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Distribution of Interplanetary Dust Detected by the Juno Spacecraft and Its Contribution to the Zodiacal Light

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Abstract The Solar System is home to a cloud of dust that orbits the Sun and makes its presence known by virtue of scattered light (Zodiacal Light) that can be seen after dusk and before dawn. Within this cloud are bands of dust orbiting near the ecliptic plane, evidenced by an excess of scattered light at discrete ecliptic latitudes. Dedicated dust detectors borne by spacecraft in transit of the solar system have detected few such particles of the appropriate size owing to limited detector aperture and sparsity of the population. Thus, the distribution, origin, and orbital evolution of the dust in these bands remains a mystery. A star camera aboard the Juno spacecraft traveling from Earth to Jupiter recorded interplanetary dust impacts on the spacecraft in numbers sufficient to characterize the spatial distribution of such particles for the first time. The observed distribution is consistent with a primary source of dust particles sharing the Mars orbit plane between Earth and the 4:1 resonance with Jupiter. We propose that the primary distribution is scattered by orbital resonances with Jupiter via the Kozai-Lidov (KL) effect into a secondary population at higher inclination to the ecliptic. The measured dust distribution, occupying a volume uniquely determined by the orbital elements of Mars and KL scattering, accounts for the observed variation of the Zodiacal Light with ecliptic latitude. Our results provide a compelling alternative to the prevailing theory of the origin and evolution of interplanetary dust observed at low ecliptic latitudes.

Plain Language Summary The Zodiacal light is sunlight reflected by dust in the inner solar system. Variations in the Zodiacal light with ecliptic latitude reveal discrete bands of dust orbiting near the ecliptic plane. The Juno spacecraft, in transit from Earth to Jupiter, recorded a sufficient number of impacts with this dust to characterize their distribution in space for the first time. The radial distribution of the dust suggests a primary source of dust with the orbital elements (inclination and eccentricity) of Mars, scattered into a secondary population at higher inclination to the ecliptic. The resulting population accounts for the observed variation of Zodiacal light with ecliptic latitude.

1. Introduction

The space between the planets is populated by objects of various sizes in orbit about the Sun, ranging from minor planets and asteroids to tiny nanograins, with smaller objects more numerous than larger objects over many decades of object radii. The smallest particles are blown out of the solar system by radiation pressure while larger particles drift inward or outward under the influence of Poynting-Robertson drag (Roberson, 1937; Wyatt & Whipple, 1950) and other nongravitational effects that influence their motion (Burns et al., 1979; Horányi, 1996; Lhota et al., 2016; Mann, 2014; Morfill & Grun, 1979). The interplanetary cloud of small particles ("dust") must be continuously replenished to offset these losses. This cloud of dust particles can be observed in reflected sunlight (the Zodiacal light), emission of thermal radiation (Hauser et al., 1984; Reach et al., 1997), and samples collected in the earth's upper atmosphere (Love et al., 1994) or in orbit about earth (Love & Brownlee, 1993), but the distribution and origin of the interplanetary dust remains an unsolved problem. However, the well documented variation of the Zodiacal light with ecliptic latitude (Low et al., 1984) evidences the persistence of discrete bands of dust concentrated near the ecliptic plane, implicating a source or sources orbiting in the same plane[s]. It has been widely assumed that the dust bands originate with a few asteroid families that deliver the dust to the inner solar system with the orbital inclination needed to reproduce features in the Zodiacal light curves (Dermott & Nicholson, 1989; Dermott et al., 1984, 1990; Espy et al., 2009; Grogan et al., 1997; Nesvorný et al., 2002, 2003, 2006; Sykes &
attention (Nesvorný et al., 2010, 2011; Pokorny, et al., 2014), particularly with regard to larger dust particles that can be tracked entering Earth's atmosphere as radar meteors.

Spacecraft equipped with specialized instruments designed to detect and characterize interplanetary dust particles (IDPs) have traversed the solar system in search of IDPs, detecting the more numerous small particles (<1 μm) that originate beyond our solar system (Humes, 1980; Grün et al., 1994; Horányi et al., 2008; Krüger et al., 2010; Landgraf et al., 2000). Larger IDPs of the size responsible for the Zodiacal light (~1–100 μm) are infrequently collected by these instruments owing to the limited detection aperture (typically ~0.1 m²) common to such instruments. Thus, larger IDPs have not been detected in-situ in sufficient numbers to characterize their distribution, and as a result, models of the interplanetary dust cloud fitted to variations of the Zodiacal light are poorly constrained.

The Juno spacecraft was launched in August 2011, bound for Jupiter orbit insertion on July 4, 2016, on a trajectory that carried it out to the asteroid belt where a deep space maneuver (DSM) sent it back toward a rendezvous with Earth en route to Jupiter (Bolton et al., 2017; Stephens, 2018). Juno was not instrumented with a dedicated dust detection instrument, but it did carry a complement of four star cameras to provide accurate attitude determination for the magnetic sensors (Connerney et al., 2017). This system is known as the Advanced Stellar Compass (ASC), and it employs four Camera Head Units (CHUs), designated A-D, that are collocated with the magnetic sensors. One of these cameras (D) was configured to search for asteroids, recording observations of luminous objects that could be tracked across multiple images in transit across the star background (Andersen et al., 2006; Soberman et al., 1974). This search for “nonstellar objects” (NSOs) turned up a number of luminous objects that were identified as impact ejecta–spallation products excavated from the body of the Juno spacecraft upon impact with an IDP traveling at ~5 to ~15 km s⁻¹ relative velocity (Benn et al., 2017). IDPs detected in this manner are thought to range from ~1 to ~100 μm in size, representative of the dust population associated with the Zodiacal light (Leinert et al., 1990). Here we present the number of IDPs detected throughout Juno’s traverse of the solar system, describe the distribution and orbital evolution of these IDPs in space, and show that this distribution can explain the observed variation of Zodiacal light with ecliptic latitude.

2. Observations

Juno’s ability to detect IDP impacts was entirely unanticipated and serendipitous, but in retrospect the spacecraft and instrument were well designed to work together to do just that. Juno was designed to use solar power at Jupiter’s orbital distance of 5.2 astronomical units (AU), requiring ~60 m² of solar arrays to provide little more than ~600 W of electrical power to the spacecraft and subsystems. The spacecraft also spins (30 or 60 s period) about the +z axis, nearly aligned with the parabolic communication antenna. The magnetometer investigation’s star cameras, mounted at the end of one of Juno’s three solar arrays, scan the sky with a 13° by 18° field of view (FOV) in the opposite direction, generally antisunward (Connerney et al., 2017). Traveling from Earth to Jupiter, Juno was moving with a velocity of >5 km s⁻¹ relative to objects in near-circular Keplerian orbit it might encounter on its way; such objects would impact the antisunward side of the solar arrays. Relatively large spallation products (~0.1 to ~1 mm in size) ejected from the IDP impact sites were imaged over multiple frames and precisely tracked utilizing the parallax afforded by spacecraft rotation. The detection method is described in detail in Benn et al. (2017) and additional information is provided in the supporting information (“Methods”). The evolution of ejecta counts as an IDP impact event, time-tagged, recorded in instrument memory, and stored on the spacecraft awaiting downlink opportunities, typically available once a week during cruise phase. The spacecraft and star camera, in this configuration, together act as a dust detector with ~60 m² of aperture, affording a vast increase in collection efficiency relative to dedicated particle detectors flown in space previously.

We confine our discussion to IDPs of the size required to eject spallation products from the target (antisunward side of the solar arrays), loosely characterized as IDPs in the size range of ~1 to ~100 μm judging by the size of the spallation products (Benn et al., 2017). The rapid decrease in IDP density with increasing size (e.g., Grün et al., 1994) suggests that most of the IDPs we detect with this method will fall to the low end of
this size range (∼few to 10 μm). In what follows, when we refer to IDPs, we mean IDPs of the size range that we are capable of detecting. The much smaller dust particles of interstellar origin that have been recorded in significant numbers by dedicated dust detectors flown previously (Humes, 1980; Grün et al., 1994; Horányi et al., 2008; Krüger et al., 2010; Landgraf et al., 2000) are evidently too small to eject spallation products large enough for the CHU to track. This conclusion is consistent with the existence of multiple extended periods of time within which no IDP impact events were recorded; this detection method is insensitive to the flux of small interstellar grains.

The number of IDP impact events experienced by Juno along its trajectory from Earth to Jupiter is illustrated in Figure 1. A series of filters was applied to the NSO candidates to isolate IDP impacts, some applied onboard the spacecraft (e.g., object motion relative to the star background) and others in ground data analysis (supporting information). This process identified a total of 15,278 IDP impacts experienced by Juno during its 5-year cruise to Jupiter. In this process we have undoubtedly discarded some actual IDP impact events that left debris trajectories too close to other luminous objects (stars, planets, and star clusters) or produced impact ejecta that receded from the spacecraft too quickly.

Interpretation of these data require a careful consideration of the operation of the instrument and Juno’s position relative to the ecliptic plane, and Mars, on its journey to Jupiter (Figure 2). The Juno mission plan included a close Earth flyby (EFB) to impart additional velocity to the spacecraft for its journey to Jupiter. This required a passage from Earth out well beyond the orbit of Mars and into the asteroid belt, where a pair of DSMs were performed prior to Juno’s return to Earth for a close flyby (Stevens, 2018). CHU-D was programmed to search for NSOs shortly after the DSMs. Prior to EFB, the search was conducted in the ecliptic plane, as illustrated in Figure 2. After EFB, Juno’s trajectory was inclined to the ecliptic by ∼4.5°.

It is instructive to examine the IDP detections shown in Figure 1 as a function of heliospheric radial distance, in the inner solar system, separately, for the inbound (Juno spacecraft traveling toward EFB) and outbound (toward Jupiter) segments of Juno’s trajectory. This part of the journey spans radial distances of 0.88 AU (Juno’s perihelion) to 2.25 AU. In Figure 3, top, the NSO search begins at the right of the figure, at 2.08 AU, following the final DSM, as Juno exits the main asteroid belt and prior to passage of the radial distance of the 4:1 resonance with Jupiter. Time increases to the left through Juno’s perihelion at 0.88 AU. The bottom panel in Figure 3 illustrates the outbound passage from perihelion to the inner asteroid belt, time increasing from left to right.
Figure 2. Illustration of the Juno spacecraft’s trajectory from launch (August 2011) through Jupiter orbit insertion (July 2016). The orbital position of Jupiter and Mars, and their inclination to the ecliptic, is indicated at 20-day intervals. Juno’s trajectory is illustrated (green) with thickened segments illustrating transit of the radial distance range bounded by Mars perihelion (1.36 AU) and aphelion (1.67 AU). Juno remained in the ecliptic while inbound to Earth from DSM and ventured out of the ecliptic while outbound from Earth to Jupiter. Earth and Mars are illustrated as positioned in their orbits at the time of Jupiter orbit insertion (JOI). AU, astronomical units; EFB, Earth flyby.

Figure 3. IDP impacts per day as a function of radial distance from the Sun during the Juno spacecraft’s transit, in the ecliptic plane, from the last deep space maneuver (DSM 2) to perihelion (upper) and from perihelion, past Earth flyby, outward (and out of the ecliptic) into the main asteroid belt (lower). Data gap indicated as in Figure 1; blue shaded region identifies instrument operation in a different mode where a correction factor (4x) was applied adjusting the impact rate for the reduced instrument sample rate. For context, asteroid density as a function of radial distance is illustrated with a gray shaded histogram (right scale). The radial positions of mean orbital resonances with Jupiter are identified (ratios, black and blue dashed lines), as are orbit perihelion and aphelion for Earth (purple) and Mars (red). IDP, interplanetary dust particle.
Both panels in Figure 3 illustrate the variation in IDP impacts as a function of heliocentric radial distance, albeit with numerous small gaps in time. The curves join estimates of the number of impacts per day from eight-hour segments of time-tagged ejecta events. The small gaps between data segments are the result of ASC instrument memory capacity limitations and infrequent data downlink opportunities. When a large number of events are experienced per unit time, the instrument memory fills to capacity and no additional time-tagged events can be logged; the gaps represent the interval of time subsequent to saturation of instrument memory and prior to the availability of memory following NSO data downlink. This instrumental limitation and the regular cadence of downlink opportunities accounts for the regularly spaced data gaps most noticeable along the inbound trajectory between \( \sim 1.6 \) and \( \sim 2.1 \) AU, where the largest impact rates were encountered.

3. Analysis

Juno experienced a greater dust impact rate along the inbound transit to EFB, during which the spacecraft remained in the ecliptic plane, relative to the outbound passage traveled at higher ecliptic latitude. Figure 3 presents the data in its most basic form, plotting the number of IDP impacts recorded per day as a function of heliospheric radial distance. To characterize the IDP population density, we require knowledge of the IDP orbits (discussed below) and the projected area of the spacecraft presented to the incoming flux. Allowing for the spacecraft relative velocity and swept volume (supporting information), the number density of IDPs observed close to the ecliptic plane varies from a few to \( 8 \times 10^{-13} \) m\(^{-3} \). This agrees well with the density of the “asteroidal” component of interplanetary dust quoted by Grün et al. (1985, 1997) for IDP masses of \( 10^{-6} \) to \( 10^{-10} \) g (\( \sim \) few to 50 \( \mu \)m radii). The number density observed at higher ecliptic latitude (4.5°) is about one third of that. Neither panel evidences the \( 1/r \) dependence expected of a population of dust drifting inward under the influence of Poynting-Robertson drag (Wyatt & Whipple, 1950; Leinert & Grun, 1990; Mann et al., 2014).

It is also clear that the distribution of events displays more structure during the ecliptic plane transit (upper panel), including an intriguing association with major Jupiter mean motion orbital resonances (MMRs) and the radial range traversed by planets Earth and Mars. In particular, Juno experienced a dearth of impact events (for \( \sim 3 \) weeks) while inbound after DSM, leaving the main asteroid belt, and approaching the strong 4:1 MMR with Jupiter (Figure 3, top panel, read right to left). A large number of events were recorded immediately upon passage inside of the 4:1 resonance, and again after passage inside of the 5:1 resonance (note a single 8-h period during which no events were recorded just beyond the 5:1 resonance). The dramatic influence of these mean motion resonances on the IDP population suggests that IDPs are efficiently scattered out of orbit at the MMR and that the IDPs we detect reside in low-inclination, near-circular Keplerian orbits. The calculated IDP density (supporting information) adopts near-circular Keplerian orbits for the IDPs, along with knowledge of spacecraft orientation, to properly account for the relative velocity and directionality of incoming IDPs.

Turning attention to Earth and Mars, as planets that might be expected to sweep dust from their path, we note that IDPs reside comfortably within a putative Mars sweeping region, whereas no IDP impact events were recorded between Earth’s aphelion and semi-major axis (“Earth Orbit-a”). We conclude that Earth does effectively sweep its path clear of IDPs with orbits evolving inwards. In contrast, a large impact rate (\( \sim 200/day \)) was recorded shortly after passage into Mars’ sweeping region and a similarly large rate was experienced just before and after passing through the Mars orbit plane (“Mars Orbit”). Therefore, Mars does not effectively sweep its path clear of IDPs. We do, however, note a systematic decrease in impact events associated with passage through the radial distance equal to Mars’s semi-major axis (“Mars Orbit-a”). This may reflect a modest depletion of the IDP population near the mean orbital radius of Mars. No such associations are evident during Juno’s outbound passage at higher ecliptic latitude.

The precipitous drop in the IDP population associated with the 4:1 resonance with Jupiter, and at Earth aphelion, suggests that the IDP source resides within this radial range (\( \sim 1.02-2.065 \) AU) and at low inclination to the ecliptic. It has been widely assumed, if not universally accepted, that the dust interior to the asteroid belt is brought in by asteroid families and/or comets with eccentric orbits, but if so, such particles would necessarily initially share the orbit eccentricity of the parent body and not survive passage of the
Figure 4 shows a substantially reduced impact rate (~10/day) observed in transit through the asteroid belt outwards to Jupiter. This traverse reveals the existence of multiple extended periods of time within which no IDP impact events were recorded. Those associated with the main asteroid belt act largely as a sink, not a source, of IDPs. This is most apparent within the radial range beyond ~2.75 to ~3.25 AU at several lengthy intervals are found within which no IDP impact events were recorded. The lack of detections during these intervals and on approach to 5 AU also demonstrates that our detection method is insensitive to the omnipresent flux of small interstellar grains detected in large numbers by dedicated dust instruments flown throughout the solar system (Mann, 2010). They are evidently too small to eject visible spallation products from the solar arrays. Figure 4 (and to some extent, Figure 3, top panel) also provides some evidence in support of a concentration of dust impact events near orbital resonances, where IDPs may be trapped for an appreciable amount of time (thousands of years) before being ejected (Liou & Zook, 1997). The relatively smaller interplanetary dust population residing in orbits beyond the 4:1 resonance may well be sourced by asteroids and comets in eccentric orbits.

3.1. Model of Interplanetary Dust Bands

The IDP population observed during the outbound passage after EFB, at an inclination of 4.5° to the ecliptic (Figure 3, bottom panel), is less dense and apparently not significantly influenced by Jupiter resonances or potential absorbers such as Mars. We believe that the population sampled by Juno at higher inclination is a secondary population of IDPs scattered from a primary population with low inclination to the ecliptic by the Kozai-Lidov (KL) mechanism (Kozai, 1962; Lidov, 1962; Naoz, 2016). The primary population we associate with Mars, which has an orbital inclination of 1.85°, is expected to fill a volume of space bounded by ±1.85° to the ecliptic and, if drifting inward and outward under the influence of nongravitational forces (Burns et al., 1979; Horányi, 1996; Lhotka et al., 2016), filling a disc as illustrated in Figure 5.

The KL effect is a dynamical gravitational interaction between an orbiting mass (IDP) and a distant third body, in this case, Jupiter. The interaction leads to a periodic exchange of orbital eccentricity and inclination; the maximum inclination (relative to the Mars orbit plane) available to an IDP with Mars’ eccentricity is:

\[ i = A \cos \left[ (1 - e^2)^{0.5} \cos(i_m) \right] = 5.5° \]
where $e$ is Mars’ orbital eccentricity ($0.0934$) and $i_m$ is Mars’ inclination relative to Jupiter’s orbital plane ($1.44^\circ$).

The secondary population, if scattered via the KL effect acting on the primary population of IDPs, will fill a disc-like volume as illustrated in Figure 5 bounded by $\pm 7.4^\circ$, summing Mars’ orbital inclination of $1.85^\circ$ and the KL scattering angle of $5.5^\circ$. KL oscillations may be expressed in terms of a pair of binaries, an inner binary consisting of masses $m_1$ and $m_3$, and an outer binary consisting of masses $m_2$ and $m_3$. In the test particle limit ($m_2 \rightarrow 0$) appropriate to IDPs, to within an order of magnitude, the period for KL oscillation over a complete cycle from maximum eccentricity (and minimum inclination) to maximum inclination (and minimum eccentricity) and back is given by Antognini (2015):

$$t_{KL} \equiv \frac{8}{15\pi} \left(1 + \frac{m_1}{m_3}\right) \left(P_{out}^2 \rho_{out}^2 \left(1 - e^2\right)^{3/2}\right) \approx 13,000 \text{ years}$$

where $m_1$ is the mass of the Sun ($1.99 \times 10^{30}$ kg), $m_3$ is the mass of Jupiter ($1.90 \times 10^{27}$ kg), $P_{out}$ is the orbital period of Jupiter (12 years) and $P_m$ is the orbital period of an IDP (1.88 years) leaving the Mars system. IDPs may be expected to achieve maximum inclination to the ecliptic in one fourth of this time, that is, in about 3,000 years, long before Poynting-Robertson drag can remove an IDP of $\sim 10 \mu$m size (see Figure S4).

3.2. Implications for the Zodiacal Light

The IDP distribution we measure, if distributed according to our model, ought to be consistent with the observed variation of Zodiacal light with ecliptic latitude, as measured by instruments on Earth orbiting satellites such as the Cosmic Background Explorer and the Infrared Astronomy Satellite (IRAS). Here (Figure 6) we model the intensity of light reflected by the IDP population with that measured (Espy et al., 2009) by IRAS, using a linear approximation to the distribution of IDPs that we measured (supporting information) and assuming that the density of IDPs that we measure is uniform throughout the primary and secondary population.
regions illustrated in Figure 5. The IRAS measures the Zodiacal light in a direction orthogonal to a radius vector from the sun as illustrated in Figure 5, left cut, and with that viewing geometry the Zodiacal light intensity is most sensitive to the nearby dust illuminated with more intense sunlight. The two closely spaced peaks near the ecliptic (±1.4°) result from the primary IDP population distributed within ±1.85° of the ecliptic. The secondary distribution within ±7.4° of the ecliptic populated by KL scattering of the primary population yields the two peaks (±9.5°) commonly referred to as “the 10° peaks” in the literature.

The approximate total number of IDPs contained in the model distribution illustrated in Figure 5 may be estimated by direct integration (supporting information), yielding $N \sim 2 \times 10^{22}$ particles. The light reflected from this model distribution is consistent with measurements (Levasseur-Regourd & Dumont, 1980) of light scattered by the IDP population (supporting information). Assuming an average IDP mass of ~15 ng, corresponding to a 10 μm radius particle, the detectable IDP population has a mass of ~2.8 × 10^{11} kg. To sustain this population against loss by Poynting-Robertson drag over a time span of ~100,000 years requires a replenishment rate of ~30 kg s⁻¹.

4. Discussion

The population of IDPs we measure presumably reflects a steady state between the supply of IDPs and the loss rate, among other factors that influence detection, such as spacecraft and IDP orbital characteristics, detector (spacecraft) projected area, image sensitivity, and so on. Little is actually known about the supply of IDPs, although it has long been assumed that the IDPs responsible for the Zodiacal light (~1–100 μm) are delivered to the inner solar system by comets and asteroids in elliptical orbit about the sun. The evidence in support of this assumption is largely circumstantial, as this population has until now escaped measurement. Morfill and Grun (1979) summed up the situation well: “an alternative to this is not perceptible to us.”

The loss rate of IDPs depends on many processes affecting their orbital motion as well as those that lead to mass loss, such as sputtering, sublimation, and collisional fragmentation. Orbital motion is influenced by forces that depend critically on particle size, such as (electromagnetic) radiation pressure, Poynting-Robertson drag, solar wind drag, and Lorentz forces on charged dust particles (Burns et al., 1979; Horányi, 1996; Lhotka et al., 2016; Morfill & Grun, 1979). Radiation pressure is effective in quickly clearing some small particles (≤~1 μm) out of the solar system, depending on mineralogy and light scattering properties (Burns et al., 1979), but even smaller particles (≤~0.1 μm) are cleared primarily by solar wind drag and Lorentz forces (Horányi, 1996; Lhotka et al., 2016; Morfill & Grun, 1979). Radiation pressure and Poynting-Robertson drag are the dominant nongravitational forces acting upon the Zodiacal IDPs (Burns et al., 1979), yielding to the Yarkovsky effect for larger objects, depending on spin period. Poynting-Robertson drag and the Yarkovsky effect (assuming prograde rotation) both circularize and collapse particle orbits, limiting orbital lifetime (supporting information). Thus, in-situ measurements of smaller (<1 μm) IDPs and inferences drawn from radar studies of larger objects (>400 μm) entering Earth’s atmosphere are of limited use in diagnosing the Zodiacal IDP population; these populations evolve under the influence of very different dynamical effects. The smaller particles measured in-situ are mostly of interstellar origin (Landgraf et al., 2000; Mann, 2010).

Our analysis of the IDP population diverges from the usual approach in one key respect: we assume the IDPs we detect are in near-circular Keplerian orbits. This assumption follows naturally from the observation that IDP density at low inclination to the ecliptic drops dramatically just beyond the 5:1 and 4:1 mean motion resonance with Jupiter; IDPs are evidently efficiently ejected from such orbits. This assumption is in contrast to that prevailing since Dermott et al. (1984) originally attributed the peaks in the Zodiacal light curve to dust bands associated with specific asteroid families (Eos, Themis, and Koronis). These families have orbital inclinations that match what was needed to reproduce the variations observed in the Zodiacal light curve. Thus began modeling efforts of increasing sophistication and complexity, using specific asteroid families (Dermott et al., 1990; Dermott & Nicholson, 1989; Espy et al., 2009; Grogan et al., 1997; Nesvorný et al., 2002, 2003, 2006; Sykes, 1987, 1990) and/or JFCs, HTCs and OCCs (Nesvorný et al., 2010; Nesvorný et al., 2011; Pokorný, et al., 2014) to deliver IDPs to the inner solar system. The population we measured demonstrates that IDPs are ejected from orbits just beyond the mean motion resonances with Jupiter; dust delivered by crossing asteroids or comets would not survive inward migration across the resonances. Of
course, there are a great many asteroid and comet families, and Reach et al., (1997) questioned why only a few of them produced dust bands. Why not, for example, the Nysa, Flora, and Eunomia/Io asteroid families?

Our IDP model fits these new (and unique) observations and benefits from simplicity. We show that the population we measure along two invariant planes near the ecliptic is consistent with a single source of IDPs with (initially) the orbital elements of Mars, scattered via the KL mechanism into a secondary volume determined by Mars’ orbital elements. KL scattering occurs well before radial migration of IDP orbits under the influence of nongravitational forces. We show that this distribution is sufficient to explain the observed variation of the Zodiacal light with ecliptic latitude (“Zodiacal dust bands”). There are no adjustable parameters in this model other than a scaling parameter for the light reflected from individual particles. In contrast, the alternative hypothesis, whereby the Zodiacal dust bands are fed by asteroids and comets transiting the inner solar system, requires identification of one or more asteroid or comet families for each pair of peaks in the Zodiacal light curve. With well over 100 individual asteroid families now identified by examination of their proper elements (Milani et al., 2014; Zappalà et al., 1995), it is perhaps not surprising that some can found with orbital inclinations close to that needed to match the Zodiacal dust bands. However, it is difficult to understand why only some asteroid families—the ones with the required inclinations—contribute dust bands, and not the many others.

What is lacking in our model is a mechanism by which the primary IDP population is created with the inclination and orbital eccentricity of Mars. Can the planet Mars source this population? A large flux of dust particles is observed by the Mars-orbiting Maven spacecraft (Andersson et al., 2015) but the source of the dust remains a mystery; the observed distribution is thought to be inconsistent with an origin at either Mars satellite Phobos or Diemos. While Mars does suffer episodic planet-wide dust storms (Martin & Zurek, 1993; Zurek & Martin, 1993) and Mars dust (so-called “rocket dust”) is launched to extraordinary altitudes (Heavens et al., 2019; Spiga et al., 2013), escape from Mars would require overcoming escape velocity ($\sim 5 \text{ km s}^{-1}$), for which no mechanism has yet been identified.

Alternatively, the Mars satellites Phobos and Diemos could potentially source dust ejected from the regolith with modest velocity ($v_{\text{esc}} \geq 12 \text{ m s}^{-1}$ for escape from Phobos, less for Diemos). This problem has been considered by many authors following S. Soter’s contemplation of a dust ring (torus) formed in orbit about Mars by micrometeoroid bombardment of the Mars satellites (Soter, 1971). Dust of appropriate size ($\sim 10$ to $\sim 100 \mu m$) launched with modest velocity from either satellite can reside in a torus about Mars for thousands of years (Krivov, 1994; Krivov & Hamilton, 1997; Soter, 1971) but escape from the Mars system remains an obstacle and micrometeoroid bombardment of Phobos and Diemos evidently yields to little ejecta mass by several orders of magnitude (Krivov, 1994). Thus we at present can identify no specific delivery mechanism by which the primary population is seeded.

5. Conclusions

We have measured the density of IDPs as a function of radial distance from the Sun between Earth and Jupiter, and we propose that the circumsolar dust bands of IDPs in our solar system originate in association with the planet Mars, initially sharing the orbital elements of Mars. IDPs are scattered to higher inclinations via the KL mechanism and subsequently evolve both inward toward the Earth and outward toward the asteroid belt. We detected a dramatic decrease in IDP density just beyond the 5:1 and 4:1 mean motion resonance with Jupiter and conclude that outward evolving IDPs in near-circular Keplerian orbits are ejected from such orbits. Fewer IDPs are detected within and beyond the main asteroid belt.

The capability to detect dust impacts on the solar arrays continues with Juno in orbit about Jupiter since July 5, 2016. However, few impacts are detected by the ASC during passage through the ring plane; ring particles in prograde orbits would thus far impact the sunward side of the arrays, producing ejecta beyond the CHU FOV. As the orbit evolves in extended mission, dust impacts occurring during a periapsis pass will occur on the darkened side of the arrays, once again visible to the ASC’s camera D. It will then be possible to compare the ring particle impact rate with that observed by the Waves investigation (Ye et al., 2020), which is sensitive to smaller particles impacting the spacecraft on any surface. The Waves instrument detects particle impacts by recording the impulsive change in spacecraft potential associated with the high temperature plasma as it evolves from the impact site.
The IDP environment recorded by Juno on its journey to Jupiter has potentially significant implications for the design and qualification of space flight hardware as well as mission design. Solar-powered spacecraft with large projected area have ventured successfully to the outer planets (e.g., Rosetta and Juno) but these spacecrafts used solar arrays backed with relatively massive aluminum honeycomb and carbon composite substrates. The flux of IDP swept by these spacecrafts was absorbed by the antisunward side of the solar array substrate. Will the new generation of ultra-thin, flexible solar arrays proposed for upcoming deep space missions tolerate IDP impacts as well?

Data Availability Statement

All data used in this article are available in the main text and in the supporting information, as well as in the permanent archival data repository, Zenodo (Jorgensen et al., 2020). JPL manages the Juno mission for the principal investigator, Scott Bolton, of SWRI. The Juno mission is funded by NASA and is part of the New Frontiers Program managed at NASA’s Marshall Space Flight Center in Huntsville, Alabama.

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