The fate of the interstellar medium in early-type galaxies I. First direct measurement of the timescale of dust removal


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The fate of the interstellar medium in early-type galaxies

I. First direct measurement of the timescale of dust removal

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ABSTRACT

Context. An important aspect of quenching star formation is the removal of the cold interstellar medium (ISM; non-ionised gas and dust) from a galaxy. In addition, dust grains can be destroyed in a hot or turbulent medium. The adopted timescale of dust removal usually relies on uncertain theoretical estimates. It is tricky to track dust removal because the dust is constantly being replenished by consecutive generations of stars.

Aims. Our objective is to carry out an observational measurement of the timescale of dust removal.

Methods. We explored an approach to select galaxies that demonstrate detectable amounts of dust and cold ISM coupled with a low current dust production rate. Any decrease of the dust and gas content as a function of the age of such galaxies must, therefore, be attributed to processes governing ISM removal. We used a sample of the galaxies detected by Herschel in the far-infrared with visually assigned early-type morphology or spirals with red colours. We also obtained JCMT/SCUBA-2 observations for five of these galaxies.

Results. We discovered an exponential decline of the dust-to-stellar mass ratio with age, which we interpret as an evolutionary trend the ISM, as well as on the assumed dust grain properties and uncertain theoretical estimates. An efficient process involves supernova (SN) forward shocks sweeping up dusty ISM material and, thereby, completely destroying existing dust grains (e.g. Jones & Nuth 2011). Depending on the environment in the ISM, as well as on the assumed dust grain properties and

1. Introduction

An important aspect of quenching star formation is the removal of the cold interstellar medium (ISM; non-ionised gas and dust) from a galaxy. This can happen by outflows, astration, or ionisation of gas by stars or active galactic nuclei. In addition, dust grains can be destroyed in a hot or turbulent medium by sputtering and sublimation (Barlow 1978; Dwek et al. 1996; Silvia et al. 2010; Slavin et al. 2015). This process must be taken into account when modelling dust evolution in nearby and distant galaxies (e.g. Gall et al. 2011a; Lagos et al. 2014, 2015; Zhukovska et al. 2016; Krumholz et al. 2017). Dust is an important constituent of the cold ISM and can be used as its tracer (e.g. Scoville et al. 2014, 2016).

The adopted timescale of dust removal usually relies on uncertain theoretical estimates. An efficient process involves supernova (SN) forward shocks sweeping up dusty ISM material and, thereby, completely destroying existing dust grains (e.g. Jones & Nuth 2011). Depending on the environment in the ISM, as well as on the assumed dust grain properties and
destruction mechanisms, the dust destruction timescale may be as short as a few thousand years up to ~1 Gyr (Draine & Salpeter 1979; Jones et al. 1994; Jones 2004; Zhukovska et al. 2016).

Observationally, estimating dust destruction rates in galaxies is impacted by substantial uncertainties. Recent investigations of SN remnants in the Milky Way and the Magellanic Clouds suggest a short destruction timescale of a few tens of Myr (Temim et al. 2015; Lakićević et al. 2015). However, for such short timescales it would be very difficult to build up and maintain significant amounts of dust in any galaxy in the absence of a very efficient dust production mechanism which would constantly replenish the dust to counter its rapid destruction (e.g. Gall et al. 2011b; Gall & Hjorth 2018). On the other hand, Barlow (1978) and Slavin et al. (2015) show that dust destruction due to SN forward shocks, mainly by sputtering, should occur on a timescale of 1–3 Gyr (see Gall & Hjorth 2018 for a recent discussion). However, the complete destruction of dust grains through SN shocks may not be the only mechanism at play in the removal of dust from a galaxy. Heating by planetary nebulae, cosmic rays, X-rays, or other high energy radiation can also alter dust grains. In addition, dust may be blown out of galaxies due to the feedback of stars and SNe (e.g. Finkelstein et al. 2012; Conroy et al. 2015; Spilker et al. 2018; Jones et al. 2019; Li et al. 2019a).

The vast amounts of molecular gas found in post-starburst galaxies indicate that the process of removing cold ISM after the starburst (e.g. through outflows, AGN or stellar feedback) is not rapid (Rowlands et al. 2015; French et al. 2015; Alatalo et al. 2016; Smercina et al. 2018). Rowlands et al. (2015) analyse a sample of starburst and post-starburst galaxies at ages up to ~1 Gyr after the starburst and do not find any clear decrease of the molecular gas content with age, indicating that ISM removal must operate on longer timescales. They did find a weak decreasing trend of dust mass with age, consistent with dust removal during the first ~200 Myr after the starburst. The dust temperature exhibits a similar decrease. On the other hand, combining all post-starburst galaxies with molecular gas measurements, French et al. (2018) and Li et al. (2019b) detect a weak anti-correlation of the molecular gas fraction with age, deriving a gas removal timescale of ~100–200 Myr.

In contrast, the only observational estimate of the lifetimes of dust grains in the hot ISM of early-type galaxies (ETGs) turned out to be very short. \(<(45 \pm 25)\) Myr, based on dust production rates derived for asymptotic giant branch (AGB) stars from mid-infrared observations and non-detections by Herschel (Clements et al. 2010). The galaxies were selected based on their lack of any sign of recent star-formation in the mid-infrared spectra, so they most likely represent the oldest early-type galaxies. Moreover, these galaxies are located in the Virgo cluster, so this timescale may apply to cluster-specific processes.

The analysis of dust removal in galaxies is complicated by the fact that most of the studies on this subject have concentrated on star-forming spiral/irregular galaxies (Draine et al. 2007; Draine 2009; Temim et al. 2015; Lakićević et al. 2015) in which both the dust destruction rate is and the dust production rate are high, due to on-going star formation. Hence, it is tricky to track the dust removal, because usually the dust is constantly being replenished by consecutive generations of stars. Here we explore an alternative approach, that is, through the selection of those galaxies which demonstrate detectable amounts of dust and cold ISM while exhibiting a low current SN dust production rate because of the extensive amount of time elapsed since the last major star-formation event. This can be achieved by selecting dusty early-type galaxies (with little star formation).

Therefore, any decrease of dust and gas content as a function of the age of such galaxies must be attributed to processes governing the ISM removal. This selection obviously limits our ability to study dust destruction processes connected with SNe (as their rate is also low), but allows us to explore dust removal in galaxies on the way to becoming passive.

With large-area Herschel surveys, the dust emission of significant numbers of such (rare) dusty early-type galaxies have been detected (Bourne et al. 2012; Rowlands et al. 2012; Smith et al. 2012; Agius et al. 2013, 2015; Wilson et al. 2013; Dariush et al. 2016). Here we use the term ETG for galaxies morphologically classified as ellipticals, lenticulars, or early-type spirals (Sa and SbA). In this series of papers, we embark on a study of the ISM removal in such early-type galaxies.

The objective for this paper is to carry out an observational measure of the timescale of dust removal. In future papers of this series, we intend to constrain the physical mechanism of this process based on the gas properties of these galaxies and theoretical modelling.

We use a cosmological model with \(H_0 = 70\) km s\(^{-1}\) Mpc\(^{-1}\), \(\Omega_\Lambda = 0.7\), and \(\Omega_m = 0.3\). All calculations are carried out assuming the Chabrier (2003) initial mass function (IMF).

## 2. Sample of dusty early-type galaxies

We use the sample of dusty early-type galaxies of Rowlands et al. (2012). This sample includes all galaxies with elliptical morphology and red spirals in the Herschel Astrophysical Terahertz Large Area Survey (H-ATLAS; Eales et al. 2010) \(\sim 14\) deg\(^2\) Science Demonstration Field (Ibar et al. 2010; Pascale et al. 2011; Rigby et al. 2011; Smith et al. 2011) that are detected at 250\(\mu\)m (all of them are also detected at 350\(\mu\)m and most of them at 500\(\mu\)m). The specific selection criteria used by Rowlands et al. (2012) are:

1. Set within the H-ATLAS Science Demonstration Field (solely positional).
2. Matched to an optical Sloan Digital Sky Survey (SDSS) source with a spectroscopic redshift (of a redshift range 0.10 < \(z_{\text{spec}} < 0.32\)) within a 10\(^\prime\) radius and with >0.8 reliability (removes high-\(z\) galaxies).
3. Herschel 250\(\mu\)m > 5\(\sigma\) detection (selects dusty galaxies).
4. Visually classified as early-type (elliptical or S0), or red spiral with near-ultraviolet (NUV) to r-band colour of NUV–r > 4.5 (morphology and colour).

We note that, in principle, criterion 2 does not remove high-\(z\) dusty galaxies lensed by nearby ones (Negrello et al. 2010, 2017; Bussmann et al. 2013; Wardlow et al. 2013). However, the maximum 500\(\mu\)m flux of the sources in the sample of Rowlands et al. (2012) is 62 mJy, and 90% of them are below 30 mJy. At such fluxes the contribution of lensed sources is expected to be, at most, only a few percent (Negrello et al. 2010, 2017; Wardlow et al. 2013).

The sample consists of 61 galaxies, including 42 ellipticals or lenticulars, and 19 red spirals, most of which are classified as Sa or SbA galaxies. These galaxies will be referred to here as ETGs. Rowlands et al. (2012) fitted the ultraviolet-to-submm spectral energy distribution [from the Galaxy And Mass Assembly (GAMA) survey; Driver et al. 2011, 2016; Hill et al. 2011; Robotham et al. 2010; Baldry et al. 2010] using Multi-wavelength Analysis of Galaxy Physical Properties (MAGPHYS; da Cunha et al. 2008). MAGPHYS fits the photometry to account for the balance between energy absorbed by dust and re-emitted in the infrared with the dust that is distributed through two components: diffuse and cloud-like.
In particular, Rowlands et al. (2012) present MAGPHYS-based estimates of stellar and dust (with an opacity coefficient of $k_{250\mu m} = 0.89 \text{ m}^2 \text{ kg}^{-1}$) masses, star formation rates (SFRs), and luminosity-weighted ages of stellar populations, which are used here. Rowlands et al. (2012) confirm the passive nature of most of these galaxies by measuring low SFRs and specific SFRs (normalised by stellar mass), which are well below those of galaxies on the star-formation main sequence (also see: Hunt et al. 2014; Hjorth et al. 2014), and a long time having passed (>1 Gyr) since the latest star-formation episode.

The basic properties of the galaxies in our sample are shown in Figs. 1–3. These figures show which properties evolve with age and which are relatively constant. The Sersic indices in the $r$-band of our galaxies are measured by Kelvin et al. (2012) within the GAMA Data Release 3 (Baldry et al. 2018). We also show the star-formation main sequence on the relevant panels (Speagle et al. 2014) and mark galaxies that are located within the main sequence with empty symbols.

3. JCMT/SCUBA-2 observations

In order to assess the robustness of the Herschel-only dust mass estimates, we obtained longer wavelength data for some of these galaxies. These data would allow us to detect very cold dust, which, in principle, could dominate the total dust mass budget. We obtained 850 $\mu$m and 450 $\mu$m data using the James Clerk Maxwell Telescope (JCMT) equipped with the Submillimetre Common-User Bolometer Array 2 (SCUBA-2; Holland et al. 2013) for a subset of the galaxies with higher ages (>2 Gyr; proposals M15AI116 and M15BI010, PI: Michalowski). We utilised the standard “daisy” scanning pattern, suitable for compact sources. The data were retrieved from the Canadian Astronomy Data Centre (CADC) and were reduced with the standard pipeline in SMURF package (Chapin et al. 2013) with the flux calibration factor (FCF) of 537 Jy $\text{pW}^{-1}$ $\text{beam}^{-1}$ for 850 $\mu$m and 491 Jy $\text{pW}^{-1}$ $\text{beam}^{-1}$ for 450 $\mu$m (Dempsey et al. 2013). The full width at half maximum (FWHM) of the resulting 850 and 450 $\mu$m maps is 14.6 and 7.9 arcsec, respectively. The observing log is presented in Table 1.

For each galaxy observed by JCMT/SCUBA-2, we performed aperture photometry on the 850 $\mu$m and 450 $\mu$m maps using aperture sizes equal to the SDSS DR7 isophotal major axis, which was also used for Herschel aperture fluxes (Pascale et al. 2011; Rigby et al. 2011). The noise was assessed by placing apertures of the same size at random around the target position. We did not detect any of the targets. The measured fluxes are presented in Table 1.

The spectral energy distributions including Herschel data are shown in Fig. 4. We overplot the grey-body models using the cold-dust temperatures derived by Rowlands et al. (2012) and the emissivity index $\beta = 2$. In three out of five cases, our upper limits are very close to the extrapolations from shorter wavelengths, ruling out the existence of a submillimetre excess, which could be the manifestation of very cold dust.

In order to assess the amount of very cold dust allowed by the data, we add a grey-body model with a temperature of 10 K (below the coldest average dust temperature ever reported for a galaxy; e.g. Chapman et al. 2005; Michałowski et al. 2008; Symeonidis et al. 2013). If we assume a higher temperature, we would infer an even lower cold dust mass. We scale this model so that its sum with the model derived using Herschel data matches the SCUBA-2 upper limits. For two galaxies with shallower SCUBA-2 limits (J085946.7−000020 and J090038.0+012810) this causes the total model to overpredict the Herschel fluxes, so we scale the very cold model further down so that the total model matches the Herschel fluxes. This allowed us to measure the limits on the dust masses of this component. Their ratios to the Herschel-derived masses are shown in the last column of Table 1. In all cases, the very cold dust component has, at most, twice the mass of the warmer dust. Even if such amounts of very cold dust are indeed present in galaxies with high ages, this effect is not sufficiently strong to explain the $M_\text{dust}$ drop of 1.5 dex shown of Fig. 1.

4. A declining dust mass with age

4.1. Timescale of dust removal

In view of the lack of any indication of the existence of very cold dust in our sample (as judged by our SCUBA-2 sub-sample), we proceed under the assumption that the Herschel-derived dust masses are reliable.

We note that the dust mass on average appears to be lower at higher ages in Fig. 1. At the same time, SFR and specific SFR are also lower on average, whereas stellar mass is somewhat higher. We do not detect any significant correlation between the age and the dust temperature or environment (Fig. 1), dust attenuation, optical size (Rowlands et al. 2012), or Sersic index (Kelvin et al. 2012; Fig. 2), indicating that these properties are not driving dust evolution. We note that dust attenuation in molecular clouds ($\tau_{MC}$ in MAGPHYS) is poorly constrained, but in these galaxies molecular clouds contribute on average only ~25% of the infrared luminosity (Rowlands et al. 2012), so this parameter is not critical.

As indicated in Fig. 1, the stellar masses of the galaxies in our sample span three orders of magnitude, so we use the dust mass relative to the stellar mass as an indicator of the dust content of a galaxy. Hence, our analysis will be based on age trends of quantities normalised to the stellar mass. In this way, the evolutionary trends we describe are not affected by the fact that the galaxies in our sample are not equally massive. For example, for galaxies with similar dust masses, but different stellar masses, we conclude that the more massive galaxy is more dust-deficient. This is true only in a mass-normalised sense.

In Fig. 5, we plot the dust masses (normalised to the stellar mass of the galaxy) as a function of age. It is evident that the dust-to-stellar mass ratio (a measure of the “dustiness” of a galaxy) decreases with stellar age. The typical uncertainties of these properties (~0.3 dex for the dust-to-stellar mass ratio $M_\text{dust}/M_*$ and ~0.2 dex for ages; Rowlands et al. 2012) are much smaller than the range of the values. The Spearman rank correlation coefficient is $-0.7$ with a very small probability ($\sim 4 \times 10^{-11}$, $\sim 5.5\sigma$, or $\sim 5 \times 10^{-7}$, $\sim 5\sigma$ without galaxies on the main-sequence) of the null hypothesis (no correlation) being acceptable.

We now assume that this can be interpreted as an evolutionary-ary sequence (we discuss alternative explanations below). There is significant scatter in the relation, so we do not propose that individual galaxies can evolve into each other; rather, it should be understood in the sense that the progenitors of galaxies with higher ages drawn from our sample were broadly similar to those with lower ages: older galaxies have less dust because they have had more time to get rid of it.

We can determine the timescale of the apparent dust decline by fitting an exponential function

$$M_\text{dust}/M_* = A e^{-\text{age}/\tau},$$

where $\tau$ is the timescale of dust removal. We note that this can be interpreted as an evolutionary sequence when $\tau$ is a constant. The best-fitting value of $\tau$ is 2.5 Gyr with a 95% confidence interval of (1.5, 3.5) Gyr.
### Fig. 1. Stellar mass (top left), SFR (top right), dust mass (middle left), specific SFR (middle right), temperature of cold diffuse ISM dust component (bottom left) and galaxy number surface density (bottom right) derived by Rowlands et al. (2012) as function of luminosity-weighted stellar age. The typical errors are shown as large crosses. On the SFR panel, the dotted line shows the exponential fit to the data, whereas on the $M_{\text{dust}}$ panel, the dotted line denotes the evolution of dust mass affected only by astration resulting from this fit and the assumption of the gas-to-dust mass ratio of 100. A higher ratio would result in an even flatter evolution. The specific SFR panel shows the star formation main sequence at $z = 0.13$ for stellar masses of $\log(M_*/M_\odot) = 10$ (higher line) and 11.5 (lower line) as solid lines, with the scatter shown as dotted lines (Speagle et al. 2014). The galaxy number surface density is defined as $\Sigma_5 = 5/\pi d_5^2$, where $d_5$ is the projected comoving distance to the 5th nearest neighbour within $\pm 1000\,\text{km}\,\text{s}^{-1}$. Open circles denote galaxies which are within the main sequence (see Fig. 3). Filled symbols which appear to be within the main sequence are galaxies at higher redshifts, for which the main-sequence is higher.

Where $A$ is the normalisation constant and $\tau$ is the lifetime. We obtain a lifetime of $\tau = (2.5 \pm 0.4)\,\text{Gyr}$ and a half-life time $t_{1/2} = \tau \ln(2) = (1.75 \pm 0.25)\,\text{Gyr}$. If we include only galaxies below the main sequence we get consistent results: $\tau = (2.7 \pm 0.5)\,\text{Gyr}$, $t_{1/2} = (1.87 \pm 0.35)\,\text{Gyr}$. This is the first time the dust removal timescale has been measured directly.

Interpreting the dust decline in Fig. 5 as an evolutionary sequence, the fitting of Eq. (1) gives an estimate of the dust removal timescale in early-type galaxies. This quantity may be
used in models in which it needs to be assumed a priori and cannot be derived. Such a long timescale (∼2 Gyr) indicates that dust is present even in very old galaxies, as long as sufficient amounts have been formed in the past (see Sect. 5 for a discussion on the source of dust in these galaxies).

Our estimate of the dust removal timescale is consistent with theoretical predictions of Barlow (1978) and Slavin et al. (2015). On the other hand, it is longer than the one measured by Clemens et al. (2010, <45 ± 25 Myr). This is not surprising as their measurement is based on the current AGB dust production rates, which leads to an underestimation of the lifetime if another dust production mechanism is dominant (e.g. grain growth in the ISM). Our derived lifetime of dust grains is much longer than that which is expected for a hot medium (Draine & Salpeter 1979; Jones et al. 1994; Jones 2004). Therefore, either there is not much hot gas in these galaxies, or dust is efficiently shielded in dense pockets. For a given stellar mass of an ETG there is spread in the X-ray luminosity of the hot gas of more than 2 orders of magnitude (Boroson et al. 2011; Su et al. 2015) so it is not possible to predict whether galaxies in our sample contain such gas.

It is helpful to compare the dust removal timescale we derive with the quenching timescale of shutting down star formation. Peng et al. (2015) and Trussler et al. (2018) measure a timescale for quenching by strangulation, a process involving halting gas inflows onto galaxies. The timescale is 2–5 Gyr for galaxies with masses $\log M_*/M_\odot = 10–11$ independent of environment (cluster or field galaxies). Similarly, simulated galaxies in the same mass range exhibit a quenching timescale of 2–4 Gyr (Wright et al. 2019). The dust removal lifetime we derive is similar or up to a factor of two shorter than these quenching timescales. This suggests that the shutdown of star formation and the decline in dust content may be connected. It is, however, unlikely that the mechanism of the dust decline is directly related to star formation, for example, by astration or SN dust grain destruction. Fig. 1 shows that the SFR is declining with age, so the effect of such processes is the weakest at high ages. Hence, if these processes were responsible for the dust decline, then the decline would be flattening at high ages, whereas Fig. 5 demonstrates a steepening with age. This effect is demonstrated in Fig. 1. We fit an exponential function to the SFR vs. age panel and calculate the resulting removal of dust, assuming a
is morphological quenching, which means it would make the gas resilient against fragmentation by the influence of the bulge (Martig et al. 2009, 2013; Bitsakis et al. 2019; Lin et al. 2019). This evolution away from the $M_{\text{dust}}$–SFR relation is also evident in Fig. 7, which shows the distance from this relation as a function of age. The galaxies in our sample, as they leave the star-formation main-sequence, are also leaving the $M_{\text{dust}}$–SFR relation.

Finally, we note that the dust removal timescale we report here is measured using a sample of early-type and red spiral galaxies with detected dust emission, so there is no evidence that it applies to different galaxies. Star-forming galaxies are likely to have quicker dust destruction processes, due to, for example, frequent SN explosions and more turbulent ISM.

The ratio of the dust mass to stellar (or baryonic) mass has been used as an indicator of the evolutionary stage of galaxies. The dust decline we detect is qualitatively consistent with models of dust evolution by Clark et al. (2015), Calura et al. (2017), and De Vis et al. (2017a,b, 2019). They conclude that the dust-to-baryon (or stellar) ratio first increases for very young galaxies and then declines in a similar way as is measured here. The evolution in some of these models is traced by the atomic gas fraction (from high to low) and indeed early-type galaxies with such measurements in these works have low gas mass fractions. On the other hand, the models of Calura et al. (2017) are difficult to compare with our study because they evolve the proto-spheroid version to less than 1 Gyr, which is not considered here.

Smercina et al. (2018) interpret the dust decline in post-starburst systems as the result of sputtering due to the ISM being dominated by hot gas. The galaxies in our sample are older but it is possible that this effect is continuing to be important for them as well.

4.2. Potential biases

The dust decline in Fig. 5 relies on the accuracy of the MAGPHYS age measurement. Such measurements based on spectral energy distributions (SEDs) are, indeed, shown to be reliable in assigning a galaxy to an age bin at least up to the age of 7 Gyr (Yi 2003), which is the oldest age of the galaxies in our sample. Moreover, the age estimate we use correlates with the D4000 spectral break strength (Balogh et al. 1999) from single fiber spectroscopy in the GAMA survey (Fig. 2; Gordon et al. 20173), which is an indicator of the age of the stellar population independent of the photometric MAGPHYS estimate. This correlation has a Spearman rank correlation coefficient of 0.6 with a very small probability of $\sim 3 \times 10^{-7}$ (4$\sigma$) of the null hypothesis (no correlation) being acceptable.

conservatively low gas-to-dust mass ratio of 100. The evolution is too flat to explain the data and it is, indeed, flattening at high ages. A higher gas-to-dust mass ratio would result in an even flatter evolution.

In Fig. 6, we show how the dust decline we measure is related to the SFR decline (the manifestation of quenching). On the $M_{\text{dust}}$–SFR diagram, the dusty early-type galaxies are evolving away from the relation measured for star-forming galaxies (da Cunha et al. 2010), diagonally towards the bottom-left\(^2\). Such behaviour has been predicted by Hjorth et al. (2014) to be a manifestation of a quenching process which does not require removing the ISM (otherwise a galaxy would evolve parallel to the $M_{\text{dust}}$–SFR relation). One possible quenching mechanism

\(^2\) See Lianou et al. (2016) for similar location of dusty ETGs but without the age information.
It is unlikely that the apparent evolution of $M_{\text{dust}}/M_*$ with age is caused by the selection effect. The sample selection was based on Herschel detections (dust-mass limit) and optical morphology. Therefore, if galaxies in the top-right corner of Fig. 5 existed (old but dusty) they would be detected and included in the sample. Indeed, the dust mass detection limit for the H-ATLAS survey at the redshifts of our galaxies is $\sim10^{5.2} M_\odot$ at $z = 0.05$ and $\sim10^{6.7} M_\odot$ at $z = 0.3^4$, much lower than the lower envelope of the measurements for our sample (see Fig. 1). Similarly, galaxies in the bottom-left corner, below the trend, would have detectable dust masses as they would be well above the detection limit.

Moreover, the trend in Fig. 5 is not driven by the different stellar masses or morphologies of galaxies in our sample. In Fig. 8, we show this trend using different colours for three stellar mass bins, or for four Sersic index bins, along with different symbols for ellipticals and red spirals. It is evident that the trend persists even in narrower mass or Sersic index bins and restricting them to a single morphological type. In fact, the trend exhibits smaller scatter keeping only galaxies with Sersic indices larger than 4, apart from two outliers. Moreover, the Sersic index in our sample is not correlated with age (Fig. 2). The trend in Fig. 5 is not a result of the age-attenuation degeneracy in the SED fitting because there is no trend in the derived ages and attenuation (Fig. 2). If this were a problem then galaxies with high attenuation would be erroneously assigned high ages. However, these galaxies are the least dusty in our sample and are not more compact than the other members in the sample (see Fig. 2), so the stellar populations are unlikely to suffer the highest attenuation. Additionally, consistent to the expectation, the stellar attenuation in the MAGPHYS modelling is weakly correlated to the total dust masses in these galaxies (Fig. 9). The Spearman rank correlation coefficient is 0.36 with a probability of $\sim0.004$ ($\sim3\sigma$) of the null-hypothesis that there is no correlation.

Finally, our sample is unlikely to be affected by the age-metallicity degeneracy because most of these galaxies span the stellar gas range of $\log(M_{*}/M_\odot) = 10.5-11.5$, at which the mass-metallicity relation is flat. Hence, the expected metallicities

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**Table 1. Observation log and results for JCMT/SCUBA-2 observations.**

<table>
<thead>
<tr>
<th>Galaxy</th>
<th>Obs. date</th>
<th>$t_{\text{int}}$ (h)</th>
<th>$\tau_{\text{CSO,225GHz}}$</th>
<th>$F_{\text{850}}$ (mJy)</th>
<th>$F_{\text{450}}$ (mJy)</th>
<th>$T_{\text{dust}}$ (K)</th>
<th>$M_{\text{cold}}/M_{\text{Herschel}}$</th>
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<td>2015 Apr. 27, May 19, 2016 Jan. 07</td>
<td>1.5</td>
<td>0.07–0.075, 0.05, 0.03</td>
<td>0.7 ± 2.0</td>
<td>–11 ± 40</td>
<td>22.2</td>
<td>0.4</td>
</tr>
<tr>
<td>J091448.7–003533</td>
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<td>1.0</td>
<td>0.023–0.051</td>
<td>–3.1 ± 4</td>
<td>47 ± 30</td>
<td>21.5</td>
<td>0.8</td>
</tr>
<tr>
<td>J085946.7–000020</td>
<td>2015 Dec. 17, 2016 Jan. 05</td>
<td>1.0</td>
<td>0.04, 0.03</td>
<td>0.2 ± 7</td>
<td>–29 ± 100</td>
<td>19.5</td>
<td>1.8</td>
</tr>
<tr>
<td>J085828.5+003814</td>
<td>2015 Apr. 21</td>
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<td>0.044–0.047</td>
<td>–1.7 ± 3.0</td>
<td>21 ± 30</td>
<td>21.7</td>
<td>0.04</td>
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<tr>
<td>J090038.0+012810</td>
<td>2016 Jan. 05, 07</td>
<td>1.0</td>
<td>0.03</td>
<td>0.2 ± 5</td>
<td>68 ± 60</td>
<td>17.5</td>
<td>0.4</td>
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</tbody>
</table>

**Notes.** The columns are: galaxy name, observation date, on-source integration time, 225 GHz opacity measured by Caltech Submillimeter Observatory (CSO), measured flux at 850 and 450 µm, ISM dust temperature derived from the Herschel data (Rowlands et al. 2012), and the ratio of the maximal very cold dust with the temperature of 10 K allowed by the data to the dust mass measured with MAGPHYS.

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**Fig. 4.** Spectral energy distribution of sub-sample observed by JCMT/SCUBA-2. Circles and arrows denote detections and 3σ upper limits, respectively. The solid lines are the grey-body models using the diffuse ISM temperature derived by Rowlands et al. (2012) and the emissivity index $\beta = 2$. In three cases, the 850 µm data are deep enough to rule out contribution of very cold dust. Dotted lines denote the models of very cold dust with the temperature of 10 K, whereas the dashed lines denote the sum of both solid and dotted curves. The very cold dust model was scaled so that the total emission matches the SCUBA-2 upper limits, or Herschel data, if they are more constraining. The panels are ordered from left to right according to their increasing ages and these galaxies are marked in Fig. 5.

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of our galaxies are within a narrow range of $12 + \log(\text{O/H}) = 9.05–9.13$ (Tremonti et al. 2004).

5. The source of dust in early-type galaxies

The origin of dust in galaxies in general, and in ellipticals in particular, is a matter of current debate. While we can interpret the trend in Fig. 5 as dust removal in action, we do not possess direct information about past dust evolution.

In general, the source of dust can be internal or external. Internal dust formation mechanisms include AGB stars and SNe, with possible significant dust mass accumulation in the ISM. This means that the detected dust is what has remained from the pre-existing (before quenching) dust amount, or it is being formed by a current population of AGB stars. Alternatively, dust might have been recently brought in from an external source following a merging with dust-rich galaxies.

We now demonstrate that the contribution of AGB stars to dust production is minor (see also Rowlands et al. 2012) based on similar arguments as those presented in Michalowski et al. (2010a,b), Michalowski (2015), and Leśniowska & Michalowski (2019) for high-redshift galaxies (see also Morgan & Edmunds 2003; Dwek et al. 2007; Dwek & Cherchneff 2011; Rowlands et al. 2014). Stellar masses of $10^{11} M_\odot$ imply around $10^{10}$ AGB stars ($1.5–8 M_\odot$), assuming the Chabrier (2003) IMF. In order to explain the measured dust masses of $\sim 10^{7} M_\odot$ for galaxies in our sample at the lowest ages, each of the AGB stars would need to produce $\sim 10^{-2} M_\odot$ of dust. This is the highest AGB dust yields to have been found theoretically (Morgan & Edmunds 2003; Ferrarotti & Gail 2006; Ventura et al. 2012; Nanni et al. 2013, 2014; Schneider et al. 2014), which is valid only for narrow mass and metallicity ranges. On average, AGB stars produce less dust than this maximum by at least a factor of ten (Gall et al. 2011c). Similarly, Rowlands et al. (2012) also conclude that the detected dust masses in these galaxies cannot be explained by AGB dust production. Hence, AGB stars could not have dominated the dust production in these galaxies.

These dust masses are larger by one order of magnitude than the prediction of recent SN dust production proposed by Gall & Hjorth (2018) for star-forming galaxies. This is expressed as $M_{\text{dust}}^{\text{SN}} = \mu_\odot \text{SFR} \Delta t$, where $M_{\text{dust}}^{\text{SN}}$ is the total dust mass formed by SNe in the current star-formation episode, $\mu_\odot$ (equal to 0.004 for the Chabrier IMF) is a product of the dust mass formed by one SN and the rate of SNe per unit SFR. $\Delta t$ is the duration of the star-formation episode, for which we conservatively assume a large value of 1 Gyr. Indeed, it is unlikely that the dust detected for the galaxies in our sample was formed recently as the current SFRs are too low. The observed dust must have formed in the past, when SFRs were much higher. In a galaxy with a stellar mass of $10^{11} M_\odot$ with the Chabrier (2003) IMF, there are $\sim 1.2 \times 10^{9}$ stars with masses $8–40 M_\odot$, which can explode as SNe. Therefore, one SN would need to produce $\sim 0.8 M_\odot$ of dust to explain the measured dust masses. This is close to the maximum SN dust yield with no dust destruction (Todini & Ferrara 2001; Nozawa et al. 2003; Matsuura et al. 2011; Gall et al. 2014; Hjorth et al. 2014; Gall & Hjorth 2018). Moreover, some stars with very high masses
ended up collapsing into a black hole without producing dust at all. Hence, SNe could be considered the source of dust if they were maximally efficient and did not destroy any dust.

ISM grain growth or very efficient SN dust production has also been invoked for this sample (Rowlands et al. 2012) and for distant galaxies (Dwek et al. 2007; Michałowski et al. 2010a,b; Michałowski 2015; Gall et al. 2011a; Hirashita & Kuo 2011; Hjorth et al. 2014; Rowlands et al. 2014; Valiante et al. 2014; Leśniewska & Michałowski 2019), some of which may be representative progenitors of early-type galaxies in our sample. Hirashita et al. (2015) shows that grain growth in the ISM of ellipticals can operate on relatively short timescales of tens of Myr. That is fast enough to produce the considerable amounts of dust detected in galaxies in our sample shortly after they were quenched. On the other hand, Ferrara et al. (2016) shows that dust grain growth is too slow in the diffuse ISM, whereas in molecular clouds, this process has no influence on the total dust mass, because accreted material forms icy mantles which do not survive when a grain leaves the cloud. Hence, the issue of grain growth is not yet understood.

Alternatively, dust in our galaxies could be of external origin, brought on by mergers with gas-rich galaxies, but the properties of the galaxies in our sample are inconsistent with this scenario. It predicts that $M_{\text{dust}}$ should be only weakly correlated with $M_*$, because the gas-rich companions would not bring in a significant amount of stars, so the final $M_*$ would depend quite weakly on the number of merger events. However, Fig. 3 shows that $M_{\text{dust}}$ is correlated with $M_*$. The Spearman rank correlation coefficient is 0.5, with a very small probability ($\sim 3 \times 10^{-5}$, $\sim 4.2\sigma$, or $\sim 5 \times 10^{-5}$, $\sim 4.1\sigma$ without the galaxies on the main sequence) of the null hypothesis (no correlation) being acceptable. This suggests that current dust reservoirs have been formed by the stellar populations present in these galaxies. In contrast di Serego Alighieri et al. (2013) finds no correlation between these quantities for dusty ellipticals in the Virgo cluster (see also Davis et al. 2011; Smith et al. 2012; Kaviraj et al. 2012, 2013; Dariush et al. 2016; Beeston et al. 2018; Kokusho et al. 2019), leading to the conclusion that the dust has been brought in by gas-rich mergers.

Moreover, in the external origin scenario, $M_*$ should be independent of or slightly decreasing in proportion to the derived age because galaxies with lower inferred age (and higher dust masses brought by mergers) have, on average, experienced more
merging and, thus, they should be more massive. However, Fig. 1 shows that the stellar mass is increasing with age, which should not be the case if the apparent dust decline with age was a result of a merging with smaller galaxies at high derived ages and with larger galaxies at low ages.

Over several billion years, the size of a galaxy is also expected to grow as a result of minor mergers (Naab et al. 2009; Hopkins et al. 2010; Trujillo et al. 2011; Furlong et al. 2017). This is not observed for our sample (Fig. 2).

Finally, the incidence of high dust masses of $\sim 10^8 M_\odot$ at lower stellar ages makes it unlikely that most of the dust has been brought in by minor mergers. Indeed, Rowlands et al. (2012) found unfeasibly high number of mergers compared to simulations of what is needed to supply enough dust. Additionally, they did not find morphological disturbances for these galaxies, which would be expected in cases of mergers.

We also note that Griffith et al. (2019) present evidence of internal production of dust for three other ellipticals by finding that the molecular gas mass and the stellar mass of these galaxies, which would not be the case if dust was brought by a merger with a low-metallicity dwarf. Moreover Babiy et al. (2018) find that the molecular content of ellipticals is correlated with their hot gas reservoir, which also suggests an internal origin of the molecular gas. Finally, Bassett et al. (2017) find the kinematics of dusty early-type galaxies consistent with recent minor mergers for only 8% of their sample. Similarly, in a small sample of early-type galaxies, Sansom et al. (2019) find carbon monoxide (CO) emission line rotational axes aligned with photometric axes, which is inconsistent with a recent merger/inflow hypothesis.

We conclude that it is likely that the dust in galaxies in our sample is of internal origin. It was either formed in the past directly by SNe (when SFRs were high) or with seeds produced by SNe, and with grain growth in the ISM significantly contributing to the dust mass accumulation shortly after the galaxies were quenched.

6. Conclusions

We have analysed the dust properties of a sample of galaxies detected by Herschel in the far-infrared based on the early-type morphology or spirals with red colours identified visually. We have found no obvious biases which could affect our analysis. We have discovered an exponential decline of the dust-to-stellar mass ratio with age, which we interpret as an evolutionary trend of dust removal from these galaxies. For the first time, we have directly measured the dust removal timescale in such galaxies to be $\tau = (2.5 \pm 0.4)\text{Gyr}$ (the corresponding half-life time is $\tau = (1.75 \pm 0.25)\text{Gyr}$). This quantity may be used in models in which it needs to be assumed a priori and cannot be derived. Any process which removes dust in these galaxies cannot, therefore, take place on shorter timescales.

This timescale is comparable to the quenching timescales found in simulations for galaxies with similar stellar masses. This suggests that the shutdown of star formation and the decline in dust content may be connected.

We find that the dust in these galaxies is likely of internal, not external origin. It was either formed in the past directly by SNe or with seeds produced by SNe, and with grain growth in the ISM significantly contributing to the dust mass accumulation.

In future papers we intend to constrain the physical mechanism of this process based on the gas properties of these galaxies and theoretical modelling.

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