The bursty star formation history of the Fornax dwarf spheroidal galaxy revealed with the HST

Rusakov, V.; Monelli, M.; Gallart, C.; Fritz, T. K.; Ruiz-Lara, T.; Bernard, E. J.; Cassisi, S.

Published in:
Monthly Notices of the Royal Astronomical Society

DOI:
10.1093/mnras/stab006

Publication date:
2021

Document version
Publisher's PDF, also known as Version of record

Document license:
CC BY-ND

Citation for published version (APA):
The bursty star formation history of the Fornax dwarf spheroidal galaxy revealed with the HST

V. Rusakov1,2,3* M. Monelli4,5 C. Gallart4,5 T. K. Fritz4,5 T. Ruiz-Lara4,5,6 E. J. Bernard7 and S. Cassisi8,9

1Department of Physics, University of Surrey, Guildford GU2 7XH, UK
2Cosmic Dawn Center (DAWN), DK-2200 Copenhagen N, Denmark
3Niels Bohr Institute, University of Copenhagen, Lyngbyvej 2, DK-2100 Copenhagen Ø, Denmark
4Instituto de Astrofisica de Canarias, E-38200 La Laguna, Tenerife, Spain
5Departamento de Astronomia, Universidad de La Laguna, E-38205 La Laguna, Tenerife, Spain
6Kapteyn Astronomical Institute, University of Groningen, Landeven 12, NL-9747 AD Groningen, the Netherlands
7Universite Côte d’Azur; Observatoire de la Côte d’Azur, CNRS, Laboratoire Lagrange, F-06304 Nice, France
8INAF – Astronomical Observatory of Abruzzo, Via M. Maggini, 1-64100 Teramo, Italy
9INFN, Sezione di Pisa, Largo Pontecorvo 3, I-56127 Pisa, Italy

Accepted 2020 December 24. Received 2020 December 7; in original form 2020 February 22

ABSTRACT
We present a new derivation of the star formation history (SFH) of the dSph galaxy Fornax in two central regions, characterized by unprecedented precision and age resolution. It reveals that star formation has proceeded in sharp bursts separated by periods of low level or quiescent activity. The SFH was derived through colour–magnitude diagram (CMD) fitting of two extremely deep Hubble Space Telescope CMDs, sampling the centre and one core radius. The attained age resolution allowed us to single out a major star formation episode at early times, a second strong burst 4.6 ± 0.4 Gyr ago and recent intermittent episodes ∼2–0.2 Gyr ago. Detailed testing with mock stellar populations was used to estimate the duration of the main bursts and study the occurrence of low-level star formation between them. The SFHs in both regions show common features, with activity at the same epochs and similar age–metallicity relationship. However, clear indications of a spatial gradient were also found, with mean age increasing with radius and star formation episodes being more prolonged in the centre. While some galaxy evolution models predict bursty SFHs in dwarf galaxies and thus a secular origin of the observed SFH cannot be excluded in Fornax, other evidence points to possible mergers or interactions as the cause of its bursty SFH. In particular, we calculated the Fornax orbit relative to the closest dwarfs and the Milky Way and observed a correspondence between the main intermediate-age and young events and peri-passages of Fornax around the Milky Way, possibly indicating tidally induced star formation.

Key words: galaxies: dwarf – galaxies: evolution – Local Group – galaxies: star formation – galaxies: stellar content.

1 INTRODUCTION
Dwarf galaxies are the most abundant type of galaxies in the Local Group (LG). Their proximity allows for detailed studies with the ability to resolve individual stars. This makes it possible to obtain quantified, resolved star formation histories (SFHs) of these galaxies, by means of recreating their observed colour–magnitude diagrams (CMDs) with models, which provides valuable information about the stellar populations present in the galaxies and when they were formed. This information is important to understand how their evolution is affected by both internal and external factors, such as supernova feedback, host dark matter structure properties, UV-reionization, tidal effects of a host galaxy, and interaction with other dwarfs (e.g. LCID project in Gallart 2012, and references therein). Different methods and techniques employed to quantify SFHs of stellar populations are described in Tolstoy & Saha (1996), Gallart et al. (1999), Hernandez, Valls-Gabaud & Gilmore (1999), Dolphin (2002), Aparicio & Gallart (2004), Aparicio & Hidalgo (2009), and Cignoni & Tosi (2010).

In this work, we investigate the SFH of the Fornax dwarf spheroidal (dSph) galaxy. Typical dSph galaxies, also referred to as early type, are characterized by low gas content in terms of neutral and ionized hydrogen and, consequently, dominated by old stars (e.g. Cetus and Tucana dSphs in Monelli et al. 2010a,b), as opposed to late-type dwarf irregular (dIrr) systems that preserve gas for longer and exhibit younger stellar populations (e.g. IC 1613 in Skillman et al. 2014). This classification is based only on the current properties of dwarfs and complemented by a transition type (dT) with mixed properties of the former two types (e.g. LGS 3 in Hidalgo et al. 2011). Yet Kormendy (1985) suggested that most early-type galaxies are not originally different from the late-type, but that their evolution diverged after their gas was stripped by external processes. This was supported by the evidence of complicated and lasting SFHs of some dSphs (e.g. Carina dSph in Hurley-Keller, Mateo & Nemec 1998; Leo I dSph in Gallart et al. 1999 and Ruiz-Lara et al. 2020a), which

* E-mail: rusakov124@gmail.com

© 2021 The Author(s)
Published by Oxford University Press on behalf of Royal Astronomical Society
retained their gas through the intermediate ages. Fornax, too, has demonstrated signs of an extended SFH, with stars formed at old- (>10 Gyr ago), intermediate- (1–10 Gyr ago) and, possibly, young-age (<1 Gyr ago) epochs (del Pino et al. 2013). Based on their full SFHs, Gallart et al. (2015) proposed a classification of dwarf galaxies into ‘fast’ and ‘slow’ types, where the first form most of their stars in the first few Gyr of evolution, while the latter type experience more extended star formation. Interestingly, these two types do not correspond directly with dSph and dIrr galaxies, a classification that is based in the current morphology. In particular, at least three dSph satellites of the Milky Way (MW), Leo I, Carina, and Fornax, are classified as slow dwarfs. The inferred past locations with respect to the MW or M31 of fast and slow dwarfs in the LG led Gallart et al. (2015) to hypothesize that the conditions of the environment in which these dwarfs formed could be an important factor imprinting their subsequent evolution.

Fornax was discovered by Shapley (1938) when studying photographic plates of the constellation of the same name. One of the peculiar findings about the galaxy is that it has five globular clusters (GCs); discovered by Baade & Hubble 1939, Shapley 1939, Hodge 1961a; a discovery of the sixth star cluster was reported in Wang et al. 2019a), which makes Fornax one of the few LG galaxies hosting GCs. It is also one of the most luminous LG dwarfs with $M_V = -13.4 \pm 0.3$ (Irwin & Hatzidimitriou 1995), and a stellar mass of $M_\star = 20 \times 10^6 M_\odot$ (assuming mass-to-light $M/L = 1$, McConnachie 2012).

The recent estimate of the distance modulus of Fornax $\mu = 20.818 \pm 0.015$ (statistical) $\pm 0.116$ (systematic) ($\sim 146$ kpc) by Karczmarek et al. (2017) was obtained by using mean near-infrared magnitudes of RR Lyrae (RRL) stars, which had the advantage of being less affected by extinction over studies performed with the photometric data, reaching $I \sim 24.5$. The authors claimed the resolution of star formation rate (SFR) from $\sim 1.7$ Gyr at the oldest ages to $\sim 0.6$ at an age of 3 Gyr ago.

Additional interest in the Fornax stellar populations was stirred by studies that combined spatial, chemical and dynamical information into their analyses, and discovered complex patterns across the galaxy. For example, del Pino et al. (2017) demonstrated these patterns in terms of angular momenta and line-of-sight velocities; Amorisco & Evans (2012) and Hendricks et al. (2014) used different approaches to fit stellar populations either based on their metallicity or the whole set of spatial and chemodynamical features. These works were built on the spectroscopic measurements analysed and published by Pont et al. (2004), Battaglia et al. (2006), Battaglia & Starkenburg (2012), Walker, Mateo & Olszewski (2009), Kirby et al. (2010), and Letarte et al. (2010).

The goal of this study is to investigate two small regions in the centre of Fornax, with little crowding, to obtain their SFH with high accuracy and resolution in age, exceeding that of the previous studies. It is done by using high-precision photometry reaching the oldest main sequence turnover (MSTO), going substantially deeper than the previous studies, down to $F814W \sim 27$ with the space-based HST observations.

In Section 2, we present the observations, describe the photometry procedure and artificial stars tests and discuss the derived Fornax CMDs. Then, the SFH procedure is briefly explained in Section 3. The SFH results, tests with mock stellar populations and comparison of the results with previous works are shown in Section 4. The main findings in the paper are discussed in Section 5, which also demonstrates the results of numerical orbit integration for Fornax and speculates about connection of the orbit with the SFH. The work concludes by the summary in Section 6.
2 OBSERVATIONS

2.1 Data overview

The observations of the two regions in Fornax at the positions \((\alpha_{2000}, \delta_{2000}) = (2^h39^m42^s, -34^\circ 31'43'\)\) and \((2^h40^m39^s, -34^\circ 32'41')\), hereafter called arbitrarily Fornax1 and Fornax2, were performed with the ACS/WFC camera of the Hubble Space Telescope (HST) as part of the proposal 13435\(^1\) (PI: M. Monelli) in 2013/14. The two ACS pointings were collected as parallel observations of the WFC3, centred on the GC 4 of Fornax. Given that half of the orbits were taken with a 180\(^\circ\) rotation, this resulted in two ACS fields separated by \(\sim 12\) arcmin. The WFC has a field of view of \((202 \times 202)\) arcsec\(^2\) and consists of two CCDs stacked side by side. The camera took nine dithered images of each region in F475W and F814W filters, three for each of the 80-, 520-, and 850-s exposures. The filters were chosen to cover a large bandwidth, as advised in Stetson (1993). Table A1 summarizes the major image characteristics. The footprints of the two camera pointings are overplotted on top of a DSS image of the core of Fornax\(^2\) on Fig. 1, together with the footprints from the previous study of SFH in Fornax by del Pino et al. (2013). Fornax1 is a near-central region, located at a radial distance of \(r = 2.8\) arcmin to the south-west of the galaxy centre, while Fornax2 is close to the core of Fornax, at \(r = 12.7\) arcmin to the east.

2.2 Photometric procedure

The photometry was obtained by using DAOPHOT and ALLFRAME software (Stetson 1987; Stetson 1994). Description of the general procedure and its routines, as well as comparison of DAOPHOT against other software can be found in Monelli et al. (2010a). Below is the summary of the procedure.

The two chips of the camera were processed separately, as the common field of view was discontinuous due to the gap of 2.5 arcsec between the chips. The FLC images (bias-corrected, dark-subtracted, flat-fielded and corrected for the charge transfer efficiency) were used here. They provided correct astrometric positions of the light sources, but their distorted pixels (due to the offset alignment of the ACS instrument within the telescope) resulted in distorted photometry. It was corrected by applying a pixel-area mask.

First, the Gaussian centroids of all sources were identified up to \(3\sigma\) level in all images and photometry was performed within 3-pixel apertures, all done with routines in DAOPHOT, which rejected the foreground stars that appeared as sources with saturated peaks. Then, the estimates of the instrumental magnitudes were improved by fitting a Lorentz point spread function (PSF) and subtracting sky values in ALLSTAR, after which coordinate transformations were applied to the images by filter using DAOMASTER. Then the individual images were median-combined to remove cosmic rays and any other spurious sources (such as satellites or bad pixels).

This clean image was used to search for stellar sources, through two iterations of the process from finding the sources with FIND and evaluating their magnitudes with ALLSTAR. The derived list of stars was adopted as input catalogue for ALLFRAME, which generated an individual catalogue of stars for each image. Then, the process was repeated re-modelling the PSFs using the ALLFRAME catalogues. We adopted the Moffat function with \(\beta = 1.5\) and \(\alpha\) linearly varying across the space in the images, and based on an average sample of 1000 brightest stars (in F814W) for Fornax1 and 620 stars for Fornax2, selected from the stars with refined positions and magnitudes from the previous step. A new stacked median image was generated using updated geometric transformations and the second and final ALLFRAME was run to obtain the final catalogues of stars. Different extended-shape sources (with [sharpness parameter] \(\geq 0.1\), see Stetson 1987), such as background galaxies, were removed from the catalogues.

Several corrections were performed to calibrate the obtained catalogues. The magnitudes were shifted according to the aperture correction to the standard ACS 0.5 arcsec aperture. The correction was calculated as the median difference between the estimated magnitudes of the PSF stars and magnitudes of the same stars obtained by using the curvature growth method with DAOGROW (Stetson 1990). The stellar catalogues were calibrated to the VEGAMAG photometric system with the zero-points \(ZP_{F475W} = 26.150\) and \(ZP_{F814W} = 25.518\) for each filter, taken from Bohlin (2016). Additionally, the absolute values of magnitude corrections from Sirianni et al. (2005) \(\Delta M_{F475W,\text{abs}} = 0.087\) and \(\Delta M_{F814W,\text{abs}} = 0.079\) were subtracted from the estimated magnitudes to account for stellar flux at the infinite aperture. The CMDs resulting from the obtained photometry are discussed in Section 2.4.


\(^2\)The background image was taken from the ESO Online Digitized Sky Survey (DSS) at https://archive.eso.org/dss/dss.

Figure 1. DSS image of the core region of Fornax dSph with camera pointings overlapped. Fornax1 and Fornax2 are rectangular regions in red solid lines. IC1, IC2, OC regions in red-dashed lines are from del Pino et al. (2013). Two blue circles of radius 3 arcmin show the locations of two overdensities, discovered by de Boer et al. (2013) and Coleman et al. (2004), respectively, were adopted from Bate et al. (2015) (two other, to the west and to the south-south-west, discovered by the latter study did not have coordinates specified, but are not expected to overlap with Fornax1 and Fornax2). The dashed ellipse signifies the core radius. The central position of Fornax is marked by the cross. Arrows at the bottom right indicate the directions to the north and east.
2.3 Artificial star test

Photometric errors in the images were estimated by means of the artificial star tests (ASTs), in which artificial stars were used to probe different levels of stellar crowding, as prescribed in Gallart, Aparicio & Vilchez (1996). This allowed evaluating the effect of crowding and other various possible errors. This step was essential to simulate the same photometric errors and completeness in the synthetic CMD, as in the observed one (see Section 3.1), which was essential for an accurate construction of the SFH.

To perform the tests a synthetic CMD with $5 \times 10^6$ stars and with stellar evolutionary models and assumptions outlined in Section 3.1 was created. The synthetic stars were injected into the images at the nodes of a regular grid by using the previously obtained PSF, with the distance between the nodes of $2R + 0.5$ pixels, such that the artificial stars interfered only with the light profiles of the real stars and not with each other. We used the value $R = 10$ pixels (for comparison, average PSF aperture radii were equal to $R_{\text{PSF}} = 6.6$ for stars observed in Fornax1 and $R_{\text{PSF}} = 6.2$ in Fornax2).

Additionally, we implemented a random change of (0–0.3 per cent) in magnitudes of artificial stars to account for the effect of charge transfer inefficiency of the ACS camera mentioned in Brown et al. (2014). Given the number of the artificial stars, distance between nodes in the grid and image dimensions, a number of iterations were required to complete the photometric measurement of all artificial stars.

Stars were recovered from the images by employing the same photometric procedure as for the observed photometry (Section 2.2), and the output artificial stars were separated out. The only difference in the photometric procedure this time was the PSF. It was defined in the same way as before, but based on fewer stars, so as to degrade it to a lower quality. It was done to avoid using the same PSF for injecting and recovering the stars, which could underestimate the errors for the artificial stars, compared with the observed. The degraded PSF was defined by using a random fraction of stars of the original PSF star list in the range of 80–95 per cent.

Most artificial stars were recovered with magnitudes differing from the input magnitudes, with the difference being due to the crowding effects and photometric errors, while some stars were lost, i.e. not recovered or recovered with a magnitude beyond a threshold value. This information was used for simulation of observational effects in the synthetic CMD later.

In conclusion, the fainter stars were recovered with the greatest magnitude difference, i.e. they were prone to larger errors. The photometric errors did not significantly affect the photometry even in the most crowded regions of Fornax1 and Fornax2, as the fraction $\sim 99.6$ per cent of stars experienced a magnitude change of $|\Delta m_{\text{in-out}}| \leq 0.5$, while for $\sim 90$ per cent of stars the maximum magnitude difference was only $|\Delta m_{\text{in-out}}| \leq 0.1$ in both magnitudes.

2.4 Observed CMD

Fig. 2 shows the observed CMDs in the planes of the calibrated apparent (left axis) and absolute (right axis) magnitudes. The CMDs of Fornax1 and Fornax2 are similar and span $\sim 10$ mag in $m_{\text{F814W}}$. They reach $\Delta m_{\text{F814W}} \sim 1.5$ and $\Delta m_{\text{F814W}} \sim 1.7$ below the oldest MSTO, to the level of 90 per cent completeness, respectively, which makes this photometry set the deepest among the recent studies of various regions in Fornax (de Boer et al. 2012; del Pino et al. 2013).

In addition, the high spatial resolution of the ACS images and the deep photometry in both fields produce very accurate CMDs that reveal various features distinctly.

![Figure 2. Observed CMDs of Fornax1 (left-hand panel) and Fornax2 (right-hand panel) in the planes of apparent (left axis) and absolute (right axis) magnitudes. Red lines indicate the levels of completeness estimated in the ASTs (top line: 90 per cent level; bottom line: 50 per cent level). The photometric errors in colour and magnitude, defined by the typical sigma values of magnitudes at all colours $m_{\text{F475W}} - m_{\text{F814W}}$ and corresponding magnitudes $m_{\text{F814W}}$, are indicated on the right side of the two panels. The grey contours along the MS show the 20, 50, 70, 80 per cent confidence regions of stellar density, respectively.](https://academic.oup.com/mnras/article/502/1/642/6067380)
using a true distance modulus \( \mu = 20.818 \pm 0.015 \) (statistical) \( \pm 0.116 \) (systematic) and extinction \( E(B-V) = 0.022 \) transformed to \( A_{F475W} = 0.079, A_{F814W} = 0.037 \) for Fornax1, with the transformations from Bedin et al. (2005) \( E(B-V) = 0.02 \) has been used for Fornax2 in the subsequent analysis.

For this comparison, we selected some solar-scaled isochrones for appropriate values of age and metallicity from the BaSTI\(^3\) stellar evolution library (Pietrinferni et al. 2004).

The left-hand panel of the figure shows isochrones for intermediate and young ages superimposed on the younger and more metal-rich part of the MS, which is represented by a relatively small number of stars. The blue limit can be approximated by the 0.2 and 0.4 Gyr old isochrones and some other populations can be traced with the 1.0 and 2.5 Gyr old isochrones.

The right-hand panel of Fig. 3 shows some intermediate-age and old isochrones superimposed on the CMD of Fornax1. The isochrone at \( t = 13.5 \) Gyr and \([\text{Fe/H}] = -1.79\) defines the oldest and most metal poor limit of simple stellar populations by edging the oldest MSTO and the blue side of the RGB. By extending to the red edge of the RGB, outlined by the 7-Gyr-old isochrone, this CMD region indicates a period of active star formation from \( \sim 13.5 \) to 7.0 Gyr ago. The isochrones for \( t = 5.5 \) Gyr, \( t = 4.0 \) Gyr, and \([\text{Fe/H}] = -0.96\), which outline the brighter SGB and a corresponding part of the RGB, evidence the star formation activity at intermediate ages. It can be presumed from the plots that the region Fornax1 was likely forming stars at different epochs throughout its lifetime from 13.5 Gyr ago until as recently as a few hundred Myr ago.

Besides indicating the presence of stellar populations of different ages and metallicities, the good fit of the isochrones to the tip of the RGB and the RC confirms the accuracy of the magnitude shift adopted to account for the distance modulus and foreground extinction.

### 3 SFH PROCEDURE

The preliminary analysis via comparison with isochrones presented in Section 2.4 suggests the diversity of stellar populations that comprise the studied regions. In this section, we describe the quantitative analysis performed through the computation of precise SFHs with the CMD fitting technique. Our procedure followed closely the one described in detail in Aparicio & Hidalgo (2009) and Hidalgo et al. (2011). Below is a brief overview of all the steps taken to obtain an SFH.

#### 3.1 Synthetic CMD

The first step of constructing a model CMD was to create a global synthetic population. The synthetic CMD presented here consists of \(50 \times 10^5\) stars over a sufficiently large range of ages and metallicities, based on the observed ranges. The age and metallicity were limited to the intervals \( t = 0.03 - 13.9 \) Gyr and \(-2.27 \leq [\text{Fe/H}] \leq -0.35\) (according to Fig. 3).

The theoretical evolutionary framework needed to obtain the synthetic CDMs was based on the solar-scaled BaSTI stellar evolutionary library (Pietrinferni et al. 2004), previously adopted for construction of a synthetic CMD for the ASTs and the preliminary comparison between the observed CDMs and theoretical isochrones. The synthetic CMD was created following a constant (uniform) SFR in the age–metallicity plane defined within the quoted ranges. The IMF from Kroupa (2002) was assumed with the fraction of the binary stars and their relative mass distribution \( f = 0.5 \) and \( q = 0.4 \), respectively. The latter were found to be the best combination from the studied ranges of these parameters: \( f = [0.3, 0.5, 0.7] \) and \( q = [0.1, 0.4] \) and agreed with an estimate of these values in Monelli et al. (2010a, appendix therein). Mass-loss along the RGB stage was set to \( n = 0.2 \) (Pietrinferni et al. 2004).

Following the creation of the synthetic CMD, observational effects (see Section 2.2) were simulated in it using the code DisPar,\(^4\) which applied the results of the AST for artificial stars (description of an alternative code called obersin is given in Hidalgo et al. 2011). It is important to account for these effects, as they scatter observed stars from their actual positions on the CMD, which is especially significant at faint magnitudes below the oldest MSTO (Gallart et al. 1999). The code investigated the colour–magnitude boxes of size (col, mag) = (0.04, 0.1) around a position of each synthetic star with the magnitudes \((V_s, I_s)\) and the same box in the space of the artificial stars with their input and recovered magnitudes \((V, I)\) and \((V'_s, I'_s)\), respectively. If the box contained at least 10 recovered artificial stars, one was selected randomly and its magnitudes were used to disperse the synthetic star as \(V'_s = V_s - (V - V_s)\) and \(I'_s = I_s - (I - I_s)\). If the artificial star was marked as unrecovered in the AST table, then the synthetic star was discarded. Otherwise, if a box had fewer than 10 AST stars, then a box of double size was used. A dispersed synthetic CMD is shown in Fig. 4.

---

\(^3\)The BaSTI library is available at the following web address: http://basti.oa-teramo.inaf.it

\(^4\)Description of the code DisPar is available at: https://riull.ull.es/xmlui/handle/915/3961
where the aim was for every non-zero bin to contain a number of intermediately dense RGB and RC. Further, the bundles were binned, below the MSTO, young MS that has significantly fewer stars and metallicity bins 

\[ \Delta t = 0.00031, 0.00038, 0.00046, 0.00056, 0.00067, 0.00081, 0.00098, 0.00119, 0.00144, 0.00174, 0.00211, 0.00255, 0.00309, 0.00373, 0.00452, 0.00547, 0.00661, 0.008\].

In order to compare the distribution of stars in the observed and model CMDs, they were divided into ‘bundles’, which outline CMD regions with different densities of stars. For example, separate bundles are usually selected to outline the heavily populated MS below the MSTO, young MS that has significantly fewer stars and intermediate dense RGB and RC. Further, the bundles were binned, where the aim was for every non-zero bin to contain a number of stars of the order \(~1\)–\(100\), with at least 50 per cent of bins containing 5 stars to ensure the balance between the SFH solutions and their signal-to-noise ratio. The adopted bundle design is illustrated in Fig. 4, where the sizes of the bins in each bundle are indicated. It is worth noting that the RC was not included in the binning due to the disagreements between the RC observations and models described in detail in Section 3.4, which also describes the implications of choosing the bundle and bin configurations differently.

The number of stars \(M_i^j\) distributed within every colour–magnitude box \(j\) of the synthetic CMD was calculated along synthetic SSPs \(i\) (with \(i\) from 1 to \(n \times m\), where \(n\) and \(m\) are numbers of bins in age \(t_i\) and metallicity \(z_i\)) and compared against the distribution of the observed stars \(O^j\) in the corresponding box \(j\) of the observed CMD. The procedure was aimed at reproducing the observed distribution \(O^j\) with the model stars \(M^j\) calculated by means of the linear relationship: 

\[ M^j = \sum a_i M_i^j, \quad \text{with } a_i > 0. \]

The best fit of the total distribution of model stars \(M\) against \(O\) for all \(n \times m\) models was found by minimizing the Cash statistic (Cash 1979), treated as a goodness-of-fit measure:

\[ \chi^2 = 2 \left( \sum_j (M^j - O^j \ln M^j) \right). \quad (1) \]

It is analogous to the Gaussian ‘chi-square’ quantity but derived from the Poisson log-likelihood to more accurately quantify the fit of the CMD bins with small numbers of stars. The reduced statistic \(\chi^2/v\) was used as a characteristic weight of the fit, where \(v = b = n \times m\) is the number of degrees of freedom, \(b\) is the number of colour–magnitude boxes defined in the CMD. The calculation of the optimal fit of the final model SFH and model CMD were performed using THESTORM code (see Bernard et al. 2018, which follows similar procedures as IAC-pop/MinnIAC used, for example, in Aparicio & Hidalgo 2009, Monelli et al. 2010a,b). The details of the algorithm procedure are provided in the following section.

### 3.3 Solving for a model CMD

A solution for the SFH can be generally described as an array of the scaled alpha coefficients \(\psi(t, z) = A \alpha(t, z)\), derived linearly, by estimating coefficients \(a_i\) for each SSP \(i\), and scaled by the constant \(A\), which is related to the stellar mass \(M_{\text{max}}\) of a model CMD; the area of the field of view \(S\) and age and metallicity range \(\Delta t\) and \(\Delta Z\), respectively, as 

\[ A = M_{\text{mod}}(\Delta t \times \Delta Z \times S). \]

This array represents a 2D distribution \(\alpha(t, z)\) with dimensions \([k, l]\) of age and metallicity, i.e. age–metallicity relation, which can be marginalized over all \(k\) to obtain the SFR \(\psi(t)\) at times 

\[ t_k: \psi_k(t) = A \sum \alpha_{kl}(t, z). \]

The metallicity enrichment \(z(t)\) can be estimated as a median of all products of \(z_i\) multiplied by normalized alpha coefficients \(\alpha_{kl}(t, z)\) at all times 

\[ z_k(t) = A \frac{\sum_k \alpha_{kl}(t, z) \Delta t_k}{\sum_k \alpha_{kl}(t, z) \Delta z_k}. \]

Distance modulus \(\mu\) and colour excess \(E(B - V)\) (Section 2.4) were used to convert the observed CMD to absolute magnitudes. A grid around the adopted \(\mu\) and \(E(B - V)\) values was sampled by the code to account for the uncertainties in distance modulus, reddening and photometric zero-points using the following colour and magnitude offsets: 

\[ \Delta M = \Delta \alpha_{F475W} = [-0.06, -0.03, 0.0, 0.03, 0.06] \]

and 

\[ \Delta M_{F814W} = [-0.15, -0.075, 0.0, 0.15, 0.225], \]

which resulted in 30 positions (marked by ‘+’ markers in Fig. 5).

After the position of the minimum \(\chi^2\) was identified, the code refined the search with smaller colour–magnitude shifts (to the left/right, top/bottom, totalling four additional shifts) around the position (shown as squares for Fornax1 and circles for Fornax2). The finer sampling was repeated around every newly identified minimum until convergence. At every offset, the calculation was repeated 5 times, each time applying small random shifts to the initial guess of the SFH solution, to avoid local minima, and the median of the 5 solutions was used to characterize each offset. Finally, this strategy resulted in 37 sampled offsets (37 \(\times 5 = 185\) SFHs) for Fornax1 and 42 (42 \(\times 5 = 210\) SFHs) for Fornax2, with the most likely positions marked by the filled square and circle for Fornax1 and Fornax2. The optimal colour and magnitude offsets were applied to the observed CMDs. As seen from the figure, both Fornax1 and Fornax2 CMDs...
shown as blue and red contours in Fig. 5 were used to propagate solution. The six resultant solutions was propagated to the error in the best solution with its optimal parameters fixed. The standard deviation step was expected to account for the uncertainties in the human-pixels). are smoothed with a Gaussian kernel (with three-pixel window with Poisson-resampled observed stars (see text for description). The likelihoods errors of the solutions with shifted bins at the best offset and solutions with the Poisson-resampled observed stars (see text for description). The likelihoods are smoothed with a Gaussian kernel (with three-pixel window with $\sigma = 3$ pixels).

were systematically forced towards fainter magnitudes for the best fit with the model CMDs, within the uncertainty in the distance modulus. Once the position of the optimal offset was determined we estimated the uncertainty in the resulting solution using the following procedure. Given the parametrization, no error was associated with each SFH calculation, as the best-fitting model CMD was computed by using a numerical iterative optimizer, which converged to the solution without sampling the whole involved parameter space. Instead, three main error sources affecting the best solution were considered: the uncertainty in the selection of the colour, magnitude, age, and metallicity bins, the Poisson uncertainty in the binning of the observed CMD and the uncertainty in the distance and reddening offset (as well as photometric zero points). This procedure followed the one described in Hidalgo et al. (2011).

Regarding the first error source, the code introduced changes to the colour, magnitude, age, and metallicity bin configurations: the bins were shifted by different fractions of their width, with the boundaries of their total ranges fixed. Overall, 12 different bin configurations were investigated, and an average of the 12 solutions weighted by their corresponding $\chi^2_0$ statistics was taken as the best model. This step was expected to account for the uncertainties in the human-selected parametrization of the CMDs.

Then, Poisson resampling of the colour–magnitude bin distribution of the observed CMD was performed six times to recalculate the solution with its optimal parameters fixed. The standard deviation of the six resultant solutions was propagated to the error in the best solution.

Finally, the regions in the space of colour and magnitude offsets shown as blue and red contours in Fig. 5 were used to propagate the corresponding error to the best SFH solution. In particular, likelihoods of solutions at every position in the map of offsets were calculated based on the employed reduced $\chi^2_0$ statistic:

$$ L = e^{\frac{\nu}{2}}. $$

and the combination in quadrature of the standard deviation of likelihood values of the solutions with the shifted bins (first error source, as above) and Poisson-resampled solutions (second error source) was used to calculate the blue and red confidence regions. Then, the standard deviation of all SFH solutions inside each of these confidence regions was calculated and propagated to the error of the corresponding best solution.

Additionally, an attempt to quantify the uncertainty associated with the SFH procedure and its ability to recover observed star formation bursts is presented in Section 4.1, where the age resolution of the procedure was estimated.

### 3.4 Solution CMDs of Fornax

The final solution CMDs are presented in the middle panel of Fig. 6 compared to the observed CMD in the left-hand panel. The figure also shows the residuals $r_j$ based on Poisson statistics, defined as

$$ r_j = \frac{O_j - M_j}{\sqrt{0.5(O_j + M_j) + 1}}. $$

where $O_j$ and $M_j$, as before, are the numbers of stars in the jth colour–magnitude bin of the observed and model CMDs, respectively. In general, the two solution CMDs demonstrate a good fit to the observations, with star counts in the CMD generally recovered within $\lesssim 2\sigma$. The larger residuals in the lower MS could be caused by the low completeness in the photometry in this faint portion of the CMD, which has not been taken into account in the fit. Other strong residuals, over 2–3$\sigma$, can be seen around the position of RC, HB, and along the RGB.

The RC and HB poor fits have a known origin related to some physical prescriptions in the adopted stellar evolution models: the brightness of both RC and HB depends on the adopted conductive opacity setting, and the presently adopted version of the BaSTI library relies on – nowadays – outdated opacities (we refer to Cassisi et al. 2007 for a detailed discussion on this issue). The use of more updated conductive opacities (such as those provided by Cassisi et al. 2007 and used by Hidalgo et al. 2018) would have the effect of decreasing the brightness of RC and HB stellar models by 0.08–0.1 mag, which would contribute to improve the match between the synthetic and observed CMDs. As far as it concerns the colour distribution along the RC/HB, it is mostly driven by the mass-loss efficiency during the previous RGB stage. To properly model this feature of the synthetic CMD along the core He-burning sequence, one should introduce two free parameters in the CMD computation, which include the average mass-loss efficiency $\eta$ and dispersion around this value, as was made for instance by Savino et al. (2019). On the other hand, the code adopted in the present investigation accounts for a fixed average mass-loss efficiency during the RGB stage (i.e. $\eta = 0.2$).

The origin of the $\sim 2\sigma$ residuals in the RGB is less clear. Fig. 6 shows that model RGB is slightly steeper than the observed RGB (i.e. redder than the observed RGB below the RC and bluer at the brighter part). Since the RGB of metal-poor populations are steeper than those of metal rich ones, this could indicate that a metallicity lower than the true metallicity is being recovered for the Fornax stellar population. However, the same trend in the residuals along the RGB can be observed in other SFH derivations using the same stellar evolution
Figure 6. Left-hand panel: observed CMDs of Fornax1 (top row) and Fornax2 (bottom row). Middle panel: final solution CMDs for the both regions. Right-hand panel: residuals of subtraction of the model CMD from the observed CMD, defined as Poisson errors. Bundles 1–4 are the same as defined in Fig. 4.

library [e.g. Monelli et al. 2010a,b; Meschin et al. 2014, for the dSphs Cetus and Tucana and for the Large Magellanic Cloud (LMC)], which could indicate some systematics in the RGB colour in the models, most likely related to inaccuracies in the colour transformations which have, for example, some dependence on the exact chemical composition adopted (Pietrinferni et al. 2006). In fact, in Meschin et al. (2014), the modelled age–metallicity distribution was shown to be accurate via comparison with that derived from spectroscopic data sets, within uncertainty. In Appendix B, we perform a comparison between the metallicity distribution function (MDF) derived from the SFH of Fornax and that derived from spectroscopic measurements by Battaglia et al. (2006), Battaglia & Starkenburg (2012), and Kirby et al. (2010). It shows that indeed the former may be slightly more metal poor, especially in comparison with the MDF from Kirby et al. (2010). However, note that they warn about a possible bias in their sample due to the fact that some stars on the extreme blue end of the RGB (where asymptotic giant branch, extremely young or extremely metal-poor red giant stars would be located) did not pass the spectroscopic selection criteria. This bias would lead to underestimating the fraction of metal poor stars in that spectroscopic sample.

4 SFH OF FORNAX

The individual solutions for the SFH of Fornax1 and Fornax2 are shown overplotted on Fig. 7. The blue and red shaded error bars are calculated as a combination of three error sources as described in Section 3.3. Both regions have a similarly extended and non-uniform history of star formation. As indicated by the SFR (upper panels), the most active and extended period of star formation occurred in the oldest and most metal-poor epoch ($-1.7 < [\text{Fe/H}] < -1.3$, from the AMR on the middle panels) in both regions.

The periods lasted until $\sim 7$ and $\sim 8.5$ Gyr ago, during which they formed almost 65 and 70 per cent of stellar mass in Fornax1 and Fornax2, respectively (as can be seen by the cumulative mass functions CMF at the bottom panels). In fact, as shown later in Section 4.2, the stars formed during $\sim 5.5$–$7$ Gyr ago in Fornax1 and, possibly Fornax2, may belong to the oldest burst instead and appear shifted towards younger ages due to the observational effects. In such case the fractions of mass in stars formed in the first star formation episode could reach 70 and 80 per cent in Fornax1 and Fornax2, respectively.
of total stellar mass.\(^{650}\) The primary distinction between the two regions is the possible presence of different separate star formation episodes in the last 3 Gyr, with measured peaks at look-back times 2.8 ± 0.3, 1.8 ± 0.2, 0.9 ± 0.1, 0.5 ± 0.1 Gyr (Fornax1); and 1.1 ± 0.3, 0.2 ± 0.1 Gyr (Fornax2). While the SFR of Fornax1 seems to occur primarily in separate bursts, star formation in Fornax2 appears to be more quiescent, with indications of at least one distinct burst, around 1 Gyr ago occurring in both regions. However, note that, given the errors, most of the young peaks and low-level star formation are compatible with zero rate (more on low-level SFR in Section 4.2). Despite the somewhat different recent star formation activity, the chemical enrichment at the recent epoch is estimated to be very similar (~−0.7 < [Fe/H] < ~−0.5) in both regions. Black dots in the middle panel of Fig. 7 indicate the age bins in which < 1 per cent of total stellar mass has formed. Clearly, the statistical significance of the populations younger than 1.5–2 Gyr old is low; therefore, the ages and metallicities of the derived young stellar populations are highly uncertain.

Additionally, the overall level of star formation is higher in Fornax1, with only the oldest burst being similar in magnitude in the two regions. The two solutions have a very similar overall trend in metal enrichment during the whole lifetime of the galaxy. The enrichment proceeded monotonically up to the current epoch in the range −1.6 < [Fe/H] < −0.4.

A very interesting feature that can be clearly seen from the overlapping SFRs and CMFs on Fig. 7 is that the old star formation epoch was more extended in time in the inner region, and later ceased in both. Then, similarly, the young star formation episodes lasted till more recently in the central field of Fornax1, compared to the Fornax2 field. This kind of spatial stellar population gradient, with central regions being younger on average, is commonly seen in dwarf galaxies (Harbeck et al. 2001). What is very remarkable in the case of Fornax is that this behaviour is observed twice, first in the old star formation epoch, and then at recent times. In either case, this behaviour may be attributed to different timescales of gas depletion in the inner and outermost regions of the galaxy.

As a test of the recovered SFH, solutions with different age and metallicity bins were trialled, although the adopted configuration provided the optimal balance between the age and metallicity resolution and uncertainties. For instance, increasing the bin sizes by 1.5–2.0 times across the whole history of Fornax recovered the major bursts at similar ages compared with the solutions presented above but the stellar bursts were smoothed out decreasing the resolution of simple stellar populations. This, for example, led to broadening of the oldest and intermediate-age bursts lowering their SFR magnitude and introducing some background SFR at the intermediate epoch during ~5.5–8 in both regions. Therefore, the initially adopted bin configuration was preferred for the higher resolution of the stellar populations in the former case. Further narrowing of the bins was not adopted, as it would decrease the signal-to-noise due to the small number statistics.

The implications of choosing different bundle strategies were illustrated in the study of the SFH of Leo A in Ruiz-Lara et al. (2018). It was shown that although alternative bundle designs have a minor effect on the derived SFH of the galaxy, with most of differences included in the uncertainty, the bundle systems with the largest coverage of the observed CMD recovered best the metallicity distribution of the RGB stars derived from spectroscopy. However, here the bundle including the RC region was omitted as it resulted in higher residuals at the faint red part of the RGB and MSTO, while still not fitting the RC and HB accurately due to the issues with the stellar evolutionary models discussed above.

### 4.1 Age resolution test

In this section, we present some tests aimed at (1) verifying the accuracy of the age of \(\psi(t)\) episodes in our solutions (hereafter called \(\psi_s(t)\)) recovered with the BaSTI stellar evolutionary models, (2) testing the time resolution of SFR at different look-back times, given the uncertainties in the SFH derived from the photometric errors, the small number of stars in the observed CMDs, and the intrinsic limitation of the SFH recovery method. The tests were performed by reproducing the positions and widths of the \(\psi_s(t)\) bursts with mock bursts created with synthetic stellar populations at a constant assumed SFR. Some examples of the test in previous works include del Pino et al. (2013), de Boer et al. (2012), Hidalgo et al. (2011), and Aparicio & Hidalgo (2009).
The mock populations representing the bursts were created with different age ranges and their corresponding metallicities (as per the obtained AMR for Fornax1 and Fornax2 in Fig. 7) and a constant input SFR $\psi_i(t)$. They were extracted from the synthetic CMD used for the SFH derivation (Section 3.1) with a number of stars such that the recovered SFH closely reproduced the height of the $\psi_i(t)$ peaks. The mock CMDs for each burst were dispersed individually, according to the observational effects in the studied Fornax regions by using the results of the ASTs. The recovered (output) SFR of the mocks, $\psi_o(t)$, was computed by using the same procedure (with the same CMD parametrization), as in Section 3.2. Thus, by eliminating the uncertainties associated with matching the stellar evolutionary models with observations, distance and reddening estimates and keeping the observational errors it was possible to test the ability of the SFH code to recover the input $\psi_i(t)$ and test the age resolution of the recovered $\psi_o(t)$, ‘smoothed’ by the simulated observational effects.

The central ages of $\psi_i(t)$ bursts were taken to coincide with the mode age of the $\psi_i(t)$ bursts at different look-back times. We performed the test for star bursts in Fornax1 and Fornax2 and show the results in Figs 8 and 9, respectively. The panels show the solutions $\psi_i(t)$ in grey dashed lines, input $\psi_i(t)$ as blue shaded regions (height of the shaded regions is not indicative of the $\psi_i(t)$ SFR magnitude, but rather shows its age range) and output $\psi_o(t)$ in black lines. Most peaks of $\psi_o(t)$ and $\psi_i(t)$ can be fitted well with a Gaussian profile. The age modes $\mu$ of the solution and output bursts and their standard deviation $\sigma$ are plotted on the corresponding subscripts. In the case of the input $\psi_i(t)$, the central age and half age width are listed. The recovered bursts were normalized to the height of the corresponding solution bursts in grey. The top rows of the figures show the tests with minimum input age ranges $\Delta t_{\text{min}}$, equal to the widths of corresponding single age bins, and the bottom rows illustrate the tests with the maximum $\Delta t_{\text{max}}$, at which the recovered mock bursts of $\psi_o(t)$ just started to become wider than the solution bursts. The experiment with the $\Delta t_{\text{min}}$ input limit was used to estimate the age resolution of the mock solutions, which was calculated as the FWHM of the $\psi_i(t)$ bursts. The values of $\Delta t_{\text{max}}$ represented the upper boundary on the duration of input star formation bursts $\psi_i(t)$, i.e. the true mock bursts. The values $\Delta t_{\text{max}}$ were obtained by incrementally increasing the width of the input bursts. Therefore, the limits $[\Delta t_{\text{min}}, \Delta t_{\text{max}}]$ for every burst represent constraints of its true duration. The results for the duration and metallicities for each input mock burst can be found in Table 1.

As a result of the tests, output median ages $\mu_o$ of mock bursts were recovered close to their input positions $\mu_i$, except for the oldest bursts, which were shifted towards younger ages by a fraction of the standard deviation of the output burst $\sim 0.2 \sigma_o$ in Fornax1 and $\sim 0.1 - 0.2 \sigma_o$ in Fornax2, and the youngest burst in Fornax2, which was projected towards younger age from $\mu_i$ by $\sim 1 - 2 \sigma_o$. These results showed that the ages of the real star formation activities at intermediate epoch were recovered accurately in the solutions, while the computed SFR of the oldest and youngest bursts may be skewed to older and younger ages from the actual bursts, respectively. The systematic effect could be caused by the observational effects or, in case of the young SFR burst, could be due to the low number of stars associated with it ($\sim 300$ stars).

As can be seen from the figures, some of the shortest mock bursts ($\Delta t_{\text{min}}$) were recovered with the same $\sigma_o$ as the mock bursts with the maximum input width ($\Delta t_{\text{max}}$). Therefore, the true lower limit could not be established for those bursts (i.e. such bursts were not fully resolved and could have a shorter true duration). For example, the peaks of $\psi_i(t)$ at $\mu = 0.46, 0.91, 1.83, 4.62$ Gyr ago in Fornax1 could be caused as well by real star formation lasting shorter than the duration of the corresponding input mocks of $\Delta t_{\text{min}} = 0.025, 0.025, 0.075, 0.125$ Gyr, but our SFH code would not be able to resolve them, given the size of our age bins. None the less, we could claim the upper duration limit for all bursts. It is important to point out that the shape of the real star formation rate bursts would be more complicated than our assumed constant $\psi_i(t)$ and, hence, the above estimates of burst duration could serve only as the first-order approximations.

Finally, based on the narrowest recovered peaks, the age resolution of our SFR, defined as $\Delta t_{\text{min}} = \text{FWHM}(\psi_i(t))$ (see Table 1), was found to deteriorate as $\sim 0.2, 0.2, 0.3, 0.9, 2$ Gyr at the look-back times $\sim 0.5, 0.9, 1.8, 4.6, 11.3$ Gyr for Fornax1; and 0.2, 0.3, 0.7, 2.5 Gyr at the ages 0.1, 1.1, 4.7, 11 for Fornax2.

### 4.2 Intermediate star formation epochs

As seen from the upper panel of Fig. 7, the complex SFH of Fornax is characterized by distinct bursts, separated by the epochs of a low-level SFR. Even though the error bars at these epochs are compatible with a null SFR, we further investigated our solutions to find out whether the low-level SFR is real or a result of stars being ‘dispersed away’ from their main star formation epoch by the observational effects in the process of SFH derivation.

To test this low-level SFR, we did an additional experiment with mock stellar populations that represented the main measured star formation bursts, not including the low-level star formation between them. For this purpose, we adopted the mock bursts from the age resolution test (Section 4.1). The input duration of the bursts was set to the maximum values $\Delta t_{\text{max}}$ from Table 1, so as to have as many stars ‘dispersing away’ by the observational effects as possible. In contrast to the age resolution test, the CMDs for all mock bursts were combined into a single CMD, from which the SFH was calculated by following the standard procedure (Section 3).

Fig. 10 shows that the mock bursts alone can account for the $\psi_i(t)$ activity at $\sim 8–5.5$ Gyr ago, between the oldest and intermediate star formation epochs, in both Fornax1 and Fornax2; however, when setting the width of the burst at $\sim 4.6$ Gyr ago to $\Delta t_{\text{min}}$ (not shown on the figure), the same background activity is still recovered with the mocks in Fornax1, but not in the case of Fornax2, suggesting that Fornax2 could have experienced star formation $\sim 8–5.5$ Gyr ago if the duration of the burst $\sim 4.6$ Gyr ago was close to $\Delta t_{\text{min}}$. Also, interestingly, in the case of $\Delta t_{\text{max}}$, the observed background SFR was not recovered in the period $\sim 3.75–2$ Gyr ago in Fornax1 and Fornax2 and in between the youngest bursts in Fornax1, indicating the possibility of active low-level star formation at that time as well.

In conclusion, the test suggests that these regions may have experienced active star formation at the periods indicated by our SFH solutions (Fig. 7), including low-level star formation at some time $\sim 8–5.5$ Gyr ago in Fornax2, $\sim 3.75–2$ Gyr ago in Fornax1 and Fornax2 and at even younger ages in Fornax1. However, a halt in the star formation activity after the first burst may have occurred in Fornax1 and possibly also in Fornax2.

### 4.3 On the impact of a different stellar evolutionary framework on the SFH retrieval

Many SFH studies of LG dwarfs, including the previous SFHs of Fornax, have been conducted using the BaSTI models, thus relying on the same theoretical framework in this study was required for proper comparison between the findings in Fornax and other LG galaxies.
The implications of adopting different models are worth noting. For example, the updated version \(^5\) of the BaSTI library (Hidalgo et al. 2018), which incorporates atomic diffusion, more updated physical inputs such as new conductive opacity tabulations and a different mass-loss efficiency during the RGB stage, is expected to provide more accurate and reliable stellar models for the advanced evolutionary stages with respect to the models used presently (Pietrinferni et al. 2004), while producing the same results for the H-burning stage, as discussed in detail in Section 3.4. A detailed comparison of the updated version with the one used here and other recent libraries is presented in Hidalgo et al. (2018). The efficiency of atomic diffusion has an effect on the old age dating: in general, for a given physical framework, in passing from models with no diffusion to fully efficient diffusion, the estimated ages decrease by 0.7–0.8 Gyr at old ages (i.e. ages larger than about 8 Gyr). Different libraries are expected to produce somewhat different results for the intermediate- and young-age populations depending on the adopted efficiency for the core convective overshooting during the H-burning stage, among other parameters (see Gallart, Zoccali & Aparicio 2005 for a general overview in the context of the SFH procedure). Some works have explicitly compared SFH results obtained adopting different star evolution libraries. For example, Monelli et al. (2010a,b) did a comparison of SFHs at old ages obtained with the Girardi et al. (2000) and BaSTI libraries finding that the former produced slightly younger (\(\sim\)1 Gyr solutions at old age); Dolphin (2012) performed a comparison between Padova (Girardi et al. 2000), BaSTI, and Dartmouth (Dotter et al. 2008) libraries for both old and intermediate age bursts using synthetic populations.

In the case of the stellar population analysed in this work, we estimate that, if models with different convective overshooting were adopted, the main star formation bursts found to occur \(\sim 4.6\) and \(\sim 10.6\) Gyr ago would likely not be affected, while the bursts estimated to occur in the past 2 Gyr could be recovered at an older age depending on the adopted overshoot prescription. However, given the low statistical weight of star formation bursts in the recent \(\sim 2\) Gyr, this would not change the main conclusions of the work. If atomic diffusion was assumed to be fully efficient (the diffusion is not implemented in BaSTI library used in this work), the two oldest, main star formation bursts could be shifted by up to \(-0.8\) Gyr towards the current epoch. Given the age resolution test (see Section 4.1), such

---

\(^5\)The updated BaSTI library is available at: http://basti-iac.oa-abruzzo.inaf.it/index.html
extreme change would still be within the uncertainty of the solutions for the oldest burst in Fornax1 and Fornax2, but just outside the error margin of the bursts at 4.6 ± 0.5 Gyr ago and 4.6 ± 0.4 Gyr ago in Fornax1 and Fornax2, respectively.

4.4 Comparison with previous results
In this section, we compare our SFH with previously published determinations from CMD analysis and with chemical evolution scenarios inferred from spectroscopy of individual stars.

Table 1. Tests with mock bursts in Fornax1 and Fornax2 (column 1). Ages of the input mock were centred on look-back ages $\mu_{i,\text{min}}$ (column 2) and $\mu_{i,\text{max}}$ (column 6) with their widths $\Delta t_{i,\text{min}}$ (column 3) and $\Delta t_{i,\text{max}}$ (column 7), respectively, where the minimum widths were defined by the age bins of the solution (i.e. at the limit of age resolution) and the maximum limit was such that the peak could still be resolved, before becoming wider than the solution burst. The recovered widths of the mock populations (FWHM ($\psi_o(t)$)) are stated as $\Delta t_{\text{o,\text{min}}}$ (column 4), defined as the age resolution, and $\Delta t_{\text{o,\text{max}}}$ (column 8), defined as the upper boundary on the burst duration. Each input mock burst has the associated metallicity range $\Delta z_{i,\text{min}}$ (column 5) and $\Delta z_{i,\text{max}}$ (column 9).

<table>
<thead>
<tr>
<th>solution</th>
<th>$\mu_{i,\text{min}}$ (Gyr)</th>
<th>$\Delta t_{i,\text{min}}$ (Gyr)</th>
<th>$\Delta t_{i,\text{max}}$ (Gyr)</th>
<th>$\Delta z_{\text{min}}$</th>
<th>$\mu_{i,\text{max}}$ (Gyr)</th>
<th>$\Delta t_{\text{i,\text{max}}}$ (Gyr)</th>
<th>$\Delta t_{\text{o,\text{max}}}$ (Gyr)</th>
<th>$\Delta z_{\text{max}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Fornax1</td>
<td>0.46</td>
<td>0.03</td>
<td>0.16</td>
<td>0.00620–0.00660</td>
<td>0.48</td>
<td>0.15</td>
<td>0.15</td>
<td>0.00578–0.00661</td>
</tr>
<tr>
<td></td>
<td>0.91</td>
<td>0.03</td>
<td>0.18</td>
<td>0.00478–0.00515</td>
<td>0.89</td>
<td>0.23</td>
<td>0.24</td>
<td>0.00452–0.00578</td>
</tr>
<tr>
<td></td>
<td>1.81</td>
<td>0.08</td>
<td>0.34</td>
<td>0.00515–0.00547</td>
<td>1.85</td>
<td>0.45</td>
<td>0.40</td>
<td>0.00426–0.00578</td>
</tr>
<tr>
<td></td>
<td>4.56</td>
<td>0.13</td>
<td>0.85</td>
<td>0.00223–0.00240</td>
<td>4.63</td>
<td>1.00</td>
<td>0.84</td>
<td>0.00184–0.00220</td>
</tr>
<tr>
<td></td>
<td>11.13</td>
<td>0.25</td>
<td>2.05</td>
<td>0.00067–0.00071</td>
<td>10.00</td>
<td>6.00</td>
<td>5.26</td>
<td>0.00010–0.00137</td>
</tr>
<tr>
<td>Fornax2</td>
<td>0.24</td>
<td>0.03</td>
<td>0.15</td>
<td>0.00547–0.00578</td>
<td>0.28</td>
<td>0.14</td>
<td>0.18</td>
<td>0.00550–0.00661</td>
</tr>
<tr>
<td></td>
<td>1.13</td>
<td>0.05</td>
<td>0.26</td>
<td>0.00373–0.00395</td>
<td>1.18</td>
<td>0.75</td>
<td>0.65</td>
<td>0.00500–0.00700</td>
</tr>
<tr>
<td></td>
<td>4.69</td>
<td>0.13</td>
<td>0.65</td>
<td>0.00240–0.00255</td>
<td>4.56</td>
<td>0.88</td>
<td>0.94</td>
<td>0.00240–0.00270</td>
</tr>
<tr>
<td></td>
<td>10.88</td>
<td>0.25</td>
<td>2.51</td>
<td>0.00067–0.00071</td>
<td>10.63</td>
<td>4.25</td>
<td>2.76</td>
<td>0.00010–0.00081</td>
</tr>
</tbody>
</table>

Figure 9. The same as Fig. 8, but for Fornax2.
4.4.1 Star formation history

The photometry and SFH results presented in this paper for Fornax1 and Fornax2 are widely consistent with the current view of Fornax as a galaxy with a complex SFH continuing to almost the present time. However, the extremely deep and precise CMDs resulting from the high signal-to-noise and high-resolution HST images allowed us to provide important new details about the SFH by accurately dating and constraining the duration of the main SFR bursts. This information allows us to characterize Fornax as a galaxy that has experienced distinct SFH bursts separated by intermediate periods of null or low-level activity.

The analysis of two fields at different galactocentric distances allowed us to see the features of the stellar population gradient, reported by every wide-field study of stellar populations of Fornax starting from Stetson et al. (1998) and followed by Battaglia et al. (2006), de Boer et al. (2012), de Boer et al. (2013), del Pino et al. (2013), del Pino et al. (2015), Bate et al. (2015), and Wang et al. (2019b). In general, the oldest stars are found to be distributed quite uniformly within the tidal radius of the galaxy and the intermediate- and young-age stars are represented by a larger fraction in the centre.

The study by Buonanno et al. (1999) argued that Fornax experienced a second major burst of star formation at ~4 Gyr ago, as judged by the separate SGBs in their CMD. This has been confirmed by the recovered peaks of $\psi(t)$ in our solutions. The age of the narrow episode at 4.6 ± 0.4 Gyr ago agreed very well in multiple averaged solutions for Fornax1 and Fornax2. However, this feature of the Fornax SFH, suggested long ago by the qualitative analysis of a high-quality WFPC2/HST CMD in the latter study, was not found in the more recent quantitative SFH determinations (de Boer et al. 2012; del Pino et al. 2013; Weisz et al. 2014), while indications of it were reported by Coleman & de Jong (2008).

The work by del Pino et al. (2013) presents the first Fornax SFH derived from CMDs that reach the oMSTO in a range of galactocentric distances. Their SFH is superimposed on ours in Fig. 11 and the map of the corresponding field locations is shown in Fig. 1. The SFR in the regions IC2, OC and, especially, the most crowded region IC1, are very smooth and do not show signs of the distinct star formation bursts found in our solutions. The most striking disagreement is found between the SFH of Fornax1 and IC1, which are, in fact, the most directly comparable regions at all ages considering the asymmetry of the spatial distribution of the stellar populations younger than ~2 Gyr, as shown in del Pino et al. (2015). The SFH from del Pino et al. (2013) demonstrates a $\psi(t)$ peak ~8 Gyr ago that lasted for the whole intermediate-age epoch, while the active period in Fornax1 stopped after ~7 Gyr ago and, likely, no stars were formed until ~5.1 Gyr ago. This cannot be attributed to the different stellar evolutionary models, as the same BAsTI isochrones were used in our work. Moreover, none of the regions from del Pino et al. (2013) show the prominent burst at 4.6 ± 0.4 Gyr ago, which could have been smoothed out by a shallower photometry. In comparison, the FWHM resolution in the present study is ~0.2, 0.2, 0.3, 0.9, 2 Gyr at the look-back times ~0.5, 0.9, 1.8, 4.6, 11.3 Gyr in the most crowded region Fornax1, while del Pino et al. (2013) reported the values equivalent to FWHM of ~1.4 Gyr at an age 3 Gyr ago and ~4 Gyr at the oldest ages achieved by their photometry. Aside from the major differences, our results for Fornax1 and Fornax2 agree with the results from del Pino et al. (2013) in terms of the declining $\psi(t)$ at increasing elliptical radii, as well as towards younger ages at all radii and in the total recovered metallicity range of $z(t)$.

4.4.2 Metallicity distribution

In independent spectroscopic studies of Fornax its metal enrichment was found to follow a radial gradient. The VLT/FLAMES data for 562 RGB stars from Battaglia et al. (2006) showed that the range of metallicity in a stellar sample located within $r < 3^\prime$ was −1.4 ≤[Fe/H] ≤ −0.4, while for $3^\prime < r < 13$ metallicities within −2.1
≤[Fe/H]≤−0.2 were measured. In comparison, the metallicity range in our solutions is −1.7 < [Fe/H] < −0.4 for both Fornax1 (r ∼ 3) and Fornax2 (r ∼ 13), and agrees with the previous results within the errors indicated in Fig. 7. Some of the most metal-poor stars with [Fe/H] = −1.7 ± 0.1 in the Battaglia et al. (2006) RGB sample had an inferred age >10 Gyr old. The metallicities of [Fe/H] = −1.0 ± 0.1 and [Fe/H] = −0.7 ± 0.1 were associated with 2–5 and 1.5–2 Gyr old isochrones. Our solutions from photometric data agree well with these estimates obtained with spectroscopic measurements, where each of the abundance levels mentioned above corresponds to the age of the onset of each of the major star formation episodes. Additionally, Battaglia et al. (2006) reported the presence of stars that were formed a few hundred Myr ago in the central regions (see also Mapelli et al. 2009 for evidence that these stars are genuine young stars and not blue straggler stars). This is also in agreement with the recovered SFH of the galaxy in our study.

In Appendix B we make a detailed comparison between the MDF of Fornax derived from the SFH, and that obtained from the spectroscopic studies by Battaglia & Starkenburg (2012) and Kirby et al. (2010). It is interesting that the MDF for Fornax1(2) derived in this paper is bimodal (with peaks at [Fe/H] ≈ −1.4 and [Fe/H] ≈ −0.9, see Fig. B2), indicating that these regions have at least two different metallicity populations associated with the distinct star formation bursts. This is in contrast with the spectroscopic MDFs which appear to be dominated mostly by the single peak around [Fe/H] ≈ −1.0. In the Appendix, we demonstrate that simulating the spectroscopic metallicity errors in the metallicities estimated here from photometry smoothens out the bimodality, resulting in an MDF closely resembling the spectroscopic one (Fig. B3, with possibly a difference in the metallicity zero-point.

Interestingly, populations with distinct spatial and chemo-dynamic properties that result in a decomposition of the spectroscopic MDF into several metallicity components, have been inferred in Fornax by previous studies. Based on the spectroscopy of Battaglia et al. (2006) with data covering the tidal radius, and using Bayesian analysis, Amorisco & Evans (2012) separated the Fornax stellar populations into metal poor (MP; [Fe/H] ≈ −1.60), intermediate metallicity (IM; [Fe/H] ≈ −0.95) and metal rich (MR; [Fe/H] ≈ −0.65), which robustly fitted the properties of the stellar spatial distribution, metallicity and line of sight velocity. The MDFs of Fornax1(2) derived from our SFH (Fig. B2) have two peaks centred around [Fe/H] ≈ −1.40 (−1.35) and [Fe/H] ≈ −0.90 (−0.74), which approximately agree with the MP and IM + MR populations inferred in the previous work. While it is hard to distinguish potential IM and MR components in our MDFs, we can only speculate that if the third population was present in our MDF solution, it would be likely more pronounced in Fornax1, at the centre of the galaxy, than in Fornax2, just outside of the core region. This agrees with the high-metallicity tails of the MDFs of Fornax1(2), as can be seen in the figure. Moreover, the three distinct episodes of star formation found in our solutions for Fornax1 (old, intermediate and young) support the hypothesis of three distinct spatial and chemo-dynamic populations of Amorisco & Evans (2012). Hendrickx et al. (2014) calculated abundances of α-elements and MDF of Fornax based on their spectroscopic CaT data at two outer regions in between one core and one tidal radii along the major axis, slightly extending the coverage of the Battaglia et al. (2006) sample with a similar selection of the RGB stars. They also identified three distinct sub-populations by fitting an analytic function. Although the sub-populations differed in metallicity from the ones in Amorisco & Evans (2012), their two most metal rich populations ([Fe/H] = −0.97, −1.45) peaked at [Fe/H] similar to the two peaks in this work ([Fe/H] ≈ −0.90 (−0.74) and [Fe/H] ≈ −1.40 (−1.35)), while their most metal poor population at [Fe/H] ≈ −2.11 end was almost absent from our solutions, given that the Fornax1(2) are dominated mostly by the intermediate-metallicity and metal-rich stars.

Driven by the photometric data and results of CMD fitting from del Pino et al. (2013), the study by Piatti et al. (2014) investigated the MDF in their regions IC1, IC2 and OC (as on Fig. 1) and observed bimodality in all, with the MDF peaks centred around [Fe/H] ≈ −1.25 and [Fe/H] ≈ −0.95 in the central field IC1, similarly to the MDFs presented here.

5 DISCUSSION. THE BURSTY SFH: WHAT DID REIGNITE STAR FORMATION?

Simulations of SFHs of isolated dwarf galaxies often predict episodic bursts of star formation owing to the effects of Supernova feedback (Shen et al. 2014; Stinson et al. 2007), which temperately deplete the central regions of star-forming material. This process leads to quenching star formation as gas flows outward into a hot halo, where it cools until subsequently sinks to the centre of the potential well, triggering a new burst of star formation. Other models (e.g. Benítez-Llambay et al. 2015; Verbeke et al. 2014) invoke the combined effects of cosmic reionization, supernova feedback and late gas accretion to explain major bursts of star formation in dwarf galaxies after some quiescent epoch following an initial star formation burst. The SFH of Fornax is not unlike some predictions by these simulations.

Among other possible scenarios, a merger of two galaxies is also expected to induce star formation. This has been studied in various simulations by Teyssier, Chapon & Bournaud (2010), Cox et al. (2008), Di Matteo et al. (2008), Miñhos & Hernquist (1994), albeit for galaxies with mass much larger than Fornax. Indications of mergers were found in dwarf galaxies such as Andromeda II (Amorisco, Evans & van de Ven 2014) and Sextans (Cicu´endez & Battaglia 2018). Various features observed in Fornax also show that it may have experienced a merger event. Amorisco & Evans (2012) review previous observations that hint at the presence of subpopulations in Fornax, indicated by age-dependent asymmetries in the isophotes of stellar populations (Wang et al. 2019b; del Pino et al. 2015; Stetson et al. 1998; Irwin & Hatzidimitriou 1995; Eskridge 1998; Hodge 1961b), stellar overdensities (Bate et al. 2015; de Boer et al. 2013; Coleman et al. 2005, 2004), as well as multiple age and chemodynamic components (Battaglia et al. 2006; Saviane et al. 2000). Therefore, as it was hypothesized by Coleman et al. (2004) and further supported by the numerical simulations of the merger in Yozin & Bekki (2012), based on proper motions available at the time, and considering the kinematic properties of different subpopulations in Amorisco & Evans (2012), the observed spatial distribution of the stellar populations could result from a merger of a bound pair of dwarf galaxies in a relatively recent period of time, ∼2–3 Gyr ago. Such scenario could be in agreement with the bursts starting to occur 2.8 and 1.9 Gyr ago in Fornax1 and continuing at a later time in Fornax2 after the shocks have reached the radius of this region.6 Other evidence based on chemokinematic properties of Fornax has been used to propose a possible merger scenario ∼6 Gyr ago in del Pino et al. (2017). These predictions coincide with the ages of the onset of the intermediate-age star formation bursts in the

6 As noted before (Section 4), the bursts at ≤2 Gyr ago have a low statistical weight; therefore, this is only a speculative statement assuming the bursts are true.
SFH solutions for Fornax1 and Fornax2. Lastly, an early merger has also been suggested by Leung et al. (2020) to explain the unusually large mass of stars deposited in five GCs of Fornax, a mismatch between the metallicity of the GC5 stars and the old field stars and the differential dynamics of different stellar populations. In their scenario, Fornax would have merged with a $\sim 3$ times less massive galaxy $\sim 10$ Gyr ago. This hypothesis could explain what contributed to the oldest star formation episode at the median age $\sim 10.6$ Gyr ago in our solutions for the central region Fornax1.

Additionally, there is theoretical and observational evidence that tidal interaction between a host and a dwarf galaxy may also ignite star formation in either of the two galaxies (e.g. Larson & Tinsley 1978, for an early reference). Indications of tidally activated star formation after interaction with a smaller satellite galaxy have been found in a M31-M33 pair and our own MW interacting with the Sagittarius dSph. Bernard et al. (2012) showed that M31 experienced a burst of star formation in its disk outskirts that could have been triggered by a close approach of M33, reproduced by numerical simulations in McConnachie et al. (2009). Ruiz-Lara et al. (2020b) have discussed recurrent episodes of enhanced star formation in the MW disk that coincide with independently inferred pericentric passages of the Sagittarius dSph galaxy (e.g. Purcell et al. 2011; de la Vega et al. 2015; Laporte et al. 2019). In turn, the SFH of the Sagittarius core shows star formation enhancements at similar times (Siegel et al. 2007). Also, numerical simulations of the SFH and orbit of the Carina dSph (Pasetto et al. 2011), which has an extended SFH and well separated star formation episodes (Smecker-Hane et al. 1994; Monelli et al. 2003) similar to Fornax, showed correlation between the peaks in SFR and pericentric passages around the MW. Interestingly, the new orbits calculated by Patel et al. (2020) based on Gaia proper motions and including the influence of the Magellanic Clouds indicate a second last pericentric passage around 5 Gyr ago when the influence of the Magellanic Clouds is taken into account, in agreement with the similar intermediate-age star formation burst of Carina observed in the SFH derived by Santana et al. (2016). Finally, a close passage of Leo I about the MW (entering the virial radius of the MW $\sim 2$ Gyr ago, and pericentre at $91 \pm 6$ kpc about 1 Gyr ago, Sohn et al. 2013) may be linked with an enhancement of the Leo I star formation activity some 2 Gyr ago and lasting to about 1 Gyr ago (Ruiz-Lara et al. 2020a).

To study whether pericentric passages about the MW could be responsible for the Fornax bursty SFH, we calculated the orbit of Fornax in the potential of the MW. For the properties of Fornax we use the Gaia proper motion of Fritz et al. (2018) (but note that the proper motion of Gaia Collaboration et al. 2018 is very similar when compared with the systematic error of 0.035 mas yr$^{-1}$). Orbit integration is performed with the GALPY$^7$ package of Bovy (2015), where for the MW potential we use variants of MWPotential14, with the halo mass modified to obtain MW masses which are more in-line with recent mass measurements. The values of $\m_{\text{MW}} = 1.45 \times 10^{12} \, \text{M}_\odot$ and $1.1 \times 10^{12} \, \text{M}_\odot$ are considered to explore the range of MW masses (Fritz et al. 2020; Bland-Hawthorn & Gerhard 2016).

A preliminary investigation included potential encounters of Fornax with other galaxies brighter than ultra-faint dwarfs in the absolute luminosity-based definition of Simon (2019). We found that within the past 1 Gyr, a close encounter of Fornax with one of the dwarf galaxies is unlikely. The closest passages within the past 2.5 Gyr are possible with Crater II (1.3 Gyr ago) and with Sculptor (1.6 Gyr ago); however, given the observational uncertainties, there is a chance of only $\sim 10$ per cent for the dwarfs to approach each other as close as 20 kpc, which significantly exceeds their stellar extent. These scenarios are not considered here in detail due to increasingly large errors at the lookback times longer than 1 Gyr. Therefore, we cannot constrain well any dwarf-dwarf encounters that could have caused the star formation bursts of Fornax.

Our main analysis concerns the separation distance between Fornax and the MW. This orbit integration included dynamic friction and two mass estimates for Fornax that bracket most mass estimates in the literature. The lower value of $2 \times 10^8 \, \text{M}_\odot$ represents the mass contained within the stellar extent of Fornax (Diakogiannis et al. 2017) and the higher mass of $2 \times 10^{10} \, \text{M}_\odot$ was obtained by Read, Walker & Steger (2019) by using a cored, untruncated NFW halo. The results are similar to the ones in Fritz et al. (2018) and are shown in terms of separation distance $r(t)$ as a function of lookback time in Fig. 12.

In the models with the more massive MW, the two most recent pericentric passages correspond to the two most recent substantial bursting epochs in Fornax1 that occurred 1.8 $\pm$ 0.2 and 4.6 $\pm$ 0.4 Gyr ago. For the higher mass of both galaxies, which are considered most likely, the last pericentre occurred 1.7 $^{+2.9}_{-0.8}$ Gyr ago at a distance of 76 $^{+65}_{-24}$ kpc and the second latest was 5.1 $^{+5.3}_{-0.7}$ Gyr ago at 87 $^{+75}_{-25}$ kpc. The timing of the passage $\sim 5.1$ Gyr ago coincided with the onset of the most intense burst, which ceased in the next $\sim 0.5$–1 Gyr. The difference of ages of the two passages is more similar to the age difference of the two considered bursts when the low Fornax mass is adopted. In this case, both bursts occur exactly at the pericentre, while for the large Fornax mass, either passage is off by 0.3–1 Gyr.

Note that, given the measurement errors, a wide range of times for the pericenter passages is possible (as indicated by the boxes depicted with thin lines in Fig. 12). Additionally, the calculations included the assumption of a spherical halo (Bovy et al. 2016), while improved simulations should account for deviations from spherical symmetry (the most prominent effects could be due to the LMC). These considerations would produce somewhat different results, as shown in Erkal et al. (2019), Sohn et al. (2017). In fact, the former study found a probability of Fornax once being a satellite of the LMC, and possible effects of such interaction should also be considered. Very recently, Patel et al. (2020) performed orbit calculations including both Magellanic Clouds and it was found that

\footnotesize
\textsuperscript{7}Available at https://github.com/jobovy/galpy.
\normalsize
although Fornax is not bound to the Clouds, their effect changed
gas times of Fornax, as compared to our results. They show the
most recent peri-centric passages at ∼1.5 and 5 Gyr ago for their
more massive MW model (see fig. 2 therein).

The latter and our calculations obtain peri-centric passages that
coincide with significant bursts of star formation in the second half
of Fornax history. It must be noted that more stringent constraints on
the mass of the MW and shape of its potential can change the inferred
dating of the peri-centric passages drastically. Therefore, the only
unambiguous conclusion from this evidence follows that if the tidal
interactions are assumed to be the mechanism that triggered recurrent
star formation in Fornax, then this would be likely if the galaxies had
the masses $M_{MW} = 1.45 \times 10^{12} M_\odot$ and $M_{For} = 2 \times 10^{10} M_\odot$ out of
the four considered mass combinations. If true, this would provide
tantalizing evidence that the MW tidal effects on Fornax at a distance
of ∼ 80 kpc may be the cause of these bursts. This distance, in fact,
places Fornax well within the MW virial radius at the time of the
tidal interaction.

If the assumption is confirmed, then Battaglia, Sollima & Nipoti
(2015) conclusion that tidal effects in Fornax were insignificant and
did not disturb its stellar or dark matter distributions provides support
to the possibility that neither ram pressure nor tidal stripping have
been able to remove the gas that may still have been gravitationally
bound to Fornax after the first epoch of star formation, but rather,
might have tidally induced star formation.

6 SUMMARY

We derived the SFH in two small fields near the centre of the Fornax
dSph with much deeper photometric data than used previously, which
resulted in a significantly improved age resolution, sufficient to date
individual episodes of star formation and estimate their intrinsic
duration. This work is thus providing an unprecedented sharp view
of the evolutionary history of the Fornax stellar populations.

The work is concentrated on two HST/ACS regions of Fornax
(∼3.4 × 3.4 squared arcmin each) at radial distances of 2.8 and 12.7
from its centre. We employed the CMD fitting technique to analyse
the stellar population composition of the galaxy and derived its SFH.
Careful tests using mock stellar populations have allowed us to put
tight limits on the duration of the bursts, and on the possibility of low-
level star formation between the bursts. We have found that the first
and most abundant populations of old, metal-poor stars, accounting
for up to ∼70 per cent of stars in Fornax1 and ∼80 per cent in
Fornax2, were formed in a star formation episode ∼10−11 Gyr ago,
that lasted between ∼4 and 6 Gyr in Fornax1 and up to ∼4.25 Gyr in
Fornax2. The second strong and sharp burst ∼4.6 Gyr ago contributed
∼10−15 per cent to the mass in stars during 1 Gyr at most in both
fields. Finally, the subsequent significant episode ∼1.8 Gyr ago in
Fornax1 and younger intermittent bursts in both fields added the
remaining mass in stars, with only 1−3 per cent of it contained in the
stars formed less than 1 Gyr ago. Due to the small number of young
stars, we cannot claim accurate ages or metallicities for the recent
star formation activity. The results of the mock stellar population
tests indicate that both regions (and more securely Fornax1) have
likely ceased star formation activity between ∼8 and 5.5 Gyr ago,
while low level star formation seems to have occurred between the
intermediate-age and the young bursts.

The SFHs for two regions of such small size are probably not
representative of the whole core region or the whole galaxy, given the
non-uniform distribution of different stellar populations. However,
the main star formation bursts are observed in both regions, with the
same age, similar inferred duration and relative burst intensity. This
is a further indication of the reliability of the derived scenario.

The two regions also seem to be connected by the spatial age
gradient, typical for dwarf galaxies, where the central region Fornax1
tends to experience more prolonged star formation episodes than
Fornax2. It is very remarkable that this outside-in quenching of star
formation, commonly observed in dwarf galaxies (e.g. Hidalgo et al.
2013, and references therein), is observed in both main epochs of star
formation: both the old burst and the young star formation activity
have a longer duration in the innermost field. That is, if star formation
would not have been reinitiated in Fornax at intermediate ages, the
galaxy would still show the usual stellar population gradient observed
in dwarf galaxies.

We have discussed the possible mechanisms that may be respon-
sible for the separated events of star formation in Fornax. Numerical
simulations often predict a bursty SFH in field dwarf galaxies, and
thus the same characteristic in Fornax could be just an example of
such behaviour owing to the normal secular evolution of this dwarf
galaxy. A second possibility is that it is related to a dwarf–dwarf
merger: indications of an early one have been discussed in Leung et al.
(2020) and would explain the peculiar Fornax GC population, while
a more recent merger could explain the multiple chemodynamic
and age components (Amorisco & Evans 2012). We discuss a third
possibility in relation to possible repeated tidal interactions of Fornax
with the MW. We show that, for the mass of the MW favoured by
the most recent determinations, and masses of Fornax within current
(wide range) estimates, the most likely times for the two most recent
pericenter passages of Fornax are in very good agreement with the
times of the intermediate-age and main young star formation
burst in Fornax. This adds another example of tidally induced star
formation to the mounting evidence of such occurrence, which can be
particularity well characterized in the LG, where not only current, but
also intermediate-age and old star formation events can be accurately
dated.

As a final remark, we would like to note that this work demon-
strates that ultra-precise CMDs of LG galaxies, such as the ones
that will be obtained by the LSST project for dwarfs in the MW
system, will allow a much sharper view of the details of their SFHs
across all epochs of their evolution, including accurate times of the
star formation events and their duration. Such studies of LG
galaxies, together with complementary information coming from
other present and upcoming galactic archaeology surveys such as
Gaia, will significantly boost our knowledge of evolution profiles of
galaxies leading to the LG as we currently observe it.

ACKNOWLEDGEMENTS

We would like to thank the anonymous referee for the detailed
report and insightful comments, which we believe greatly helped
to improve the quality of the manuscript. We thank Giuseppina
Battaglia for kindly providing their spectroscopic catalogues and
useful discussions, and Andrés del Pino for providing the results
of their SFH work on Fornax. This work has made use of data
stored at Barbara A. Mikulski Archive for Space Telescopes (STScI),
NASA Astrophysics Data System Bibliographic Services, and the
NASA/IPAC Extragalactic Database (NED), as well as ESO Science
Archive Facility. Used facilities: HST (ACS).

TKF, CG, MM, and TRL acknowledges support through Ministerio de Economía, Industria y Competitividad and AEI/FEDER
(UE) grant no. AYA2017-89076-P, and by the Consejería de Economía, Industria, Comercio y Conocimiento of the Canary
Islands Autonomous Community, through the Regional Budget
(including IAC projects ‘Galaxy Evolution in the Local Group’ and ‘TRACES’). TRL additionally acknowledges financial support through grant nos AYA2016-77237-C3-1-P (RAVET project), a Juan de la Cierva – Formación grant (FJCI-2016-30342) and support from a Spinova grant (NWO) awarded to A. Helmi, grants (AEI/FEDER, UE) AYA2017-89076-P, AYA2015-63810-P, AYA2016-77237-C3-1-P (RAVET project), and a MCIU Juan de la Cierva – Formación grant no. (FJCI-2016-30342). SC acknowledges support from Premiale INAF MITIC, from INFN (Iniziativa specifica TAsP), PLATO ASI-INAF contract no. 2015-019-R0 and no. 2015-019-R.1-2018.


DATA AVAILABILITY

The data used in this article are available in Barbara A. Mikulsi Archive for Space Telescopes at http://archive.stsci.edu/. The data are part of the proposal 13435 (PI: M. Monelli). The text of the proposal is available at http://archive.stsci.edu/proposal_search.php?id=13435&mission=hst.

REFERENCES

Astropy Collaboration et al., 2018, AJ, 156, 123
Baade W., Hubble E., 1939, PASP, 51, 40
Harris C. R. et al., 2020, Nature, 585, 357
Letarte B. et al., 2010, A&A, 523, A17
Letarte B. et al., 2010, A&A, 523, A17
McConnachie A. W. et al., 2009, Nature, 461, 66
McConnachie A. W., 2012, AJ, 144, 4
McKinney W., 2010, in van der Walt S., Millman J., eds, Proceedings of the 9th Python in Science Conference, Scipy Conference, Austin, TX, p. 56

Downloaded from https://academic.oup.com/mnras/article/502/1/642/6067380 by Det Kongelige Bibliotek user on 10 June 2021
Table A1. Log table indicating the images taken as part of HST/ACS proposal ID 13435 used in this paper. Columns (left to right): names of image files, target of observations indicated in the image header, date and time of observations, right ascension and declination coordinates (J2000), ACS filter, exposure time, and modified Julian date. Top half of data is for Fornax1, bottom is for Fornax2.

<table>
<thead>
<tr>
<th>Image</th>
<th>Target</th>
<th>Date</th>
<th>Time</th>
<th>RA(^h), m(^s), (s)</th>
<th>DEC(^\circ), (\prime), (\prime)</th>
<th>Filter</th>
<th>Exposure (s)</th>
<th>Julian</th>
</tr>
</thead>
<tbody>
<tr>
<td>jcb009c7q.flt</td>
<td>Fornax1</td>
<td>2014-01-04</td>
<td>03:07:04</td>
<td>2:03:42.55</td>
<td>-34:31:43.25</td>
<td>F475W</td>
<td>80</td>
<td>5.6661E+04</td>
</tr>
<tr>
<td>jcb009c9q.flt</td>
<td>Fornax1</td>
<td>2014-01-04</td>
<td>03:54:37</td>
<td>2:03:42.43</td>
<td>-34:31:44.58</td>
<td>F814W</td>
<td>850</td>
<td>5.6661E+04</td>
</tr>
<tr>
<td>jcb009c9q.flt</td>
<td>Fornax1</td>
<td>2014-01-04</td>
<td>04:21:06</td>
<td>2:03:42.55</td>
<td>-34:31:43.25</td>
<td>F475W</td>
<td>80</td>
<td>5.6661E+04</td>
</tr>
<tr>
<td>jcb101cnq.flt</td>
<td>Fornax1</td>
<td>2014-01-04</td>
<td>05:37:33</td>
<td>2:03:42.55</td>
<td>-34:31:44.58</td>
<td>F814W</td>
<td>850</td>
<td>5.6661E+04</td>
</tr>
<tr>
<td>jcb101cnq.flt</td>
<td>Fornax1</td>
<td>2014-01-04</td>
<td>06:09:14</td>
<td>2:03:42.55</td>
<td>-34:31:43.25</td>
<td>F814W</td>
<td>850</td>
<td>5.6661E+04</td>
</tr>
<tr>
<td>jcb111v7q.flt</td>
<td>Fornax1</td>
<td>2013-12-27</td>
<td>14:21:06</td>
<td>2:03:42.55</td>
<td>-34:31:43.25</td>
<td>F814W</td>
<td>80</td>
<td>5.6653E+04</td>
</tr>
<tr>
<td>jcb111v9q.flt</td>
<td>Fornax1</td>
<td>2013-12-27</td>
<td>14:26:20</td>
<td>2:03:42.43</td>
<td>-34:31:44.58</td>
<td>F475W</td>
<td>850</td>
<td>5.6653E+04</td>
</tr>
<tr>
<td>jcb111vq.flt</td>
<td>Fornax1</td>
<td>2013-12-27</td>
<td>14:52:47</td>
<td>2:03:42.55</td>
<td>-34:31:43.25</td>
<td>F814W</td>
<td>850</td>
<td>5.6653E+04</td>
</tr>
<tr>
<td>jcb006d2q.flt</td>
<td>Fornax2</td>
<td>2014-06-22</td>
<td>03:34:03</td>
<td>2:04:39.53</td>
<td>-34:32:41.60</td>
<td>F814W</td>
<td>80</td>
<td>5.6830E+04</td>
</tr>
<tr>
<td>jcb006d4q.flt</td>
<td>Fornax2</td>
<td>2014-06-22</td>
<td>05:05:41</td>
<td>2:04:39.53</td>
<td>-34:32:41.60</td>
<td>F814W</td>
<td>850</td>
<td>5.6830E+04</td>
</tr>
</tbody>
</table>

**APPENDIX A: HST/ACS OBSERVATIONS LOG**

Table A1 details the HST/ACS observations that were used in this paper.
APPENDIX B: METALLICITY DISTRIBUTION FUNCTION

This Appendix describes the comparison of the MDF of Fornax obtained from the SFH retrieval procedure and from some of the previous spectroscopic studies of Fornax.

Among the considered spectroscopic samples with the spatial coverage including Fornax1(2) (medium resolution Kirby et al. 2010; medium and high resolution Letarte et al. 2010; Ca II triplet spectroscopy from Battaglia & Starkenburg 2012), the samples from (Kirby et al. 2010; Battaglia & Starkenburg 2012, hereafter, K10, B12) were used as the most numerous. Since the number of stars measured spectroscopically in the Fornax1 and Fornax2 fields is small, for the comparison we are using larger samples of stars within the Fornax central region. Although no substantial stellar population gradients are expected in the Fornax core that can affect the metallicity distribution, particularly of RGB stars, which are older than \( \approx 2 \) Gyr (del Pino et al. 2015), the spatial variation of the MDF samples is investigated here.

To test how the MDFs can be affected by random sampling of stars within an elliptical radius equal to that of Fornax2, 1000 Monte Carlo samples (46 stars each, which is similar to the number of stars spectroscopically measured in Fornax1 + 2 by K10) were drawn from the full spectroscopic samples. Fig. B1 shows, for K10 (top) and B12 (bottom), the MDF of the stars measured within the core radius (black line), the median MDF from the Monte Carlo samples (blue line) and the MDF of stars measured in the fields Fornax1 + 2 (for K10) or the stars measured within elliptical shells covering the regions Fornax1 and Fornax2 (for B12, as only 6 stars from their sample coincided with the exact regions Fornax1 and Fornax2). The median of the Monte Carlo samples is consistent with the MDF samples of Fornax1 + 2 within the 90 per cent confidence interval, meaning that the metallicity data are not expected to have a significant spatial bias.

In spite of the fact that the results of the conducted test showed little dependence of the MDF on spatial positions of stars, the spectroscopic data used for the analysis below were selected from the elliptical shells enclosing Fornax1 and Fornax2 and not from the whole core to further reduce the possibility of a spatial bias. Fig. B2 illustrates the comparison of the MDF derived from the SFH solutions in this work and the MDF from spectroscopic samples with similar number of stars. The photometric selection criteria for the stars in the solution CMD were adjusted according to the spectroscopic data sets: K10 includes stars just above the RC (\( M_{F814W} < -0.7 \) according to Fig. 2); B12 covers the \( \sim 2 \) brightest magnitudes of the RGB only. As can be seen from Fig. B2, Fornax1 and Fornax2 MDFs derived from the SFH are similar along most of the metallicity range, with Fornax1 being slightly more metal rich. Both are characterized by a bimodal MDF, with a metal-poor peak at \( \text{[Fe/H]} \approx -1.4 \) and a metal-rich peak at \( \text{[Fe/H]} \approx -0.9 \), which correspond with the stellar populations that are \( \sim 6–14 \) Gyr old (Fornax1) and 8–14 Gyr old (Fornax2), respectively. This distribution is unlike the K10 and B12 counterparts, which peak at \( \text{[Fe/H]} \approx -1.0 \) and are skewed towards the more metal-
rich end. Another clear disagreement, particularly with K10, occurs at $-1.8 < \text{[Fe/H]} < -1.3$, with this range represented by a larger number of stars in Fornax1(2).

The MDF showing at least two distinct modes in this work is consistent with the derived SFH, characterized by major old and intermediate bursts of star formation, likely separated by an epoch of lower activity around 6–8 Gyr ago, and metallicity increasing with time as a consequence of star formation. At this point, we consider the following question: if the true MDF of Fornax was indeed bimodal, could the measurement errors of the spectroscopic metallicities erase the bimodality?

To answer this question, we simulated the spectroscopic errors (as drawn from the K10 and B12 data tables) in the MDFs derived from the SFH. The results are shown in Fig. B3, where the Fornax1(2) MDFs represent medians of 1000 randomly drawn Gaussian error samples ($\sigma \approx 0.14$ dex in K10 and B12) combined with the MDFs from the SFH. This step resulted in smoothing of the Fornax1(2) distributions and partial resolution of the bimodality disagreement, particularly in comparison with B12.

With regards to the second disagreement with the spectroscopy, the metal-poor stars at $-1.8 < \text{[Fe/H]} < -1.3$ are represented more in the Fornax1(2) solutions even after introducing the errors, as can be seen in Fig. B3. The difference is more prominent in the case of K10. Note that K10 warn about a possible bias in their sample because some stars on the extreme blue end of the RGB (populated mainly by asymptotic giant branch, extremely young or extremely metal-poor red giant stars) did not pass the spectroscopic selection criteria. This bias could underestimate the fraction of metal poor stars. Finally, it is possible that a zero-point in the metallicity scale could account for this remaining difference.

From the above analysis, we conclude that the MDF of the solutions presented in this work is reasonably consistent within the errors with the MDF spectroscopically derived for the core region of Fornax, with a possible zero-point difference. The analysis shows that the MDF derived for Fornax1(2) from the fit to the CMD may have a higher resolution in metallicity than achieved by the K10 and B12 data sets, revealing a possible bimodality in the metallicity distribution of the stars in the core of Fornax.

This paper has been typeset from a \TeX/\LaTeXX file prepared by the author.

The spectroscopic metallicities presented in B12 were kindly made available by G. Battaglia.