3(0)t} + 4e^{\gamma_3(0)}}{4}\right).
\end{align*}
\]

A.2. $(\cdot,\cdot)$-equidistant surfaces to translators. Let $\Sigma^{2}$ be a translator in $\mathbb{R}^{3}$, $p = (p_1, p_2, p_3) \in \Sigma^{2}$ and denote by $\nu = (\nu_1, \nu_2, \nu_3)$ its unit normal. Consider the geodesic $\gamma : \mathbb{R} \to (\mathbb{R}^{3}, \langle \cdot, \cdot \rangle)$ with initial conditions $\gamma(0) = p$ and $\dot{\gamma}(0) = e^{-p_3/2}\nu$. By the translator equation (1) we have that
\[
\dot{\gamma}_3(0) = He^{-p_3/2}.
\]
Therefore
\[
\dot{\gamma}_1(0)^2 + \dot{\gamma}_2(0)^2 = e^{-p_3}(1 - H^2),
\]
so that
\[
\frac{\dot{\gamma}_1(0)}{\sqrt{\dot{\gamma}_1(0)^2 + \dot{\gamma}_2(0)^2}} = \frac{\nu_1}{\sqrt{1 - H^2}}.
\]
Substituting in (4), we therefore obtain that the surface at distance $t$ from $\Sigma$ in the direction of $\dot{\gamma}(0) = e^{-p_3/2}\nu$ is parametrized by
\[
\begin{align*}
x_{1,t} &= p_1 + 2\frac{\nu_1}{\sqrt{1 - H^2}} \arctan\left(\frac{\sqrt{1 - H^2}}{H + 2e^{p_3/2}}\right) \\
x_{2,t} &= p_2 + 2\frac{\nu_2}{\sqrt{1 - H^2}} \arctan\left(\frac{\sqrt{1 - H^2}}{H + 2e^{p_3/2}}\right) \\
x_{3,t} &= \ln\left(\frac{t^2 + 4He^{p_3/2}t + 4e^{p_3}}{4}\right).
\end{align*}
\]
We assume now that $H \equiv 0$, that is, wlog,
\[
\Sigma = \Sigma_0 = \{(0, x_2, x_3) \in \mathbb{R}^{3} : x_2, x_3 \in \mathbb{R}\}.
\]
By (6), for any $t \in \mathbb{R}$ the equidistant surface $\Sigma_t \subseteq \mathbb{R}^{3}$ at (signed) distance $t$ from $\Sigma_0$ is parametrized by
\[
\begin{align*}
x_{1,t} &= 2 \arctan\left(\frac{t - x_3/2}{2}\right) \\
x_{2,t} &= x_2 \\
x_{3,t} &= \ln\left(\frac{t^2 + 4e^{x_3}}{4}\right).
\end{align*}
\]
Let us study some features of the family of equidistant surfaces $\{\Sigma_t\}_{t \in \mathbb{R}}$:

(i) $\{\Sigma_t\}_{t \in \mathbb{R}}$ foliate $(-\pi, \pi) \times \mathbb{R} \times \mathbb{R}$.

(ii) Let’s focus here on the case $t > 0$. Setting $x_1 = 2 \arctan\left(\frac{t - x_3/2}{2}\right)$ we can reparametrize $\Sigma_t$ as follows:
\[
\left\{(x_1, x_2, \ln\left(\frac{t^2}{4}\right) - 2 \ln\sin\left(\frac{x_1}{2}\right)) : (x_1, x_2) \in (0, \pi) \times \mathbb{R}\right\}.
\]
Therefore the unit normal to $\Sigma_t$ is
\[
\nu_t = \left(\cos\left(\frac{x_1}{2}\right), 0, \sin\left(\frac{x_1}{2}\right)\right)
\]
and the first and the second fundamental form of $\Sigma_t$ are given by
\[
g_{\Sigma_t} = \begin{pmatrix}
\frac{1}{\sin^2(x_1/2)} & 0 \\
0 & 1
\end{pmatrix}, \quad H_{\Sigma_t} = \begin{pmatrix}
\sin^2(x_1/2) & 0 \\
0 & 1
\end{pmatrix} \begin{pmatrix}
\frac{1}{2 \sin(x/2)} & 0 \\
0 & 0
\end{pmatrix}.
\]
In particular,

\[ H_{\Sigma_t} = \frac{1}{2} \sin(x_1/2) = \frac{1}{2} \langle e_3, \nu_{\Sigma_t} \rangle. \]

Note that any \( \Sigma_t \) looks like an open piece of a rescaled (with a factor 2) grim reaper cylinder. In particular \( \Sigma_t \) is not a translator.

References


[HMW] David Hoffman, Francisco Martín, and Brian White. Translating annuli for mean curvature flow. *In preparation*.


