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On-chip nano-electro-mechanical switching of deterministic single photons

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Abstract: We demonstrate a nano-electro-mechanical single-photon router integrated with semiconductor quantum emitters, showing an extinction ratio of >20 dB and operation speed of MHz with insertion loss of 0.67 dB and footprint <30 µm². © 2019 The Author(s)


Quantum photonic integrated circuits enable scaling up quantum optics experiments by integrating multiple functionalities, such as deterministic photon-emitter interfaces, routers, and detectors, on a chip. To this end, a key requirement is to develop reconfigurable circuitry that can be operated at a speed compatible with the emitter qubit coherence time (0.1–1 µs) and introduces neither loss nor emitter decoherence. The existing tuning mechanisms, based on thermo-optic and electro-optic effects, do not meet these demands as they are either too slow or require very large footprints, limiting both scalability and efficiency. In this work, we demonstrate a different approach to on-chip routing by using nano-electro-mechanical devices integrated with deterministic single-photon sources, i.e., InAs quantum dots (QDs) in GaAs. Unlike refractive-index-tuning methods, such electro-opto-mechanical interaction can be very strong and material-independent [1], thus offering the advantage of a much smaller attainable device footprint, lower insertion loss, faster switching speed, and better scalability.

Fig. 1. (a) Scanning electron micrograph image of the nano-electro-mechanical router on GaAs membrane with embedded InAs quantum dots. (b) Schematic of the device working principle. (c) Transmission at two output ports as a function of gap distance.

The integrated photon router is fabricated in 160-nm thick GaAs membranes with embedded InAs QDs, as shown by a scanning electron micrograph in Fig. 1(a). It consists of two identical suspended nanobeam waveguides each attached to capacitive nanomechanical actuators, forming a directional coupler with controllable gap distance d and, therefore, tunable coupling strength g(d). In the case of illumination from one waveguide, photons are fully transferred to the other waveguide after a length \( L_t(d) = \pi/(2g(d)) \). The splitting ratio at the end of the coupling region (length \( L_c \)) is \( SR = \tan^2((\pi L_c)/(2L_t(d))) \). By tuning the gap distance with electro-static force, photons can be deterministically switched between two output ports, as shown in Figs. 1(b) and 1(c). The effect of tuning is more pronounced with a smaller starting gap distance \( x_0 \), showing the advantage of working at the nanoscale.

We characterize our device in a Helium flow cryostat at 10 K. Multiple QDs in one arm of the waveguide (Port 1) are optically excited and the emission spectra, collected from Port 3 and Port 4 at different driving voltages, are
shown in Fig. 2(a). A clear anti-correlation between the two ports can be found. Unlike thermo-optical methods, the absence of de-tuning of the excitonic emission energy confirms that disturbance-free single-photon routing through electro-mechanical approach is achieved. The integrated and normalized photon counts of an emission line at a wavelength of 927.26 nm, shown in Fig. 2(b), indicate a maximum extinction ratio of $(23.4 \pm 1.7) \text{ dB}$ at a bias voltage of 7 V, corresponding to a displacement $x = 24 \text{ nm}$ (see ref. [2]). The single-photon nature of the routing is confirmed by a pulsed Hanbury-Brown and Twiss experiment, showing a single-photon purity $g^{(2)}(0) = 0.18 \pm 0.03$ (see ref. [2]).

The time response of the router is limited by the fundamental resonant frequency $\nu_m$, which, from numerical simulation (inset of Fig. 2(c)), is estimated at around 1.36 MHz, enabling sub-microsecond switching time. To verify this, we perform a ring-down measurement by applying a voltage of 4 V with a ramp-down time $\Delta t = 1 \mu s$. To avoid damaging the device due to excessive oscillation at resonance, the sample is tested at room temperature and at ambient pressure to dampen the mechanical quality factor. As shown in Fig. 2(c), the measured optical signal, which probes the mechanical motion of the waveguides, follows the dynamics of the input bias signal, indicating that the device can be driven at MHz speed. Moreover, resonant actuation schemes could be used to route photons from port to port at $1/(2\nu_m) \sim 370 \text{ ns}$, for applications such as on-chip de-multiplexing [3].

By comparing the loss of the router to suspended waveguides, we estimate a total insertion loss of 0.67 dB/switch. The demonstrated high performance of the single-photon router opens new opportunities for building de-multiplexed sources of $N$ single photons. Given the low insertion loss, a 10-photon source can be readily achieved by cascading the nanomechanical routers. With the use of QDs at telecom wavelengths and further realistic device optimization, the technology could be scaled up to achieve $N > 50$, required to show quantum advantage for boson sampling. Furthermore, adding the spin quantum memory could lead to a new generation of non-trivial photonic quantum resources.

In summary, we present a nano-electro-mechanical reconfigurable single-photon router integrated with QDs as deterministic sources. The demonstrated extinction ratio of $>20 \text{ dB}$, low insertion loss of 0.67 dB/switch and operation speed of MHz, combined with the small footprint and the planar design, allow cascading multiple devices in a single chip, opening new opportunities for complex on-chip quantum experiments such as de-multiplexing single-photon sources and programmable gates for boson sampling and quantum simulation.

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