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An improved land biosphere module for use in the DCESS Earth system model (version 1.1) with application to the last glacial termination

Roland Eichinger1, Gary Shaffer2,3, Nelson Albarrán1, Maisa Rojas1, and Fabrice Lambert5

1Department of Geophysics, University of Chile, Blanco Encalada 2002, Santiago, Chile
2GAIA-Antarctica, University of Magellanes, Avenida Bulnes 01855, Punta Arenas, Chile
3Niels Bohr Institute, University of Copenhagen, Blegdamsvej 17, Copenhagen, Denmark
4Department of Physics, University of Santiago de Chile, Avenida Ecuador 3493, Santiago, Chile
5Department of Physical Geography, Catholic University of Chile, Vicuña Mackenna 4860, Santiago, Chile

Correspondence to: Roland Eichinger (roland.eichinger@dlr.de, roland@dgf.uchile.cl)

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Abstract. Interactions between the land biosphere and the atmosphere play an important role for the Earth’s carbon cycle and thus should be considered in studies of global carbon cycling and climate. Simple approaches are a useful first step in this direction but may not be applicable for certain climatic conditions. To improve the ability of the reduced-complexity Danish Center for Earth System Science (DCESS) Earth system model DCESS to address cold climate conditions, we reformulated the model’s land biosphere module by extending it to include three dynamically varying vegetation zones as well as a permafrost component. The vegetation zones are formulated by emulating the behaviour of a complex land biosphere model. We show that with the new module, the size and timing of carbon exchanges between atmosphere and land are represented more realistically in cooling and warming experiments. In particular, we use the new module to address carbon cycling and climate change across the last glacial transition. Within the constraints provided by various proxy data records, we tune the DCESS model to a Last Glacial Maximum state and then conduct transient sensitivity experiments across the transition under the application of explicit transition functions for high-latitude ocean exchange, atmospheric dust, and the land ice sheet extent. We compare simulated time evolutions of global mean temperature, \( pCO_2 \), atmospheric and oceanic carbon isotopes as well as ocean dissolved oxygen concentrations with proxy data records. In this way we estimate the importance of different processes across the transition with emphasis on the role of land biosphere variations and show that carbon outgassing from permafrost and uptake of carbon by the land biosphere broadly compensate for each other during the temperature rise of the early last deglaciation.

1 Introduction

On centennial to millennial timescales, ocean processes may largely determine variations of atmospheric \( CO_2 \) concentrations (Fischer et al., 2010; Sigman et al., 2010). Such processes include changes in ocean dynamics as well as in biogeochemical properties like variations in the phosphate inventory or iron fertilization (Martin et al., 1990; Maher et al., 2010). However, interactions between atmosphere and land can also have an important impact on the overall change in the carbon cycle and thus on the Earth’s climate system. Net primary production on land takes up \( CO_2 \) from the atmosphere at a rate that increases with the \( pCO_2 \) itself (CO2 fertilization; Saugier et al., 2001). Remineralization in the soils increases with increasing temperature (Davidson and Janssens, 2006). Different vegetation zones advance and retreat due to varying climate conditions, thereby changing the terrestrial biomass budget and thus the carbon amount being stored in vegetation (Ciais et al., 2012). Moreover, changes in permafrost area and, during glacial conditions, changes in
areas covered by ice sheets also have the potential to modify atmospheric $\rho$CO$_2$ significantly (Schuur et al., 2008). The release of carbon into the atmosphere through the thawing of permafrost in a warming future climate has been assessed in a number of studies (e.g. Schaefer et al., 2011; Schuur et al., 2008; Khvorostyanov et al., 2008) and carbon storage and release in and from permafrost can also help explain glacial–interglacial cycles (Zech, 2012; Ciais et al., 2012; Crichton et al., 2016). A land biosphere module within an Earth system model should be able to address these processes.

For this reason, we here extend the Danish Center for Earth System Sciences (DCESS) Earth system model (Shaffer et al., 2008) by a new terrestrial biosphere scheme. This parameterization features the three vegetation zones – tropical forests (TF); grasslands, savanna and deserts (GSD); and extratropical forests (EF) – through definition of their characteristic values of biomass reservoirs and net primary production (NPP). The dynamic accounting of the latitudinal boundaries of the different zones and thereby their area extents is approximated by fitting polynomial functions of global mean temperature ($T_{\text{glob}}$) to results of a complex vegetation model study by Gerber et al. (2004). For completeness we also developed a simple approach to vegetation albedo based on the relative sizes of the three vegetation zones. Moreover, we present a component that accounts for carbon being stored in permafrost and below terrestrial ice sheets to allow extensive carbon storage on land during glacial climate conditions and its release across deglaciation events. In DCESS model simulations, these new developments considerably improve the estimates of amount and timing of land–atmosphere carbon exchanges, including the carbon isotopes $^{13}$C and $^{14}$C.

For a first application of this new module, we furthermore developed a set of explicit functions that describe the transitions of high-latitude ocean exchange, atmospheric dust and land ice sheet extent within the last 25 kyr BP. This allows us to simultaneously simulate time series of global mean temperature, $\rho$CO$_2$, atmospheric and oceanic carbon isotopes as well as ocean dissolved oxygen concentrations across the deglaciation after the Last Glacial Maximum (LGM, $\sim$21,000 years ago). Hitherto, the DCESS model has been used mainly for future climate projections (see e.g. Shaffer et al., 2009; Shaffer, 2010) and evaluated for pre-industrial (PI) climate conditions (see Shaffer et al., 2008). For the present application, the model is calibrated to glacial conditions by adapting physical and biogeochemical parameters guided by proxy data records. This includes a physically simple method to generate isolated deep water in the high-latitude model ocean (as had been hypothesized by several studies, e.g. Francois et al., 1997; Sigman and Boyle, 2000; Broecker and Barker, 2007) through the imposition of a depth profile for the vertical exchange intensity. Transient sensitivity simulations across the last 25 kyr BP are then performed. These demonstrate the impact and timing of various processes on atmospheric temperatures, $\rho$CO$_2$ and the carbon isotopes $^{13}$C and $^{14}$C at the beginning of the last glacial termination (“mystery interval” – MI, from 17.5 to 14.5 ka BP; Broecker and Barker, 2007).

2 A new land biosphere in the DCESS model

The DCESS model features components for the atmosphere, ocean, ocean sediment, land biosphere and lithosphere and has been designed for global climate change simulations on timescales from years to millions of years (Shaffer et al., 2008). Its geometry consists of one hemisphere, divided into two 360° wide zones by 52° latitude. The model ocean is divided into a low/mid-latitude and a high-latitude sector (as in the HILDA – high-latitude exchange/interior diffusion advection – model, developed by Shaffer and Sarmiento, 1995) and features a continuous vertical resolution of 100 m, to a depth of 5500 m. The near-surface atmospheric mean temperature is described by a simple, zonal mean, energy balance model in combination with sea ice and snow parameterizations. The atmosphere is assumed to be well mixed for gases and air–sea gas exchange fluxes, and transports via weathering, volcanism and interactions with the land biosphere are considered for carbon dioxide (CO$_2$) and methane (CH$_4$) in $^{12}$, $^{13}$, $^{14}$C species, respectively, as well as for nitrous oxide (N$_2$O) and oxygen (O$_2$). Ocean dynamics are characterized by high-latitude sinking and low/mid-latitude upwelling as well as horizontal and vertical diffusion between the latitude zones and the ocean layers. For the ocean biogeochemical cycling, a number of tracers are considered (namely, phosphate – PO$_4$; dissolved oxygen – O$_2$; dissolved inorganic carbon – DI $^{12}$, $^{13}$, $^{14}$C; alkalinity – ALK), which are forced by new production, air–sea exchange, remineralization of organic matter, dissolution of CaCO$_3$, river inputs and evaporation/precipitation (Shaffer, 1996; Shaffer et al., 2008). There is a sediment section for each of the ocean model layers addressing CaCO$_3$ dissolution/burial and organic matter remineralization/burial.

A land biosphere scheme accounts for the $^{12}$, $^{13}$, $^{14}$C cycling with leaf, wood, litter and soil boxes (Shaffer et al., 2008). NPP on land takes up CO$_2$ from the atmosphere and is distributed between leaves and wood. Leaf loss goes to litter, wood loss is divided between litter and soil, and litter loss is divided between the atmosphere (as CO$_2$) and the soil. Soil loss goes to the atmosphere as CO$_2$ and CH$_4$. Losses from all land reservoirs are taken to be proportional to reservoir size and, for litter and soil, to depend upon the mean atmospheric temperature according to $\lambda_Q \equiv Q^{(T_{\text{glob}}-T_{\text{glob},0})/10}$, where $Q_{10}$ (a biotic activity increase for a 10°C increase of $T_{\text{glob}}$) is chosen to be 2 (Friedlingstein et al., 2006).

In an attempt to remedy these deficiencies while retaining simplicity on the level of the rest of the model, we here present the extension of this scheme to three different vegetation zones. We define a latitudinal distinction of these three vegetation zones and their latitudinal boundaries on a global scale. The zones we consider are tropical forests (TF);
2.1 Description of the vegetation zones

The three vegetation zones (TF, GSD, EF) were defined on the basis of a study by Gerber et al. (2004). In that study, the complex LPJ terrestrial biosphere model (Lund–Potsdam–Jena Dynamic Global Vegetation Model) was applied to distinguish between a number of vegetation zones based on several variables. The latitudinal limits of these vegetation zones are dynamically defined. In general, the extent of certain vegetation zones depends mainly on temperatures and precipitation. However, the limitations of the DCESS model (no explicit computation of precipitation and restriction to two latitudinal sections) require a somewhat more general approach. We therefore determine the division of the three vegetation zones solely by the deviation of the global mean atmosphere temperature from its PI value (15°C). For this purpose, we derived two polynomial functions from a study by Gerber et al. (2004). We started from the total tree cover frame of their Fig. 4 by reading off, at 2°C intervals from −10 to 10°C deviation from pre-industrial global mean temperature, the latitudes in the Northern Hemisphere of 50% tree cover both above and below the subtropical zone of lower tree cover. Each of these two sets of 11 points formed the basis of our curve fitting. We found that fifth-order polynomials provided good fits to each of these sets. This emulation of a complex vegetation model thereby implicitly includes the role of precipitation in the temperature dependence of the vegetation zone boundaries. The two latitudinal limitations of the vegetation zones are described by the two fifth-order polynomials

\[
L_{\text{TF-GSD}} = -1.83 \times 10^{-5} \cdot \delta T_{\text{glob}}^5 - 0.0005809 \cdot \delta T_{\text{glob}}^4 + 0.005168 \cdot \delta T_{\text{glob}}^3 + 0.0497 \cdot \delta T_{\text{glob}}^2 + 1.092 \cdot \delta T_{\text{glob}} + 11.28
\]

and

\[
L_{\text{GSD-EF}} = 1.152 \times 10^{-5} \cdot \delta T_{\text{glob}}^5 - 0.0001785 \cdot \delta T_{\text{glob}}^4 - 0.004557 \cdot \delta T_{\text{glob}}^3 + 0.04156 \cdot \delta T_{\text{glob}}^2 + 1.017 \cdot \delta T_{\text{glob}} + 37.77,
\]

which depend only on the deviation of the global mean atmosphere temperature \(\delta T_{\text{glob}}\) from the calibrated PI steady state.

\[L_{\text{TF-GSD}}\] denotes the latitude of the border between the TF and the GSD zones and \(L_{\text{GSD-EF}}\) the latitude between GSD and EF. These two fifth-order polynomials are illustrated in Fig. 1.

The EF vegetation zone additionally is limited by either the model snow line or the line of the terrestrial ice sheet extent, depending on which one of the two lines expands the farthest from the pole at the current time step (see Sect. 2.4 for definition of “snow line” and further explanations). The snow line is also included in Fig. 1 – the zone poleward of the snow line is taken to be permafrost area in our simplified approach. Based on these latitudinal limits, the total \(\text{CO}_2\) and \(\text{CH}_4\) fluxes between the terrestrial biosphere and the atmosphere are now determined by the sum of the three vegetation zones, and thereby depend on the areas and mean temperatures of each zone as well as their values of NPP and stored biomass.

Table 1 shows the characteristic global values of biomass reservoirs and NPP of those vegetation zones at PI climate conditions \((T_{\text{glob}} = 15\,\text{°C}, \rho \text{CO}_2 = 280\,\text{ppm})\) (Gower et al., 1999; Saugier et al., 2001; Sterner and Elser, 2002; Zheng et al., 2003; Chapin III et al., 2011). The values in Table 1 have been constrained such that the sum over the three vegetation zones adds up to global PI values of the original biosphere model (Shaffer et al., 2008).

2.2 Vegetation albedo

For completeness and consistency, we also extended the model albedo calculation to account for the new biosphere
scheme with the three vegetation zones. In the DCESS model, albedo, $\alpha$, is taken to be constant and equal to 0.62 for all snow- or ice-covered areas. For non-snow/ice-covered areas, $\alpha$ is expressed as

$$
\alpha = a + b \cdot \left[ 0.5 \cdot (3 \cdot \sin(\Theta))^2 - 1 \right],
$$

(3)

where $\Theta$ is the latitude, $b = 0.175$ and $a = 0.3$ for present-day conditions. This functional form and these constant values have been based on present-day observations (Hartmann, 1994). The albedo of non-snow/ice-covered areas should vary with vegetation type since forested areas have lower albedo than non-forested areas (Bonan, 2008). As seen in Fig. 1, as the Earth cools from present day, both forested model areas (EF and TF zones) contract while the non-forested model area (GSD zone) expands slightly, in part in response to drier conditions (Gerber et al., 2004). This would lead to higher albedo and a positive feedback on the cooling. For completeness in our new treatment of the role of the land biosphere in climate and to capture such albedo variations within the context of our new land biosphere module, we assume that $a$ in Eq. (3) may be related to vegetation type such that

$$
a = 0.3 - \gamma \cdot \left( 1 - \frac{\text{frac}(\delta T_{\text{glob}})}{\text{frac}_0} \right),
$$

(4)

where the factor 0.3 is the present-day value of $a$, the $\gamma$ is a multiplier, the value of which is determined by calibration (see below), “frac” is the ratio of the area of the GSD zone to the total non-snow/ice-covered area (i.e. the sum of the areas of the EF, GSD and TF zones) and $\text{frac}_0$ is this ratio for present day. Note that $\text{frac}(\delta T_{\text{glob}})$ can be taken from Fig. 1 or calculated explicitly using Eqs.(3) and (4) and the snow-line/ice-sheet dependence on $\delta T_{\text{glob}}$ (see Fig. S3 in the Supplement). Figure 2a shows a plot of $\text{frac}(\delta T_{\text{glob}})/\text{frac}_0$.

The vegetation albedo forcing for the LGM ($\delta T_{\text{glob}} = -3.5^\circ$) relative to present day has been determined in more complex models from which we choose the value of $-0.7$ W m$^{-2}$ as being representative (Köhler et al., 2010). Together with Eq. (4) and the model latitudinal distribution of solar forcing, we find that this LGM vegetation albedo forcing anomaly is obtained in our model simulation for a $\gamma$ value of 0.02, a value we adopt here. Figure 2b illustrates new albedo distributions with latitude for the specific cases of $\delta T_{\text{glob}} = -4$, 0 and $4^\circ$ C for which $a = 0.3027$, 0.3 and 0.2976, respectively.

### 2.3 Extension of the carbon flux equations

In the original version of the DCESS terrestrial biosphere module (Shaffer et al., 2008), the global vegetation NPP is determined by

$$
\text{NPP} = \text{NPP}_{\text{PI}} \left( 1 + f_{\text{CO}_2} \cdot \ln \left( \frac{p_{\text{CO}_2}}{p_{\text{CO}_2,\text{PI}}} \right) \right),
$$

(5)

Now, we subdivide this equation into three equations:

$$
\text{NPP}_{\text{TF}} = \text{NPP}_{\text{TF,PI}} \cdot A_{\text{TF}} \cdot \left( 1 + f_{\text{CO}_2} \cdot \ln \left( \frac{p_{\text{CO}_2}}{p_{\text{CO}_2,\text{PI}}} \right) \right),
$$

(6)
NPP_{GSD} = NPP_{GSD,PI} \cdot \Delta_{GSD} \cdot \left(1 + f_{CO_2} \cdot \ln \left( \frac{pCO_2}{pCO_2,PI} \right) \right) \quad (7)

and

NPP_{EF} = NPP_{EF,PI} \cdot \Delta_{EF} \cdot \left(1 + f_{CO_2} \cdot \ln \left( \frac{pCO_2}{pCO_2,PI} \right) \right) \quad (8)

for the different vegetation zones, respectively. Thus, the global NPP is now determined by the sum of the NPP of the three vegetation zones:

NPP = NPP_{TF} + NPP_{GSD} + NPP_{EF}. \quad (9)

The factors $A_{TF}$, $A_{GSD}$ and $A_{EF}$ are calculated by

$A_{TF} = \frac{\sin(L_{TF-GSD})}{\sin(L_{TF-GSD,PI})},$ \quad (10)

$A_{GSD} = \frac{\sin(L_{GSD-EF} - L_{TF-GSD})}{\sin(L_{GSD-EF,PI} - L_{TF-GSD,PI})}$ \quad (11)

and

$A_{EF} = \frac{\sin(L_s) - \sin(L_{GSD-EF})}{\sin(L_s,PI) - \sin(L_{GSD-EF,PI})}$ \quad (12)

and scale the contributions of the respective NPP by the current area of the individual vegetation zone. The index PI stands for reference PI conditions and $f_{CO_2}$ for the CO$_2$ fertilization factor. In the original configuration, this factor was set to 0.65, which was in good agreement with results by Friedlingstein et al. (2006). However, a revision of this value in a model intercomparison study yielded a lower value of 0.37 to be a more suitable value for the terrestrial biosphere (Zickfeld et al., 2013; Eby et al., 2013), and this has also been used in the present study. Analogously, the land biosphere methane production (LBMP) (see Shaffer et al., 2008) is now calculated separately for the three vegetation zones as well.

Now, the four conservation equations per carbon isotope ($^{12}C, ^{13}C, ^{14}C$) (see Shaffer et al., 2008) have to be calculated for each vegetation zone separately. The losses for reservoir size of litter and soil were dependent on the mean global atmosphere temperature in Shaffer et al. (2008) for the uniform vegetation. In order to achieve a more realistic dependence of this process in the three vegetation zone scheme, we now approximate a mean atmosphere temperature for each vegetation zone separately by making use of the DCESS model latitudinal temperature profile expressed as a second-order Legendre polynomial in sine of latitude (Shaffer et al., 2008). This yields

$T_{DF} = \frac{(T_{atm,LL} - 0.5 \cdot T_{atm,HL}) \cdot \sin(L_{DF-GSD}) + 0.5 \cdot T_{atm,HL} \cdot \sin(L_{DF-GSD})^2}{\sin(L_{DF-GSD})}.$ \quad (13)

Here, $T_{atm,LL}$ denotes the mean atmospheric temperature in the DCESS model low/mid-latitude sector (0–52°) and $T_{atm,HL}$ in the model high-latitude sector (52–90°). $L_{DF-GSD}$ stands for the minimum of the latitude of the snow and the ice sheet line (see next section). Now, $\lambda_Q$, which influences the decay of litter and soil, can be calculated for each vegetation zone separately with $\lambda_Q(I) = Q(I)^{10}/T_{atm}$, where the index $i = 1, 2, 3$ stands for the three vegetation zones TF, GSD and EF. The conservation equations for the land biosphere reservoirs of $^{12}C$ from Shaffer et al. (2008) for leaves ($M_G$), wood ($M_W$), litter ($M_D$) and soil ($M_S$) are thus split into 12 equations, four for each vegetation zone:

$\frac{dM_G^T}{dt} = \frac{35}{60} \cdot NPP^T \cdot \frac{M_G^T}{M_G^T},$ \quad (16)

$\frac{dM_W^T}{dt} = \frac{25}{60} \cdot NPP^T \cdot \frac{M_W^T}{M_W^T},$ \quad (17)

$\frac{dM_D^T}{dt} = \frac{35}{60} \cdot NPP^T \cdot \frac{M_D^T}{M_D^T} + \frac{20}{60} \cdot NPP^T \cdot \frac{M_W^T}{M_W^T} - \frac{55}{60} \cdot NPP^T \cdot \lambda_Q \cdot \frac{M_D^T}{M_D^T},$ \quad (18)

$\frac{dM_S^T}{dt} = \frac{5}{60} \cdot NPP^T \cdot \frac{M_W^T}{M_W^T} + \frac{10}{60} \cdot NPP^T \cdot \lambda_Q \cdot \frac{M_D^T}{M_D^T} - \frac{15}{60} \cdot NPP^T \cdot \lambda_Q \cdot \frac{M_S^T}{M_S^T}.$ \quad (19)

Analogously, these equations are extended for the rare carbon isotopes $^{13}C$ and $^{14}C$, where fractionation factors for land photosynthesis and, for $^{14}C$, radioactive sinks are considered (Shaffer et al., 2008). The flux of carbon dioxide between the terrestrial biosphere and the atmosphere is then determined by

$F_{CO_2} = \sum_{i=1}^{3} -NPP^T \cdot \frac{45}{55} \cdot NPP^T \cdot \lambda_Q \cdot \frac{M_D^T}{M_D^T} + \frac{15}{60} \cdot NPP^T \cdot \lambda_Q \cdot \frac{M_S^T}{M_S^T}.$ \quad (20)
The flux is then given by

\[ F_{13,14} = \frac{3}{13} \cdot \frac{\Delta M_{13,14} \cdot \Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Delta t} \cdot \frac{\Delta M_{13,14}}{\Delta t} \cdot \frac{\Delta M_{12,14}}{\Δ\delta
LGM reconstructions show less carbon in the land biosphere than for warmer, pre-industrial conditions (Peng et al., 1998; Prentice et al., 2011). This simplistic model behaviour can be seen in Fig. 3a, which shows the steady-state terrestrial biomass as a function of $\rho CO_2$ and $T_{glob}$. These results are generated through prescribing various $\rho CO_2$ and $T_{glob}$ values in numerous 2 kyr model simulations.

In the new version (Fig. 3b), biomass decreases when temperatures sink as vegetation types shift and the snow line moves equatorward (note, however, that a prescribed ice sheet line is not included in these simulations). The permafrost biomass, however, increases in the course of that process. The figure also shows that further cooling only slightly reduces the land biosphere carbon storage. This shows that the general land carbon storage is represented more realistically in the new model version.

Furthermore, we show the response of the model vegetation zones and the different vegetation reservoirs to a reduction of atmospheric temperatures and $\rho CO_2$ to LGM conditions and compare the results with complex vegetation models as well as with data reconstructions. To evaluate the vegetation scheme for LGM conditions, we carried out cooling simulations with the new and with the old biosphere scheme. For these, we started from a PI steady state ($T_{glob} = 15^\circ C$, $\rho CO_2 = 280$ ppm), but prescribed the global mean temperature to $T_{glob} = 11.5^\circ C$ (see Shakun et al., 2012) and the atmospheric $\rho CO_2$ concentration to 190 ppm (e.g. Monnin et al., 2001). A third cooling simulation was conducted with the new biosphere scheme and conditions as described but with an additional prescription of the ice sheet line to 47$^\circ$ latitude. Figure 4a shows the global sum of total land biomass (LB) carbon (without carbon stored in permafrost) for the three cooling simulations, as well as LB carbon for the three individual vegetation zones, and Fig. 4b shows LB carbon of the vegetation reservoirs above ground (leaves + wood) and below ground (litter + soil) for the three
simulations. Since the system seems to be in equilibrium after around 1 ka, we integrated these simulations over 2 ka.

As already presented in Fig. 3, the cooling experiment again demonstrates that LB carbon increases in the old model version and decreases with the new biosphere scheme. Figure 4b shows that the unrealistic increase of LB carbon is due to an increase in litter/soil carbon (i.e. biosphere below ground). In the simulation with the new biosphere scheme, this does not happen. The EF zone is dominated by biosphere below ground and due to the limitation of the poleward expansion of the EF zone through the snow line, this carbon reservoir is now decreasing. Also, the figures show that the timing of the change is represented more nuanced with the new biosphere scheme. The biospheres in the three vegetation zones show different reaction times according to their distinct temperatures and the dominating pool of vegetation in the respective area. When we also include the expansion of ice sheets ($L_{\text{ice}} = 47^\circ$; for explanation see Sect. 3.1), covering larger areas of the EF vegetation zone than the snow line, the total land carbon pool decrease is stronger (dotted lines). It is mainly the biosphere below ground, exclusively in the EF zone, that accounts for this.

A poleward limitation of the biosphere in the old vegetation scheme also leads to a reduction of LB carbon in the cooling simulation. To confirm this, we performed an additional simulation with the old vegetation scheme, but with the crude vegetation area limiting approach $A = (\sin(L_{\text{snow}}) / \sin(L_{\text{snow},\text{pt}}))^2$. In this cooling experiment the total LB carbon decreases, but not as much as with the new biosphere scheme, and the decrease happens faster than with the new biosphere scheme (not shown). The LB change in the EF zone mainly depends on variations in soil, which has a slow response time and is the largest biomass reservoir (Fig. 4a). The TF zone adapts much more quickly to the new climate conditions because in this vegetation type the biomass is dominated by leaves and wood. This shows that not only the quantitative but also the temporal description of land biomass changes is represented more accurately now. The GSd vegetation zone shows the smallest change in biomass, because in the cooling simulations the area of this vegetation zone changes only slightly, rather shifting just latitudinally.

We calibrated the latitudinal dependence of the vegetation zone borders to match the LPJ model results. However, the calculation of carbon stored in the terrestrial biosphere at different climate conditions also depends on other parameters. Hence, we also evaluate the performance of the new DCESS vegetation scheme by comparing it to the results of the LPJ model study by Gerber et al. (2004). For this, Table 2 shows the percentage change of biomass carbon in the cooling experiment for the new vegetation scheme with and without ice sheet prescription and for the old vegetation scheme with and without the biosphere area limit (old bio plus) as well as for the LPJ model.

### Table 2. Percentage change of biomass carbon in the cooling experiment for total biomass and divided into reservoirs above and below ground. DCESS model with the old biosphere scheme, with and without the crude approach for vegetation area limiting (see text), and with the new biosphere scheme with and without prescribed ice sheet expansion, and LPJ model study presented in Gerber et al. (2004).

<table>
<thead>
<tr>
<th>ALB/%</th>
<th>Total</th>
<th>Litter + soil</th>
<th>Leaves + wood</th>
</tr>
</thead>
<tbody>
<tr>
<td>Old bio</td>
<td>+3.0</td>
<td>+9.5</td>
<td>-14.5</td>
</tr>
<tr>
<td>Old bio plus</td>
<td>-10.8</td>
<td>-5.2</td>
<td>-26.0</td>
</tr>
<tr>
<td>New bio</td>
<td>-18.0</td>
<td>-14.2</td>
<td>-28.5</td>
</tr>
<tr>
<td>New bio ice</td>
<td>-27.6</td>
<td>-25.3</td>
<td>-33.6</td>
</tr>
<tr>
<td>LPJ</td>
<td>-24.8</td>
<td>-24.7</td>
<td>-25.0</td>
</tr>
</tbody>
</table>

This comparison demonstrates that in relation to the LPJ model, the adaptation of the LB to different climate conditions is captured much better with the new biosphere scheme. While biomass carbon increased with the old model version, the new biosphere scheme produces most of the change that the LPJ model shows. Most of the improvement in LB variations through the new vegetation scheme is due to the snow line, that limits the poleward expansion of the biosphere. Using the old biosphere with additional vegetation area limitation, LB carbon decreases under LGM climate conditions. However, with the new vegetation scheme, the snow line particularly limits the EF zone and this largely improves the overall representation of biomass below ground. When vegetation area reduction is applied to the old biosphere module, the biomass change above ground was already in good agreement with the LPJ model. Hence, the reason for the much larger changes in overall biomass between the old and the new model version as shown in Fig. 4 is mainly due to the better representation of the slow change of the soil biomass in the EF zone. This more accurate representation of soil in the EF zone, however, is also due to the fact that now the biomass reservoir of each vegetation zone depends on the specific temperature of the zone in question and not on the global mean temperature as in the old model version. The prescribed ice sheet line at $47^\circ$ latitude generates a further drawdown of the land carbon stock. The percentage change is then close to the LPJ model, about 3 % higher. Vegetation above ground changes too much, although this type of vegetation is not affected much by the ice line (see Fig. 4b), but due to the low total amount, the percentage change is high.

Peng et al. (1998) provide an overview of various studies that estimate the reduction in global land biomass for the LGM compared to present day. Those are based on either global circulation model (GCM) simulations, marine carbon isotope data changes or vegetation mapping approaches. Altogether, these studies show a large spread from 0 (Prentice and Fung, 1990) to $-1350\text{ Gt C}$ (Adams et al., 1990). The majority of the studies show values between $-300$ and
700 Gt C; a more recent modelling study by Prentice et al. (2011) provides values of −550 to −694 Gt C. Through the implementation of the new vegetation scheme, the DCESS model biomass carbon change between PI and LGM improves from +43 to −408 Gt C. Thus, results with the new model version agree well with other estimates, albeit at their low end. Carbon stored in permafrost is around 600 Gt for PI conditions and around 1000 Gt for LGM climate when the ice sheets are included. Hence, the total amount of carbon on land is about 2800 Gt for either climate state. Ciais et al. (2012) estimate the LGM global carbon stock to be 3640 ± 400 Gt, so somewhat higher than in our model. Partly this is due to our estimation of 30 Gt m$^{-2}$ for permafrost from Schuur et al. (2015). This apparent underestimation of the permafrost carbon inventory has to be kept in mind when analysing the model results and it will be addressed in the following with a sensitivity experiment using 60 Gt C m$^{-2}$ for permafrost, for which the total amount of carbon on land will be about 3800 Gt.

Overall, it can be stated that the new biosphere scheme with the three vegetation zones constitutes a significant improvement for the representation of the terrestrial biomass as well as the estimates of the size and timing of carbon exchanges between the terrestrial biosphere and the atmosphere. This new implementation better captures the complex interactions between the terrestrial and the atmospheric carbon exchange, as is required for a better understanding of the processes that determine climate changes on glacial-interglacial timescales.

3 Application to LGM and deglaciation

As a first application of the new DCESS terrestrial biosphere module, we simulate the deglaciation after the LGM, when global atmospheric temperatures rose by around 3.5 K (Shakun et al., 2012) and atmospheric pCO$_2$ increased from 190 ppm during the LGM to Holocene conditions of 260 ppm in a series of steps (e.g. Monnin et al., 2001). The most marked of these steps is a steep 38 ppm rise near the onset of the deglaciation, the MI (Broecker and Barker, 2007). In the Supplement, we provide a literature review with details about the MI, including current hypotheses for the explanation of that climate change. Earlier studies found considerably greater LGM global mean cooling (Schneider von Deimling et al., 2006); recent estimates based on much improved temperature data, however, have shown LGM cooling of 3.2–4 K (Schmittner et al., 2011; Shakun et al., 2012; Annan and Hargreaves, 2013).

A complete explanation for the pCO$_2$ and temperature increase at the onset of the last glacial termination must be able to reproduce a simultaneous decrease by 0.3 and 160‰ of atmospheric $\delta^{13}$C (Schmitt et al., 2012) and $\Delta^{14}$C (Reimer et al., 2013), respectively. Furthermore, it should also include how LGM deep water with high salinity (Adkins et al., 2002), low $\delta^{13}$C (Curry and Oppo, 2005) and $\Delta^{14}$C (Burke and Robinson, 2012) and low dissolved oxygen concentrations (but not widespread anoxia) (Jaccard et al., 2014) was formed during the last glacial. Hence, it requires the consideration of a globally comprehensive picture of the physical and biogeochemical processes in atmosphere, ocean and on land, as well as their interactions on various timescales. With its new biosphere scheme, the DCESS model is now better suited for investigations of that kind. However, a number of further adaptations need to be made to simulate LGM conditions and the transition to the Holocene. These are presented next, followed by transient simulations across the last 25 kyr BP. For these, the model was initialized and forced with the conditions described in Sect. 3.1. Since we focus on the MI (17.5–14.5 ka BP), we mainly present and discuss the time period from 20 to 10 ka BP. We assess the impact of various processes on the overall climate change with a focus on the new biosphere scheme and permafrost. In the process, we also evaluate proposed time series for the production of $^{14}$C in the atmosphere.

3.1 Model LGM and transition

Guided by proxy data records, we first modified several biogeochemical and physical parameters to generate a model steady state that represents the LGM well. For this, a number of parameters can be considered as possible candidates (see e.g. Kohfeld and Ridgwell, 2009). However, under consideration of the possibilities provided by the enhanced model and knowledge about candidate parameters, we decided upon the adaptations described below.

Increased iron supply and thereby ocean fertilization (Martin et al., 1990) through enhanced atmospheric dust concentrations during the LGM (see e.g. Mahowald et al., 1999, 2006b; Maher et al., 2010), particularly in the high southern latitudes (e.g. Lambert et al., 2013, 2015), probably led to enhanced new production of organic matter in the Southern Ocean (SO) by way of iron fertilization (see also Lamy et al., 2014; Martínez-García et al., 2014). To account for this, we modified the efficiency factor for new production in the model high-latitude ocean sector from 0.36 (standard value for PI conditions; see Shaffer et al., 2008) to 0.5. This leads to a reasonable productivity increase of around 40% for the area of the SO and induces an atmospheric pCO$_2$ reduction of around 20 ppm, consistent with the DCESS model iron fertilization results in Lambert et al. (2015). Moreover, an additional radiative effect of −1 W m$^{-2}$ (Mahowald et al., 2006a) for glacial conditions through atmospheric dust during the LGM is considered. For the transient simulations from the LGM to the Holocene, we have developed a transfer function between temperature and dust fluxes from proxy data records that we applied to the efficiency factor and to the radiative effect. It yielded an exponential dependence of dust with temperature; details can be found in the Supplement.
The lower sea level during the LGM (around 130 m; see e.g. Waelbroeck et al., 2002; Lambeck et al., 2014) and a thereby reduced ocean volume by around 3.5% (see e.g. Adkins and Schrag, 2002) is accounted for by increasing phosphate concentrations (the nutrient-limiting source in the DCESS ocean biochemistry) and the ocean salinity (see Adkins et al., 2002) by 3.5%. For the transition of these parameters across the last 25 ka BP, we use the latest sea level reconstruction time series from Lambeck et al. (2014). We do not account for the expansion of land mass and vegetation due to reduction of sea level, which causes additional carbon storage (Joos et al., 2004). Although Joos et al. (2004) found that this effect is less important than the effect through climate/CO$_2$-caused vegetation changes or the ice sheet area effect, it can still have a considerable impact in deglaciation simulations and should be kept in mind when evaluating results. To generate LGM conditions for $\Delta^{14}$C in atmosphere and ocean, we applied the average cosmogenic $^{14}$C production rate from 25 to 26 ka BP (PR$_{14C} = 2.1 \times 10^{5}$ atoms cm$^{-2}$ s$^{-1}$). For this and in most of the transient simulations, we use the most recent production rate time series developed by Hain et al. (2014). In a sensitivity analysis, the $^{14}$C production rates from the studies by Laj et al. (2004) and Muscheler et al. (2004) are applied as well. A description of the main characteristics of these data is given in the Supplement.

LGM climate reconstructions show that the Laurentide ice sheet expanded as far south as 38° N (see e.g. Peltier, 2004). To account for this and the lack of an ice sheet in large parts of Siberia, and within the constraints of our zonally averaged one-hemisphere model, we prescribe the southernmost ice sheet extent to be 47°. For the transient simulations we impose the temporal retreat of the ice line to the disappearance of the ice sheets at 70° latitude during the Holocene. For this, we linearly prescribe $L_{\text{ice}}$ (see Sects. 2.3 and 2.4) to a data set presented in Shakun et al. (2012) showing the Northern Hemisphere ice sheet expansion from 100% (ice line at 47°) at the LGM to 0% (ice line at 70°) at present day. An example case for $L_{\text{ice}}$ and $L_{\text{snow}}$ in a transient simulation is given in the Supplement.

A model analogy to isolated deep water in the SO (see e.g. Watson and Naveira Garabato, 2006) is generated through application of a depth-dependent function for vertical exchange intensity in the high-latitude ocean sector. For this, we impose a sharp decrease in vertical diffusion at around 1800 m ocean depth which limits mixing of the upper ocean layers with intermediate and deep ocean waters. The transition depth of this profile was varied to obtain LGM climate conditions that constrain all required oceanic and atmospheric variables. Through the application of this diffusion profile, the isolated ocean waters below the transition change towards high dissolved inorganic carbon (DIC) and alkalinity values as well as towards low oxygen concentrations and $^{13,14}$C isotope ratios. This variation in vertical exchange intensity should not be understood as a change in real oceanic vertical diffusion, but rather as a model analogy for LGM conditions of the SO that were likely due to some combination of weakened or equatorward-shifted westerly winds (Toggweiler and Russel, 2008; Anderson et al., 2009; d’Orgeville et al., 2010) and increased stratification through brine-induced effects (Bouttes et al., 2010, 2011; Mariotti et al., 2013). With its wide latitudinal extent and the land bounding poleward of 70°, the high-latitude ocean sector of the DCESS model bears considerable resemblance to the SO. During the transient simulations, we slowly restore this modification back toward PI conditions between 17.5 and 14.5 ka BP to apply the entire effect of this process to the MI. In this process, deeper layers in the high-latitude ocean sector are again brought into contact with surface layers promoting outgassing, and ocean profiles go back toward the initial PI state shown in Shaffer et al. (2008). An illustration of the profile as well as a detailed technical description of the procedure and some additional information are presented in the Supplement.

When all these adaptations, plus a few minor changes (described in the Supplement), are applied, an 80 ka DCESS simulation leads to a steady climate state with conditions close to data-based LGM reconstructions. Atmospheric $p$CO$_2$ decreases to 187.9 ppm and the global mean atmosphere temperature to 11.70°C. For $p$CO$_2$, proxy data records by Lüthi et al. (2008) provide a range of 186–198 ppm and Shakun et al. (2012) present LGM global mean atmosphere temperatures between 11.5 and 11.8°C. Moreover, atmospheric isotope ratios of $\delta ^{13}$C = −6.41‰ and $\Delta ^{14}$C = 414.5‰ and low oxygen values but no widespread anoxia in the deep ocean are achieved. This agrees well with proxy data records presented by Schmitt et al. (2012), Reimer et al. (2013) and Jaccard et al. (2014). An overview of these data and the ocean profiles for LGM conditions of various variables for the high- and the low/mid-latitude sector are shown in the Supplement. In the following sections, we present analyses of the transient simulations from the LGM to the Holocene, using the transition functions described above.

3.2 Transient simulation results

To evaluate the impact of the individual new model developments, we carried out six transient simulations starting from LGM conditions as described above, varying the following parts of the model land biosphere: the first simulation features a null-vegetation model (Nul_veg), meaning that the vegetation does not change from LGM conditions and land–atmosphere carbon (including the rare isotopes) fluxes are suppressed. Then, we use the old uniform land biosphere scheme (Old_bio) (without the snow/ice line-based reduction of biosphere area; see Sect. 2.5, and no permafrost parameterization), and subsequently the new scheme without the permafrost (permafrost–atmosphere carbon fluxes set to zero) and the new albedo features (NoPF_alb) and...
Table 3. Overview of the DCESS model simulations with short descriptions.

<table>
<thead>
<tr>
<th>Simulation</th>
<th>Long name</th>
<th>Setup</th>
</tr>
</thead>
<tbody>
<tr>
<td>Nul_veg</td>
<td>Null vegetation</td>
<td>Suppressed land–atmosphere fluxes</td>
</tr>
<tr>
<td>Old_bio</td>
<td>Old biosphere</td>
<td>Original uniform DCESS land biosphere scheme (no permafrost and no land area change)</td>
</tr>
<tr>
<td>NoPF_alb</td>
<td>No permafrost</td>
<td>Suppressed fluxes from permafrost and old (not vegetation-dependent) albedo</td>
</tr>
<tr>
<td>NoPF</td>
<td>No permafrost</td>
<td>Suppressed fluxes from permafrost</td>
</tr>
<tr>
<td>REF</td>
<td>Reference</td>
<td>Including all new developments as described in the text</td>
</tr>
<tr>
<td>2xPF</td>
<td>Doubled permafrost</td>
<td>As REF but with two times the estimate for permafrost carbon reservoir</td>
</tr>
</tbody>
</table>

Figure 5. Atmospheric values for the DCESS simulations with null vegetation model (Nul_veg, red line with dots), old biosphere scheme (Old_bio, red line), deactivated permafrost component and old albedo scheme (NoPF_alb, blue line with dots), deactivated permafrost component and new albedo scheme (NoPF, blue line), reference simulation with all new components (REF, light blue line with dots) and sensitivity experiment with doubled permafrost carbon reservoir (2xPF, light blue line) and data-based reconstructions (black). \( p\text{CO}_2 \) by Lüthi et al. (2008), temperatures by Shakun et al. (2012), \( \delta^{13}\text{C} \) by Schmitt et al. (2012) and \( 1^{14}\text{C} \) by Reimer et al. (2013).

then the same but including the new albedo (NoPF). Last, we performed simulations with all the new model developments (REF) plus a further sensitivity experiment with a doubled (60 Gt C m\(^{-2}\)) permafrost carbon reservoir, as already mentioned. An overview of these simulations is provided in Table 3.

The results of these model simulations as well as data-based reconstructions are presented in Fig. 5 from 20 to 10 ka BP. As our transition functions (in particular the upwelling of the deep ocean) are tailor-made for simulating the MI between 17.5 and 14.5 ka BP, we particularly focus on these 3 years of the last glacial termination in the analysis.

The null vegetation model shows the largest atmospheric changes across the MI. Uptake of carbon through the land biosphere does not take place in this simulation; therefore, all outgassed carbon stays in the atmosphere and amplifies global warming. This also reflects in the \( \delta^{13}\text{C} \) and \( 1^{14}\text{C} \) curves – isotopically strongly depleted carbon from the deep ocean decreases the atmospheric isotope ratios. Especially the far too strong drop in \( \delta^{13}\text{C} \) in the null vegetation sim-
ulation indicates that the regrowth of the biosphere and its preferential uptake of $^{12}$C keeps $\delta^{13}$C at a reasonable level in the other simulations, although the increase after 12 ka BP is not represented well in the model. The simulation with the old land biosphere scheme shows rather small changes across the MI; the expansion of the biosphere leads to uptake of atmospheric carbon. Due to its reaction on vegetation changes, the new albedo diversification leads to stronger warming. This also generates some stronger $p$CO$_2$ increase. When we enable the permafrost parameterization in the REF simulation, $p$CO$_2$ rises by around 2.6 ppm more and the global mean atmosphere temperature by around 0.1 $^\circ$C. The results of the simulations start diverging at around 19 ka BP. This is when the change in ice sheet extent leads to first clear variations through its effect on the permafrost parameterization in the model (see Fig. S3). The isotope ratios are only slightly affected by these new features; in particular, $\Delta^{14}$C is controlled mostly by the changes of the stratospheric production rate of $^{14}$C. The sensitivity experiment with a doubled permafrost reservoir shows a further increase of $p$CO$_2$. The difference between the 2xPF and the REF simulation is larger than between the REF and the noPF simulation. The biosphere regrowth and its carbon uptake is only slightly enhanced in the 2xPF simulation. However, some more change already happens before, i.e. after 19 ka BP. Therefore, this shows that uncertainties of that kind can have a considerable impact on climate change simulations. In comparison to data-based reconstructions, the MI atmospheric changes are closest in the 2xPF simulation (disregarding the Nul_veg simulation). More than half of the $p$CO$_2$ and the global mean temperature changes are represented and the drop in $\delta^{13}$C is almost reached. $\Delta^{14}$C shows only little sensitivity to our new model developments.

Figure 6 shows the changes of permafrost carbon, land biosphere carbon and their sum for the REF and the 2xPF simulations. In the REF simulation, carbon uptake through the regrowth of the biosphere across the MI slightly exceeds (by 70 Gt C) carbon outgassing through ice sheet retreat and permafrost thawing then. In the 2xPF simulation, the permafrost carbon change slightly outweighs the vegetation effect. This demonstrates that the two mechanisms broadly compensate for each other and provide an estimate of its uncertainty. As mentioned above, the land biosphere carbon reservoir change in the REF simulation is at the low end of the range found in other studies (Peng et al., 1998; Prentice et al., 2011). Also, model carbon release of 337 Gt C from permafrost is lower than that of Ciais et al. (2012), who found a 700 Gt C difference between LGM and present-day global permafrost carbon reservoir. Our lower estimate seems to be related to our simplified permafrost treatment and the simple assumption of 30 kg of available carbon per square metre of permafrost-covered area (Schuur et al., 2008). The sensitivity simulation with 60 kg C m$^{-2}$ in permafrost provides more realistic values for permafrost carbon release (667 Gt C) and also for the global carbon reservoir ($\sim$ 3800 Gt C; see also Sect. 2.5).

Additionally, we conducted four transient simulations to assess the impacts of the individual transition functions on atmospheric $T_{glob}$, $p$CO$_2$, $\delta^{13}$C and $\Delta^{14}$C changes (see Supplement). The transition functions described above were applied sequentially to better assess the impact of each process. These simulations show that during the 3 ka of the MI, most of the simulated changes can be attributed to the resumption of the ocean high-latitude vertical diffusion and the thereby induced outgassing of the carbon-rich and isotopically depleted deep waters. Our DCESS simulations reproduce only some aspects of the early last deglaciation, while others are underestimated because important processes are either missing or not adequately represented.

As has been mentioned, the change in $\Delta^{14}$C during the MI in the REF simulation is not as large as in the data-based reconstructions. Apart from atmospheric $^{14}$CO$_2$ itself and the release of deep ocean waters, $\Delta^{14}$C is strongly influenced by the cosmogenic production rate of $^{14}$C. This production rate is determined with rather large uncertainties and there are different ways to derive it. In the Supplement, we present the three $^{14}$C production rate time series of the studies by Laj et al. (2004), Muscheler et al. (2004) and Hain et al. (2014) across the last 25 kyr BP. Here, we present an evaluation of the three $^{14}$C production rate data applied to the ALL_TF simulation. In Fig. 7, we show the simulations with the three different production rates, as well as for a simulation with constant LGM-value production rate (Mus_PR, Muscheler et al., 2004 production rate; Laj_PR, Laj et al., 2004 production rate; LGM_PR, constant LGM-value production rate). The proxy data record by Reimer et al. (2013) is also included in the figure.

The simulation with constant $^{14}$C production rate at LGM level shows a $\Delta^{14}$C drop by 80% from the beginning to the end of the MI, almost entirely through the outgassing of isotopically depleted deep ocean waters. Neither of the $^{14}$C pro-
duction rates can account for the remaining 80\% reduction to explain the $\Delta^{14}C$ decrease of 160\% across the MI that can be seen in the data-based reconstruction by Reimer et al. (2013). With the data set by Hain et al. (2014), $\Delta^{14}C$ drops by 96\%, using the Laj et al. (2004) data, a 105\% decrease can be explained, and the Muscheler et al. (2004) time series only leads to −58\% change. Furthermore, the proxy data do not show the production-rate-caused variations within the MI, and also, in the Mus PR simulation, atmospheric $\Delta^{14}C$ shows a large and sudden drop of around 150\% shortly after the MI between 14.3 and 13.7 ka BP.

### 3.3 Discussion of transient simulations

The model reproduces more than half of the MI changes in atmospheric $pCO_2$, $T_{\text{glob}}$, $\delta^{13}C$ and $\Delta^{14}C$ as shown in data-based reconstructions. Overall, the representation of the land biosphere is shown to play an important role in the interplay of many processes. The model results reach from 12 to 31 ppm change in $pCO_2$ across the MI, i.e. from less than a third of the change presented in data-based reconstructions to more than 80\%. The “best” results are reached with the least complex vegetation model version, unambiguously for the wrong reasons. The missing uptake of carbon through the land biosphere leads to too-high $pCO_2$ and temperature values. The $\delta^{13}C$ isotope ratios reveal this model deficiency. $\delta^{13}C$ further decreases in the Nul_veg simulation after the MI, while in all other simulations, $\delta^{13}C$ stagnates. In the data-based reconstructions, $\delta^{13}C$ even rises again. Schmitt et al. (2012) mainly attribute this rise to the continuing growth of the land biosphere, which does not have such a strong effect on atmospheric $\delta^{13}C$ in the model. According to Crichton et al. (2016), peatlands could also account for this effect – however, those are not included in our vegetation scheme. When we apply the doubling of the estimate of 30 Gt C m$^{-2}$ from Ciais et al. (2012), the model results considerably improve in comparison with data-based reconstructions. In consideration of the apparent underestimation of total land biosphere carbon as shown in Sect. 2.5 and the large uncertainties in the estimation by Ciais et al. (2012), the usage of 60 Gt C m$^{-2}$ is still reasonable.

The impact of the land biosphere on $\Delta^{14}C$ is very small, even though we assume carbon released from permafrost to be radiocarbon free. The expected radiocarbon decrease generated through permafrost thawing can apparently be compensated for by ocean–atmosphere exchange and subsequent mixing to the deeper ocean. It has to be considered that the carbon buried below permafrost seems to be underestimated in our model approach compared to a study by Ciais et al. (2012) and that interhemispheric seesaw effects can affect the timing of extensive permafrost ($^{14}C$ depleted) carbon release, especially during HE1 (see e.g. Köhler et al., 2014). The much discussed sharp $\Delta^{14}C$ drop of 160\% (see Reimer et al., 2013) (note that in previous studies by Broecker and Barker, 2007, or Reimer et al., 2009, this was referred to as 190\%) at the early stages of the last deglaciation is not entirely reproduced by this modelling study. By applying a constant LGM $^{14}C$ production rate, all the above-described processes can account for about 70\% change. None of the three different time series of the $^{14}C$ production rate can account for the rest of the $\Delta^{14}C$ change. At most, the data of Laj et al. (2004) lead to an additional 25\% decrease. However, the determination of the $^{14}C$ production rate is obviously subject to large uncertainties. For example, the drop in the Muscheler et al. (2004) time series at around 14 ka BP leads to a sudden 150\% decrease in $\Delta^{14}C$ in our model simulation but cannot be seen in $\Delta^{14}C$ proxy data. In this context, it should be mentioned, that recent revisions to ice core timescales have not yet been applied for revising the reconstructed snow accumulation rates and $^{10}Be$ fluxes and their influence on the $^{10}Be$ based $^{14}C$ production rate (R. Muscheler, personal communication, 2015).

Most of the MI changes are caused by the upward transport of carbon-rich and isotopically depleted waters from the deep ocean through prescription of the vertical diffusion profile and its resumption. The dust component accounts for about 0.3 \degree C global temperature change during the MI. Since the other atmospheric quantities are only moderately affected by dust, most of that can be related to the direct dust radiative forcing. To account for the other half of changes that our simulations cannot reproduce, several processes can be thought of as being insufficiently represented in the model. Moreover, this could be due to the timing of one or more of the transition functions, under-representing effects during the MI. Brovkin et al. (2007), Kohfeld and Ridgwell (2009) and Mariotti et al. (2013) discuss a number of processes that, combined, can account for the entire deglaciation, although with sometimes large uncertainties. Not all of them were captured in our study. For example, enhanced ocean remineralization length scales during the glacial, due

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**Figure 7.** $\Delta^{14}C$ in transient simulations with all changes (see Sect. 3.1) applying different $^{14}C$ production rates: from Hain et al. (2014) (red, REF), Muscheler et al. (2004) (magenta, Mus PR), Laj et al. (2004) (blue, Laj PR) and fixed LGM production rate (cyan, LGM PR) and data-based reconstructions from Reimer et al. (2013) (black).
to less active bacteria at low temperatures, could trap more DIC in the deep ocean, which then could account for additional CO$_2$ outgassing but would also reduce deep-ocean dissolved oxygen concentrations. Also the volume of isolated deep waters in the SO is uncertain; moreover, water masses in other oceans may also have contributed to the overall atmospheric $p$CO$_2$ change (Rose et al., 2010; Okazaki et al., 2010; Kwon et al., 2012; Huiskamp and Meissner, 2012). The $T_{\text{gl}}$ and $p$CO$_2$ changes after the MI across the BA, the Younger Dryas and the Holocene are not expected to be simulated in detail by the DCESS model. Due to the model’s simplified geometry, interactions between the hemispheres and thus the bipolar seesaw cannot be represented. The simplicity of DCESS model ocean dynamics also limits feedbacks of ocean–atmosphere interactions that may have contributed to the overall carbon cycle change during the MI. For instance, Mariotti et al. (2016) discuss the effect of North Atlantic freshening through ice sheet melting inducing upper water stratification and subsequent prevention of carbon uptake by the ocean to contribute to enhanced $p$CO$_2$ during HE1 at the end of the MI. An alternative approach would be to use 3-D modelling to deal specifically with one or more of the processes listed above. However, this would involve other types of uncertainties, like the strength and position of the southern westerly winds and the parameterization of diapycnal mixing.

4 Summary and conclusions

The land biosphere scheme that accounts for $^{12,13,14}$C cycling with leaf, wood, litter and soil of the reduced complexity Earth system model DCESS has been extended to three different vegetation zones. Based on a complex land biosphere model study, we defined dynamically varying vegetation borders on a global scale that depend on temperature variations. We also introduce a parameterization that accounts for carbon, including its rare isotopes, that is being trapped below the permafrost as well as below terrestrial ice sheets for glacial conditions and released during deglaciation events. In an evaluation, the new terrestrial biosphere scheme is shown to simulate more realistic global biomass size and timing in climate change experiments, and thereby significantly improves the representation of land–atmosphere carbon exchange rates in the DCESS model. For climate change studies on glacial–interglacial timescales, these aspects can be crucial when analysing the contributions and interactions of processes controlling carbon exchange between land, atmosphere and ocean.

For a first application of the new biosphere parameterization, the model is first tuned to LGM conditions to subsequently carry out transient simulations across the last glacial termination. Along with a number of established adaptations of physical and biogeochemical parameters, the DCESS model successfully reproduces proxy data records of glacial conditions in the ocean and atmosphere when we impose the isolation of high-latitude deep ocean waters. For the transient model simulations, we have additionally developed a set of explicit functions that describe the transitions of atmospheric dust, ocean volume and terrestrial ice sheet extent across the last 25 kyr BP. These sensitivity experiments show that large parts of the exceptional change in atmospheric $p$CO$_2$, $\delta^{13}$C, $\Delta^{14}$C and $T_{\text{gl}}$ at the onset of the last glacial termination (MI, 17.5–14.5 ka BP) can be represented by this approach. Some variations as seen in data-based reconstructions cannot be reproduced by our model study. These remaining changes could possibly be captured by applying a dynamically more complex model including distinct water masses and a second hemisphere for representing bipolar seesaw effects, or by revising and/or adding one or more model parameterizations. New insights into these mechanisms can help to improve our understanding of global carbon cycle changes on centennial to millennial timescales.

The thawing of permafrost due to atmospheric warming and retreat of ice sheets, as well as the regrowth of the terrestrial biosphere, are found to play moderate, but important roles in explaining the climate change of this period of the last deglaciation. We found that these two processes broadly compensate for each other in the model in terms of CO$_2$ exchange with the atmosphere, making little net contribution to atmospheric $p$CO$_2$ changes across the last transition. However, since our simulation bears considerable uncertainties, we also found that particularly the permafrost component could be strongly underestimated. Simulations across the transition using the original DCESS land biosphere model also showed essentially no net contribution to atmospheric $p$CO$_2$ change as reflected in the very small change in land biomass between LGM and present day. But with the new biosphere module (including permafrost) this result is obtained in a more correct manner, in better agreement with proxy data and more complex modelling results.

Data availability. The basic DCESS model code is available at http://www.dcess.dk/ and all applied data are available as referenced.

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Competing interests. The authors declare that they have no conflict of interest.

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